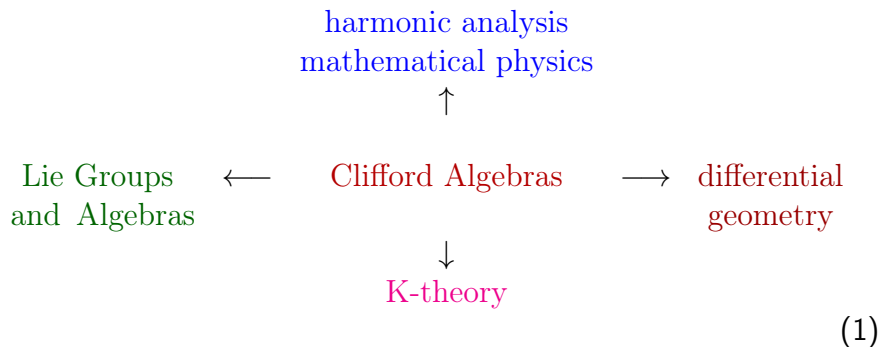


Clifford, Spin and All That

Mariano Echeverria

Clifford Algebras and Mathematics



The Dirac Equation

- In special relativity we have the (famous) relationship

$$E = \sqrt{p_x^2 + p_y^2 + p_z^2 + m^2} \quad (2)$$

- In Quantum Mechanics one “quantizes” the previous equation by turning, E, p_x, p_y, p_z into differential operators

$$E \longrightarrow i \frac{\partial}{\partial t} \quad p_x \longrightarrow -i \frac{\partial}{\partial x} \quad p_y \longrightarrow -i \frac{\partial}{\partial y} \quad p_z \longrightarrow -i \frac{\partial}{\partial z} \quad (3)$$

- Equation 2 becomes

$$i \frac{\partial}{\partial t} = \sqrt{-\frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial z^2} + m^2} \quad (4)$$

Square Root for the Laplacian

- One might try first to find first a square root for the Laplacian
$$\Delta = \sum_{i=1}^3 \frac{\partial^2}{\partial x_i^2}$$
- $\sqrt{\Delta}$ should be a linear differential operator of first order

$$D = \sum_{i=1}^3 c_i \frac{\partial}{\partial x_i} \quad (5)$$

and from the condition $D^2 = -\Delta$ we find that the c_i must satisfy (formal manipulation)

$$\begin{cases} c_1^2 = c_2^2 = c_3^2 = -1 \\ c_1 c_2 + c_2 c_1 = c_1 c_3 + c_3 c_1 = c_2 c_3 + c_3 c_2 = 0 \end{cases} \quad (6)$$

- D is called the Dirac Operator

Pauli Matrices and Clifford Algebras

- Clearly the previous relations can't be satisfied by ordinary numbers (real or complex)
- A representation for the previous relations are the matrices

$$c_1 = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \quad c_2 = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \quad c_3 = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \quad (7)$$

- Since we may associate to $-\Delta = -\sum_{i=1}^3 \frac{\partial^2}{\partial x_i^2}$ the quadratic form $q(\mathbf{x}) = -x_1^2 - x_2^2 - x_3^2$ on \mathbb{R}^3 we may restate our problem as:

Given a quadratic form $q(x)$ defined on a vector space of dimension d is it possible to embed the vector space V into an algebra A in such a way that for elements $e_1, \dots, e_d \in V$ we have $e_i^2 = q(e_i) \cdot 1_A$?

Hamilton's Quaternions

- After many attempts Hamilton found the quaternions $a + ib + jc + kd$ subject to

$$\begin{cases} i^2 = j^2 = k^2 = -1 \\ ij + ji = ik + ki = jk + kj = 0 \end{cases} \quad (8)$$

- To close the algebra, he introduced also the relations

$$ij = k \quad jk = i \quad ki = j \quad (9)$$

- The Clifford Algebra can be considered as a generalization of the quaternions!

Definition Clifford Algebra

- V : finite dimensional vector space over \mathbb{K} (in practice \mathbb{R} or \mathbb{C})
- $\beta : V \times V \rightarrow \mathbb{K}$ bilinear form
- A Clifford algebra $\text{Cliff}(V, \beta)$ is an associative algebra with unit 1 over \mathbb{K} and a linear injective map $\gamma : V \rightarrow \text{Cliff}(V, \beta)$ such that $\{\gamma(x), \gamma(y)\} = 2\beta(x, y)$, $1 \notin \gamma(V)$ and $\gamma(V)$ generates $\text{Cliff}(V, \beta)$ as an algebra
- Given any associative algebra \mathcal{A} with unit element 1 and a linear map $\varphi : V \rightarrow \mathcal{A}$ such that $\{\varphi(x), \varphi(y)\} = 2\beta(x, y)$ there exists an associative algebra homomorphism $\tilde{\varphi} : \text{Cliff}(V, \beta) \rightarrow \mathcal{A}$ such that $\varphi = \tilde{\varphi} \circ \gamma$

$$\begin{array}{ccc} V & \xrightarrow{\gamma} & \text{Cliff}(V, \beta) \\ \varphi \searrow & & \downarrow \tilde{\varphi} \\ & & \mathcal{A} \end{array} \quad (10)$$

Some Properties and Examples

- We start with the tensor algebra $\mathcal{T}(V)$ and take the two-sided ideal $\mathcal{J}(V, \beta)$ generated by $x \otimes y - y \otimes x - 2\beta(x, y)$ for $x, y \in V$. We define $\text{Cliff}(V, \beta) = \mathcal{T}(V)/\mathcal{J}(V, \beta)$
- We might identify V with $\gamma(V)$ and \mathbb{K} can be identified with $\text{span}(1)$.
- If $\dim V = n$ then as a vector space $\dim \text{Cliff}(V, \beta) = 2^n$
- Over \mathbb{R} define $q_{n,m}(x) = \sum_{i=1}^n x_i^2 - \sum_{i=n+1}^{m+n} x_i^2$. Then

$$\begin{aligned} \text{Cl}(0, 1) &\simeq \mathbb{C} \\ \text{Cl}(0, 2) &\simeq \mathbb{H} \end{aligned} \tag{11}$$

- **Remark:** $\text{Cl}(0, 3)$ are not the octonions since they are not even associative.

Superalgebra

- If we define $\Pi : V \rightarrow V$ by $\Pi(v) = -v$ then by the universal property Π extends to an automorphism $\Pi : \text{Cliff}(V, \beta) \rightarrow \text{Cliff}(V, \beta)$
- We define the even part $\text{Cliff}^+ = \{v \in \text{Cliff} : \Pi(v) = v\}$ and the odd part $\text{Cliff}^- = \{v \in \text{Cliff} : \Pi(v) = -v\}$. We have that

$$\begin{aligned}\text{Cliff}^+ \text{Cliff}^+ &\subseteq \text{Cliff}^+ \\ \text{Cliff}^- \text{Cliff}^- &\subseteq \text{Cliff}^+ \\ \text{Cliff}^+ \text{Cliff}^- &\subseteq \text{Cliff}^- \\ \text{Cliff}^- \text{Cliff}^+ &\subseteq \text{Cliff}^-\end{aligned}\tag{12}$$

- Moreover,

$$\text{Cliff} = \text{Cliff}^+ \oplus \text{Cliff}^-\tag{13}$$

This turns the Clifford Algebra into a **Superalgebra**.

Clifford Algebras as a Quantization Procedure

- If we take $q(x) = 0$ then the Clifford algebra is the exterior algebra $\wedge V$
- In this way, if we take a non-degenerate symmetric bilinear form $b(\cdot, \cdot)$ by a parameter t gives a Clifford Algebra $\text{Cliff}(V, tb(\cdot, \cdot))$ which can be considered as a deformation of the exterior algebra $\wedge V$
- As a superalgebra, the exterior algebra $\wedge V$ is supercommutative while the Clifford Algebra is not. Therefore, the Clifford Algebra is a noncommutative version of the exterior algebra in this other sense

Bott Periodicity

- For $n \geq 0$ we have the following isomorphism of complex associative algebras

$$\text{Cliff}(\mathbb{C}^{n+2}) \simeq \text{Cliff}(\mathbb{C}^n) \otimes \text{Mat}_{2 \times 2}(\mathbb{C}) \quad (14)$$

- To see this write $\mathbb{C}^{n+2} = \mathbb{C}e_1 \oplus \mathbb{C}e_2 \oplus \mathbb{C}^n$ where e_1, \dots, e_{n+2} generate $\text{Cliff}(\mathbb{C}^{n+2})$ and e_1^*, \dots, e_n^* generate $\text{Cliff}(\mathbb{C}^n)$

$$\phi : \mathbb{C}^{n+2} \longrightarrow \text{Cliff}(\mathbb{C}^n) \otimes \text{Mat}_{2 \times 2}(\mathbb{C})$$

$$\phi(e_1) = 1 \otimes \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \quad \phi(e_2) = 1 \otimes \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \quad (15)$$

$$\phi(e_j) = (ie_{j-2}^*) \otimes \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \quad 3 \leq j \leq n+2$$

Classification of Clifford Algebras over \mathbb{C}

- From Bott periodicity and the fact that

$$\begin{cases} \text{Cliff}(\mathbb{C}^0) = \text{Cliff}(1) \simeq \mathbb{C} \\ \text{Cliff}(\mathbb{C}) \simeq \mathbb{C} \oplus \mathbb{C} \end{cases} \quad (16)$$

it follows that for $k \geq 0$

$$\begin{cases} \text{Cliff}(\mathbb{C}^{2k}) \simeq \mathbb{C} \otimes_{i=1}^k \text{Mat}_{2 \times 2}(\mathbb{C}) \simeq \text{End}(\mathbb{C}^{2^k}) \\ \text{Cliff}(\mathbb{C}^{2k+1}) \simeq (\mathbb{C} \oplus \mathbb{C}) \otimes_{i=1}^k \text{Mat}_{2 \times 2}(\mathbb{C}) \simeq \text{End}(\mathbb{C}^{2^k}) \oplus \text{End}(\mathbb{C}^{2^k}) \end{cases} \quad (17)$$

- For $n = 2k, 2k + 1$

$$\Delta_n = \mathbb{C}^{2^k} \quad (18)$$

is called the **spinor space**

The Spin Group

- Over \mathbb{R}^n consider $q(x) = -x_1^2 - \dots - x_n^2$. Observe that $x^2 = -\|x\|^2$.
- The **Pin Group** $\text{Pin}(n)$ consists of the group generated under the Clifford Multiplication by all vectors $x \in S^{n-1}$
- The **Spin Group** is

$$\text{Spin}(n) = \text{Pin}(n) \cap \text{Cliff}^+ \quad (19)$$

- It can be shown that

$$\begin{aligned} \lambda : \text{Pin}(n) &\longrightarrow O(n) \\ \lambda(x)y &= xyx^T \end{aligned} \quad (20)$$

is a group homomorphism which is a double cover of $O(n)$

- For example,

$$\begin{aligned} \text{Spin}(2) &= SO(2) \\ \text{Spin}(3) &= SU(2) \end{aligned} \quad (21)$$

Thank you!

