

Last Time For our purposes, a Spin^c -structure is a nowhere-vanishing vector field v on Y .

• $v_1 \sim v_2$ if they are homotopic after deleting a ball \emptyset from Y .

• $\{v \text{ nonvanishing vector fields}\} \longleftrightarrow \{ \text{maps } F_v : Y \rightarrow S^2 \}$

$$\longleftrightarrow H^2(Y; \mathbb{Z})$$

after choosing a trivialization

• $v \mapsto \delta^*(v) = f_v^*(\alpha)$ a generator of $H^2(S^2)$.

• $\delta(v_1, v_2) = \delta^*(v_1) - \delta^*(v_2) \in H^2(Y; \mathbb{Z})$ independent of trivialization

• Action of $H^2(Y)$ on $\text{Spin}^c(Y)$ via letting $a \cdot v$ be st $\delta(a \cdot v, v) = a$

• $\text{Spin}^c(Y)$ comes w/ an involution $s \mapsto \bar{s}$ via $v \mapsto -v$

• We also have a first Chern class

$$c_1 : \text{Spin}^c(Y) \rightarrow H^2(Y; \mathbb{Z})$$

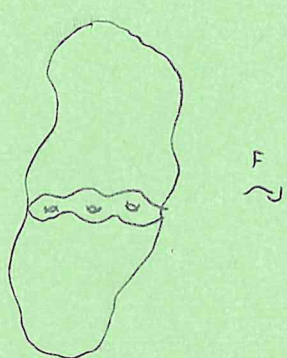
$$s \mapsto s - \bar{s}$$

Back to Heegaard diagrams.

(2)

We have a map $s_z: \pi_{\alpha}^{-1} \cap \pi_{\beta}^{-1} \rightarrow \text{Spin}^c(Y)$ which refines ϵ .

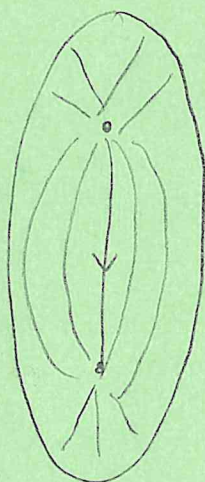
Let



Picking \vec{x} and \vec{y} picks a set of g trajectories from index one to index 2 critical pts, we also have a trajectory from index 0 to index 3. Delete nibbles and get a complement of disjoint unions on which ∇F does not vanish.

We can fill this in to get a nowhere-vanishing vector field:

Indeed $\frac{\text{Thm}}{s_z(y) - s_z(x) = PP[\bar{\epsilon}(\vec{x}, \vec{y})]}$ i.e. $s_z(y) = s_z(x) \Leftrightarrow \pi_2(\vec{x}, \vec{y})$ nonempty.



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In general one can do something of this ilk for any case where the indices of the critical pts differ.

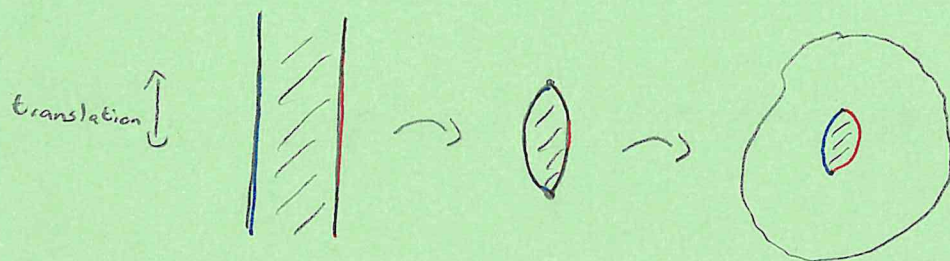
Boundary operator

(3)

• We consider disks for which $n(\phi) = 1$.

• For a given homotopy class of disks $\phi \in \pi_2(x, \gamma)$, we let $M(\phi)$ be the moduli space of (pseudo)-holic representations of ϕ .

• This comes to us w/ an action of \mathbb{R}



• We quotient out to get rigid pseudoholic disks

$$\hat{M}(\phi) = \frac{M(\phi)}{\mathbb{R}}$$

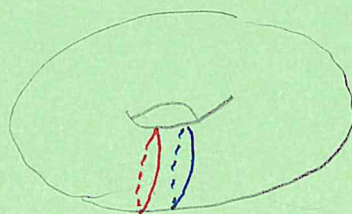
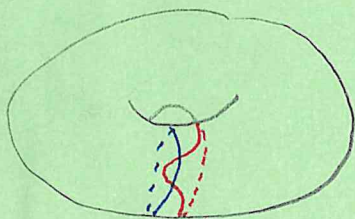
Add one more condition:

We ask for Heegaard diagrams which are weakly admissible
i.e. any disk whose ^{shadow's} boundary is $\partial D(\phi) = \sum n_i \alpha_i + m_i \beta_i$ contains domains of both positive and negative multiplicities.

Yes boundary

No

A periodic domain



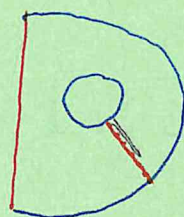
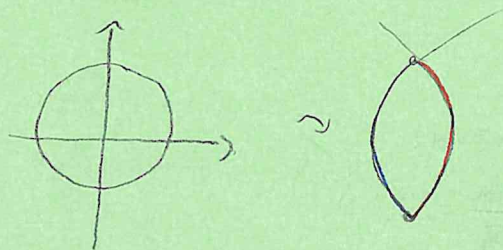
Propn There is a family of perturbations of the complex structure with the property that if $n(\varphi) = 1$ then $\hat{M}(\varphi)$ is a compact oriented zero-dimensional mfd.

Helpful Note If we have $u \in M(\varphi)$, there is a map

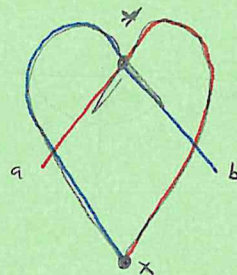
$$\begin{array}{ccc} g\text{-Fold} & \rightarrow & \\ \text{branched} & & \\ \text{cover} & & \end{array} \quad \begin{array}{ccc} S & \xrightarrow{\hat{u}} & \Sigma \\ \downarrow P & & \downarrow P \\ D^2 & \xrightarrow{u} & \text{Sym}^g(\Sigma) \end{array}$$

We get a hol'ic map
for $\hat{u}: S \rightarrow \Sigma$

Examples

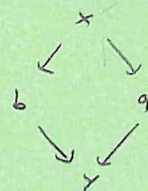


$$n(\varphi) = 1$$



$$n(\varphi) = 2$$

There is a one-dim'l family of maps by the Riemann mapping thm; we quotient and get a 0-dim'l family.



\exists a biholomorphism between any two simply-ctd domains in \mathbb{C} or \mathbb{R}^n ; in fact given that the space of maps

$\begin{pmatrix} a & b \\ c & d \end{pmatrix}$ is $3d$, fixing $ad-bc=1$ two points gives a $1d$ set



The boundary operator is

$$\partial x = \sum_{\gamma \in \pi_\alpha \cap \pi_\beta} \left(\sum_{\substack{\phi \in \pi_2(x, \gamma) \\ u(\phi) = 1}} \# \hat{M}(\phi) \cdot U^{n_2(\phi)} \gamma \right)$$

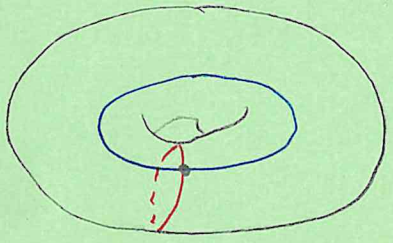
$$\hat{\partial} x = \sum_{\gamma \in \pi_\alpha \cap \pi_\beta} \left(\sum_{\substack{\phi \in \pi_2(x, \gamma) \\ u(\phi) = 1 \\ n_2(\phi) = 0}} \# \hat{M}(\phi) \gamma \right) \quad \left. \vphantom{\sum_{\gamma \in \pi_\alpha \cap \pi_\beta}} \right\} \text{This is equivalent to counting disks in } \text{Sym}^0(\Sigma \setminus \{x, y\})$$

In the case $b_i = 0$,

This comes with a grading: $gr(x, y) = u(\phi) - 2n_2(\phi)$ where ϕ is any disk connecting them. Ergo the differential lowers grading by 1.

More generally $\mathbb{Z}/2\mathbb{Z}$ -graded

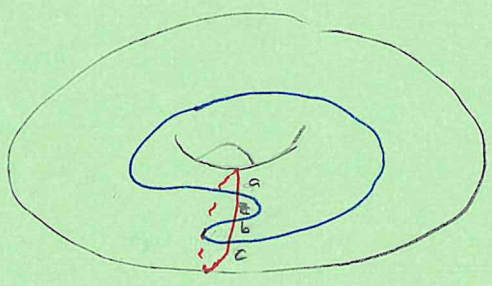
Examples



$$CF^-(S^3) = \mathbb{F}[U] \langle x \rangle$$

$$HF^-(S^3) = \mathbb{F}[U]$$

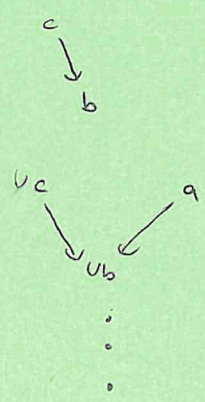
$$\hat{HF}(S^3) = \mathbb{F}$$

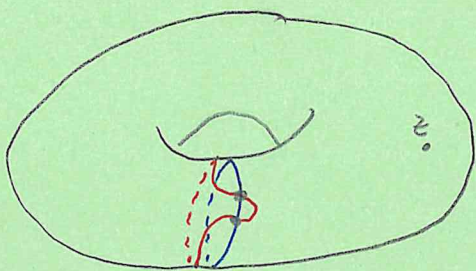


$$CF^-(S^3) = \mathbb{F}[U] \langle a, b, c \rangle$$

$$HF^-(S^3) = \mathbb{F}[U] \langle [Uc + a] \rangle$$

$$\hat{HF}(S^3) = \mathbb{F} \langle [a] \rangle$$



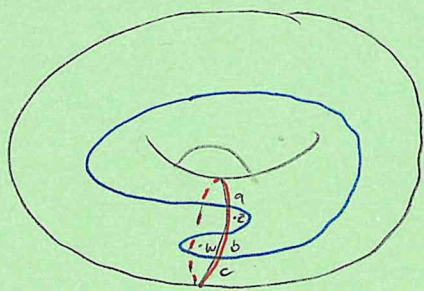


$$\bullet \text{HFK}^-(S^3) = \mathbb{F}_2[u] \langle [a], [b] \rangle$$

• Every other spin^c structure is empty.

IF we have a diagram for (Y, K) , we have a filtration coming from the second basepoint, via counting intersections w/ v_w .

$$\begin{aligned} \partial_K \vec{x} &= \sum_{\gamma} \sum_{\substack{\phi \in \pi_2(\gamma, \gamma) \\ n(\phi)=1, \\ n_w(\phi)=0}} U^{n_z(\phi)} \bar{x}^{n(\phi)} \vec{\gamma} & \hat{\partial}_K \vec{x} &= \sum_{\gamma} \sum_{\substack{\phi \in \pi_2(\vec{x}, \vec{\gamma}) \\ n(\phi)=1, \\ n_w(\phi)=n_z(\phi)=1}} \# \hat{M}(\phi) \vec{\gamma} \end{aligned}$$



$$\text{KFK}^-(3,1) = \mathbb{F}[u] \langle a, b, c \rangle$$

$$\text{HFK}^-(3,1) = \mathbb{F}[u] \langle [a] \rangle \oplus \mathbb{F} \langle [b] \rangle$$

$$\hat{\text{HFK}}(3,1) = \mathbb{F} \langle [a], [b], [c] \rangle$$

Note

c a

b

v_u

v_b

a

$$0 \rightarrow \text{CFK}^-(Y, K) \xrightarrow{u} \text{CFK}^-(Y, K) \rightarrow \hat{\text{CFK}}(Y, K) \rightarrow 0$$

So you expected

$$\hat{\text{HFK}}(3,1) = \ker(u) \oplus \text{coker}(u)$$

What is a spin^c -structure?

$$\begin{aligned} \bullet \text{ } SO(3) &= SU(2) / \{\pm 1\} = U(2) / U(1) \\ &\parallel \\ &\mathbb{R}P^3 \end{aligned}$$

$$\begin{aligned} \bullet \text{ } U(2) \text{ principal circle bundle } &\Leftrightarrow [x, SU(2)] = H^2(x) \\ &\downarrow \\ &SO(3) \\ &\underbrace{\hspace{2cm}} \\ &\text{This bundle} \\ &\text{corresponds} \\ &\text{to the nontrivial element} \\ &\text{in } H^2(SO(3); \mathbb{Z}) = \mathbb{Z}_2 \end{aligned}$$

$$\begin{aligned} 1 \rightarrow \mathbb{Z}_2 \rightarrow \text{Spin}(3) \rightarrow SO(3) \rightarrow 1 \\ \parallel \\ SU(2) = S^3 \end{aligned}$$

$$\begin{aligned} 1 \rightarrow \mathbb{Z}_2 \rightarrow \text{Spin}^c(3) \rightarrow SO(3) \times U(1) \rightarrow 1 \\ \parallel \\ (SU(2) \times U(1)) / \{\pm 1\} \\ \parallel \\ U(2) \end{aligned}$$

$$\bullet \text{ } \text{Spin}(3) = SU(2)$$

$$\text{Spin}^c(3) = (U(1) \times \text{Spin}(3)) / \{\pm 1\} = (U(1) \times SU(2)) / \{\pm 1\} = U(2)$$

\bullet Y closed oriented 3-mfd, has a metric

$$\begin{array}{ccc} TY & \xrightarrow{\sim} & F_\pi \text{ principal } SO(3) \text{ bundle.} \\ \downarrow & & \downarrow \\ Y & & Y \end{array}$$

\bullet A spin^c -structure is a lift of F_π to a principal $\text{Spin}^c(3) = U(2)$ bundle

$$\begin{array}{ccc} \bullet \text{ So a } U(2)\text{-bundle } F & \text{w/ an isomorphism } & F/U(1) \cong F_\pi \\ \downarrow & & \downarrow \\ Y & & Y \end{array}$$

\leadsto A Spin^c -structure is an element of $H^2(F_\pi)$ whose reduction to every fiber is the nonzero element of $H^2(F_\pi)$. $\text{Spin}^c(Y)$

via the element α associated to F

$$\downarrow$$

$$F/U(1) \cong F_r$$

$H^2(Y)$ acts on $H^2(F_r)$ via pullback $H^2(Y) \rightarrow H^2(F_r)$. This is free and transitive [$F_r = Y \times SO(3)$ after picking a trivialization of TY] so $H^2(F_r) = H^2(Y) \oplus (\mathbb{Z}/2\mathbb{Z})$.

We also have the description given previously, which we can rewrite as an element of $H^2(SY)$ whose restriction to every fiber is the oriented generator of $H^2(S^2)$.

$p: SO(3) \rightarrow S^2$ is a circle bundle w/ fiber corresponding to the nonzero element of $H_1(SO(3)) = \mathbb{Z}/2\mathbb{Z}$

$(e_1, e_2, e_3) \mapsto e_1$

Pullback $p^*: H^2(S^2) \rightarrow H^2(SO(3))$ sends a generator of $H^2(S^2)$ to a generator of $H^2(SO(3))$

Given

$$\begin{array}{ccc} (e_1, e_2, e_3) & \mapsto & e_1 \\ F_r & \longrightarrow & SY \\ \downarrow & & \downarrow \\ Y & \longrightarrow & Y \end{array}$$

$p^*: H^2(SY) \rightarrow H^2(F_r)$ sends a nonzero vector field to a Spin^c -structure.

Note - $TY = u \oplus u^\perp \leadsto$ Structure group $(1) \oplus U(1) \subseteq U(2)$ } of course has a Spin^c -structure

$\circ c_1(s) = c_1(u^\perp)$ is $\sigma^*(v) - \sigma^*(\cdot, \cdot) = s - \bar{s}$