# Hilbert's Monkey Saddle and other Curiosities All 3-(Point-)Particle Riesz Equilibria on a Circle 

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## Optimal $N$-Particle Configurations on $\mathbb{S}^{n}$

- Pick $N$ distinct points $\mathbf{p}_{k} \in \mathbb{R}^{n+1}, k \in\{1, \ldots, N\}$, with $\left|\mathbf{p}_{k}\right|=1$
- Take any pair, say $p_{i}$ and $p_{j}$.
- Denote their Euclidean distance by $r_{i j}>0$.
- Assign each ij pair a Riesz s-energy $V_{s}\left(r_{i j}\right)$, with

of an $N$-pt. configuration $\omega^{(N)}$,

- Optimal average $s$-Riesz pair energy of $N$ point particles:
$v_{s}(N):=\inf _{\omega(N)}\left\langle V_{s}\right\rangle\left(\omega^{(N)}\right)$


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\begin{array}{ll}
V_{s}(r):=\frac{1}{s}\left(\frac{1}{r^{s}}-1\right), & s \in \mathbb{R}, \quad s \neq 0 \\
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## Optimal $N$-Particle Configurations on $\mathbb{S}^{n}$ when $s \geq-2$

- An optimizer $\omega_{*}^{(N)}$ exists whenever $s \geq-2$ :

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v_{s}(N)=\left\langle V_{s}\right\rangle\left(\omega_{*}^{(N)}\right)
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- Optimizers $\omega_{*}^{(N)}\left(/ O(n+1)\right.$ and $\left./ S_{N}\right)$ generally not unique!


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- Asymptotic large $-N$ expansion for $n=2$ and $s=0$ :

$$
\frac{N(N-1)}{2} v_{0}(N)=a N^{2}+b N \ln N+c N+d \ln N+\mathcal{O}(1)
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a=\frac{1}{4} \ln \frac{e}{4}, \quad b=-\frac{1}{4}, \quad c=\ln \left(2(2 / 3)^{1 / 4} \pi^{3 / 4} / \Gamma(1 / 3)^{3 / 2}\right), \quad d ?
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For each $N$, can one find $\omega_{\dot{\alpha}}^{(N)}$ on $\mathbb{S}^{2}$ such that

$$
\left|v_{0}(N)-\left\langle V_{s}\right\rangle\left(\omega_{\infty}^{(N)}\right)\right|<D \frac{\ln N}{N(N-1)}
$$

using not more than $\operatorname{Poly}(N)$ many steps?

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## - 3 points inevitably on GREAT CIRCLE: $\mathbb{S}^{1}$

- We speak of "optimal N-particle arrangement" in this case.
- $N=2 M$ : optimizer $=M$ points each in two antipodal points
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## Optimal 3-Particle Arrangements on $\mathbb{S}^{1}$

- Average pair energy of 3-particle arrangement $\omega^{(3)} \equiv(\alpha, \beta, \gamma)$

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\begin{gathered}
\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma):=\frac{1}{3}\left(V_{s}(|a|)+V_{s}(|b|)+V_{s}(|c|)\right) \\
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- Compactify via:

$$
\begin{aligned}
& V_{s}(0):=-s^{-1}, \quad s<0 \\
& V_{s}(0):=+\infty, \quad s \geq 0
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$$

and find all extremal / critical points of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$ !

## Optimal 3-Particle Arrangements on $\mathbb{S}^{1}$

The absolute minimizers (Nerattini-Brauchart-K.):
Theorem
Set $s_{3}:=\ln (4 / 9) / \ln (4 / 3)$.

- Then for $s \neq s_{3}$ the optimal minimal Riesz s-energy $N=3$ arrangement on $\mathbb{S}^{1}$ is unique (up to rotation/permutation).
- For $s<s_{3}$ it is given by the antipodal arrangement $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$.
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The relative minimzers of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$ which are not absolute:

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- For $s_{3}<s<-2$ the antipodal arrangement $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$ (up to permutation) is a relative minimizer which is not absolute.



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- For $s_{3}<s<-2$ the antipodal arrangement $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$ (up to permutation) is a relative minimizer which is not absolute.
- For $-4<s<s_{3}$ the equilateral configuration $(\alpha, \beta, \gamma)=\left(\frac{\pi}{3}, \frac{\pi}{3}, \frac{\pi}{3}\right)$ is a relative minimizer which is not absolute.


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The relative maximizers of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$ which are not absolute:

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For $s<-4$ the equilateral configuration $(\alpha, \beta, \gamma)=\left(\frac{\pi}{3}, \frac{\pi}{3}, \frac{\pi}{3}\right)$ is a relative maximizer which is not absolute. This exhausts all relative maximers of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$ which are not absolute.

## Optimal 3-Particle Arrangements on $\mathbb{S}^{1}$

The absolute maximizers of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$ :

## Theorem

For all $s<0$ the completely degenerate triangular configuration (i.e. the one-point arrangement) given by $(\alpha, \beta, \gamma)=(0,0, \pi)$ (up to permutation) is the unique (up to rotation) absolute maximizer of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$. Contracting and compactifying all energies to $[-1,1]$, the one-pt arrangement is the unique (up to rotation) absolute maximizer for all s.

## Optimal 3-Particle Arrangements on $\mathbb{S}^{1}$

The saddle points of $\left\langle V_{s}\right\rangle(\alpha, \beta, \gamma)$ :

## Theorem

The following list exhausts all the saddle points:

- The equilateral configuration $(\alpha, \beta, \gamma)=\left(\frac{\pi}{3}, \frac{\pi}{3}, \frac{\pi}{3}\right)$ is a saddle point at $s=-4$.
- The antinodal arrancement $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$ (up to permutation) is a saddle for $-2 \leq s<0$. Contracting and compactifying all energies to $[-1,1]$, the antipodal arrangement is a saddle for all $s \geq 0$, too.
- For $\{s<-2\} \cap\{s \neq-4\}$ there are two families (disjoint open sets) of non-universal isosceles triangular equilibrium configurations, and all these are saddle points.


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## Optimal 3-Particle Arrangements on $\mathbb{S}^{1}$

Non-proper pseudo equilibria:

## Theorem

The following are the only non-proper pseudo Riesz s-force equilibria of $N=3$ point particles on $\mathbb{S}^{1}$.

- The completely degenerate triangular configuration given by $(\alpha, \beta, \gamma)=(0,0, \pi)$ (up to permutation) is a non-proper pseudo Riesz s-force equilibrium for all $s \geq-1$
- The antipodal arrangement given by $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$ (up to permutation), is a non-proper pseudo Riesz s-force equilibrium for all $s \geq-1$.


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The proper Riesz s-force equilibria:

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This list exhausts all proper Riesz s-force 3-particle equilibria:

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- The antipodal arrangement given by $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$ (up to permutation), is a proper Riesz s-force equilibrium for all $s<-1$.
- The equilateral configuration $(\alpha, \beta, \gamma)=\left(\frac{\pi}{3}, \frac{\pi}{3}, \frac{\pi}{3}\right)$ is a proper Riesz s-force equilibrium for all $s \in \mathbb{R}$.
- For each $s<-2$ except $s=-4$, there exists a
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This list exhausts all proper Riesz s-force 3-particle equilibria:

- The completely degenerate triangular configuration given by $(\alpha, \beta, \gamma)=(0,0, \pi)$ (up to permutation) is a proper Riesz $s$-force equilibrium for all $s<-1$.
- The antipodal arrangement given by $(\alpha, \beta, \gamma)=\left(\frac{\pi}{2}, \frac{\pi}{2}, 0\right)$ (up to permutation), is a proper Riesz s-force equilibrium for all $s<-1$.
The equilateral configuration $(\alpha, \beta, \gamma)=\left(\frac{\pi}{3}, \frac{\pi}{3}, \frac{\pi}{3}\right)$ is a
proper Riesz s-force equilibrium for all $s \in \mathbb{R}$.
For each $s<-2$ except $s=-4$, there exists a
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## The two isosceles families of non-universal equilibria:

## Theorem

- The family of isosceles triangular Riesz s-force equilibria for $s \in(-\infty,-4)$ interpolates continuously and monotonically between a right triangular configuration ( $\gamma=\pi / 2$ ), to which it converges when $s \downarrow-\infty$, and the equilateral configuration ( $\gamma=\pi / 3$ ), to which it converges when $s \uparrow-4$.
- The family of isosceles triangular Riesz s-force equilibria for $s \in(-4,-2)$ interpolates continuously and monotonically between the equilateral configuration ( $\gamma=\pi / 3$ ), to which it converges when $s \uparrow-4$, and the antipodal arrangement $(\gamma=0)$, to which it converges when $s \uparrow-2$.

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The two isosceles families of non-universal equilibria (cont.d):

## Theorem

The asymptotics of $\gamma$ as function of s for the isosceles triangles is given by the following:
(a) in a left neighborhood of $\gamma=\pi / 2$ (as $s \downarrow-\infty$ ),

$$
\gamma(s) \asymp \frac{\pi}{2}-\sqrt{ } 2^{1+s}
$$

(b) in a neighborhood of $\gamma=\pi / 3$ (for $s \approx-4$ ),

$$
\gamma(s)=\frac{\pi}{3}-\frac{1}{2 \sqrt{3}}(4+s)+\mathcal{O}\left((s+4)^{2}\right)
$$

(c) in a right neighborhood of $\gamma=0$ (as $s \uparrow-2$ ),

$$
\gamma(s) \asymp 0+2^{\frac{1}{2+s}}
$$

## Equilibrium BIFURCATION diagram:



Figure: Bifurcation diagram ( $\gamma$ vs. $s$ ) of the Riesz $s$-force equilibria. Color code: MINIMUM, minimum, saddle, maximum, MAXIMUM.

$\gamma$
Figure: $U_{s}(\gamma): s \in\left\{-30,-15,-10,-7,-5,-4, \frac{\ln (4 / 9)}{\ln (4 / 3)},-2,-1\right\}$.
The graphs are monotonically ordered with $s$, decreasing with $s$

The Figure shows the graph of $\left\langle V_{-4}\right\rangle(\alpha, \beta, \pi-\alpha-\beta)$ over the isosceles triangle $(\alpha, \beta) \in[0, \pi]^{2} \cap\{\alpha+\beta \leq \pi\}$ ( $\leftarrow$ this is the projection of the fundamental triangle into the $(\alpha, \beta)$ plane. Note that this projected illustration somewhat distorts the three-fold symmetry of Hilbert's Monkey Saddle.)


The Figure shows the contour lines of $\left\langle V_{-4}\right\rangle(\alpha, \beta, \gamma)$ in the fundamental triangle in $(\alpha, \beta, \gamma)$ space.


## BIBLIOGRAPHY

- Rachele Nerattini, Johann S. Brauchart, and M.K.: Optimal N-point configurations on the sphere: "Magic" numbers and Smale's 7th problem, J. Stat. Phys. 157:1138-1206 (2014).

Hilbert's Monkey Saddle and other Curiosities Preprint, Rutgers 2017.

- D. Hilbert and S. Cohn-Vossen, "Geometry and the Imagination" (2nd ed.), Chelsea, N.Y. (1952).


## BIBLIOGRAPHY

- Rachele Nerattini, Johann S. Brauchart, and M.K.: Optimal N-point configurations on the sphere: "Magic" numbers and Smale's 7th problem, J. Stat. Phys. 157:1138-1206 (2014).
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