

Chapter 1

Superconnection

1.1 \mathbb{Z}_2 -graded algebra

Definition 1.1.1. A is a \mathbb{C} -algebra, A is \mathbb{Z}_2 -graded, if there is a splitting: $A = A_+ \oplus A_-$, such that $A_+A_+ \subset A_+, A_-A_- \subset A_+, A_+A_- \subset A_-, A_-A_+ \subset A_-$.

$A_{\text{even}} = A_+, A_{\text{odd}} = A_-$ are called the even and odd part of A . The usual algebra is trivially \mathbb{Z}_2 -graded with $A = A_+$.

Definition 1.1.2. for homogeneous elements $a, b \in A_{\pm}$, the supercommutator is

$$[a, b] = ab - (-1)^{\deg a \cdot \deg b} ba, \quad \deg a = \begin{cases} 0, & a \in A_+ \\ 1, & a \in A_- \end{cases}$$

we extend linearly to A . $a, b \in A$ is said to be supercommutative if $[a, b] = 0$. Note that $[A_+, A_+] \subset A_+, [A_-, A_-] \subset A_+, [A_+, A_-] \subset A_-, [A_-, A_+] \subset A_-$.

Proposition 1.1.1 (Generalized Jacobi Identity).

$$[a, [b, c]] + (-1)^{\deg a(\deg b + \deg c)} [b, [c, a]] + (-1)^{\deg c(\deg a + \deg b)} [c, [a, b]] = 0, \text{ or equivalently}$$

$$[a, [b, c]] = [[a, b], c] + (-1)^{\deg a \cdot \deg b} [b, [a, c]] \text{ (derivation property)}$$

Example 1.1.1. 1. E is a finite dimensional vector space. The exterior algebra: $\wedge^* E^* = \wedge^{\text{even}} E^* \oplus \wedge^{\text{odd}} E^*$ is a commutative \mathbb{Z}_2 -graded algebra, since $\omega \wedge \eta = (-1)^{\deg \omega \cdot \deg \eta} \eta \wedge \omega$.

2. $\tau : E \rightarrow E$ is an endomorphism, s.t. $\tau^2 = id$, then E is \mathbb{Z}_2 -graded: $E = E_+ \oplus E_-$, $\tau|_{E_{\pm}} = \pm 1$. $End(E)$ is a \mathbb{Z}_2 -graded algebra:

$$End(E)^{\text{even}} = \{A \in End(E) : A\tau = \tau A\}, End(E)^{\text{odd}} = \{A \in End(E) : A\tau = -\tau A\}$$

3. q is a nondegenerate quadratic form on E , the Clifford algebra is \mathbb{Z}_2 -graded: $\alpha : Cl(E, q) \rightarrow Cl(E, q)$ is the involution defined by $\alpha(v) = -v, \forall v \in E$.

Definition 1.1.3. A, B are \mathbb{Z}_2 -graded algebra, $A \otimes B$ has a product:

$$a \otimes b \cdot a' \otimes b' = (-1)^{\deg b \cdot \deg a'} aa' \otimes bb'$$

for homogeneous elements $a, a' \in A, b, b' \in B$. It's well defined and $A \otimes B$ becomes a \mathbb{Z}_2 -graded algebra, denoted by $A \hat{\otimes} B$:

$$(A \hat{\otimes} B)_+ = (A_+ \otimes B_+) \oplus (A_- \otimes B_-), (A \hat{\otimes} B)_- = (A_+ \otimes B_-) \oplus (A_- \otimes B_+)$$

Remark 1. If A, B are unitary ($1_A \in A, 1_B \in B$), then A, B are embedded in $A \hat{\otimes} B$. We can identify a with $a \otimes 1_B$, and b with $1_A \otimes b$, then $a \cdot b = (-1)^{\deg a \cdot \deg b} \cdot a$, so A and B are supercommutative in $A \hat{\otimes} B$.

Example 1.1.2. $(E_1, q_1), (E_2, q_2)$ are two vector spaces with nondegenerate quadratic forms. Let $E = E_1 \oplus E_2, q = q_1 \oplus q_2 : q((v_1, v_2)) = q_1(v_1) + q_2(v_2)$, then $Cl(E, q) = Cl(E_1, q_1) \hat{\otimes} Cl(E_2, q_2)$.

Definition 1.1.4. assume $E = E_+ \oplus E_-$ is \mathbb{Z}_2 -graded defined by involution τ . $\forall A \in End(E)$, the supertrace of A is $Tr_s(A) = Tr(\tau A)$. In blocked matrix, $\tau = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, if $A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$, then $Tr_s(A) = Tr(a) - Tr(d)$.

Proposition 1.1.2. $\forall A, B \in End(E), Tr_s[A, B] = 0$.

assume $(E, \tau), (F, \sigma)$ are two \mathbb{Z}_2 -graded vector space, then $E \hat{\otimes} F = (E \otimes F, \tau \otimes \sigma)$ is a \mathbb{Z}_2 -graded vector space. $\forall A \in End(E), B \in End(F)$, we have

$$Tr_s(A \otimes B) = Tr(\tau A \otimes \sigma B) = Tr(\tau A) Tr(\sigma B) = Tr_s(A) Tr_s(B).$$

Definition 1.1.5. E is a finite dimensional, we define the supertrace on the \mathbb{Z}_2 -graded algebra $\wedge^* E^* \hat{\otimes} End(E)$ (matrix of forms):

$$\begin{aligned} Tr_s : \wedge^* E^* \hat{\otimes} End(E) &\longrightarrow \wedge^* E^* \\ \omega \otimes A &\longmapsto \omega Tr_s(A) \end{aligned}$$

By the universal property of tensor product, this is well defined. For simplicity, we write ωA for $\omega \otimes A$. We have

$$[\omega A, \omega' A] = (-1)^{\deg A \cdot \deg \omega} \omega \omega' [A, A']$$

1.2 Superconnection

Definition 1.2.1. E is a finite dimensional \mathbb{Z}_2 -graded complex vector bundle on the manifold M . A superconnection A is an odd differential operator acting on $\Gamma(\wedge^* T^* M \hat{\otimes} E)$, which satisfies the Leibniz rule:

$$A(\omega s) = d\omega \cdot s + (-1)^{\deg \omega} \omega \wedge A s$$

Remark 2. $\wedge^* T^* M \hat{\otimes} E$ is a \mathbb{Z}_2 -graded vector bundle, $\Gamma(\wedge^* T^* M \hat{\otimes} E)$ is the set of E -valued forms which is naturally \mathbb{Z}_2 -graded. Locally on an open set U , we choose a cotangent frame $\{e^i\}$ and a frame of $\{s_\alpha\} = \{s_\lambda^+\} \cup \{s_\mu^-\}$ of $E = E_+ \oplus E_-$, then locally $\tau \in \Gamma(\wedge^* T^* M \hat{\otimes} E)$ is:

$$\tau = f_{i_1, \dots, i_r}^\alpha e^{i_1} \wedge \dots \wedge e^{i_r} \otimes s_\alpha, \tau_{i_1, \dots, i_r}^\alpha \in C^\infty(U)$$

$A\tau = g_{j_1, \dots, j_s}^\beta e^{j_1} \wedge \dots \wedge e^{j_s} \otimes s_\beta, g = D(f), D$ is a matrix of differential operators. τ is even iff locally it's the sum of $\omega^{even} \cdot s_\lambda^+$ and $\omega^{odd} \cdot s_\mu^-$.

$$\Gamma = \Gamma_{even} \oplus \Gamma_{odd} = \Gamma((\wedge^* T^* M \hat{\otimes} E)_{even}) \oplus \Gamma((\wedge^* T^* M \hat{\otimes} E)_{odd})$$

By definition, $A(\Gamma_{even}) \subset \Gamma_{odd}, A(\Gamma_{odd}) \subset \Gamma_{even}$.

Example 1.2.1. ∇^E is a connection preserving the splitting $E = E_+ \oplus E_-$, i.e. $\nabla^E(\Gamma(E_\pm)) \subset \Gamma(E_\pm)$.

$$\text{Write } \nabla^E = \nabla^{E_+} \oplus \nabla^{E_-}, \text{ or in blocked matrix: } \nabla^E = \begin{pmatrix} \nabla^{E_+} & 0 \\ 0 & \nabla^{E_-} \end{pmatrix}$$

We extend ∇^E to act on $\Gamma(\wedge^* T^* M \hat{\otimes} E)$ satisfying the Leibniz rule: $\nabla^E(\omega s) = d\omega \cdot s + (-1)^{\deg \omega} \omega \wedge \nabla^E s$, Then ∇^E is a superconnection, since it increases the degree of forms by 1, while keeps the degree of sections of the bundle E , thus odd.

let $D \in \Gamma(\text{End}(E)^{\text{odd}})$, i.e. $D_x \in \text{End}(E_x)^{\text{odd}}$, in matrix $D = \begin{pmatrix} 0 & D_- \\ D_+ & 0 \end{pmatrix}$, then $\nabla^E + D$ is a superconnection, since D preserves the degree of forms, and interchanges the even and odd part of E .

In particular, when $M = \{pt\}$, E is a \mathbb{Z}_2 -graded vector space, $D \in \text{End}(E)^{\text{odd}}$ is a superconnection.

Proposition 1.2.1. *Any superconnection A can be written as $A = \nabla^E + B$, with $\nabla^E = \nabla^{E_+} \oplus \nabla^{E_-}$, and $B \in \Gamma(\mathcal{B})$, where $\mathcal{B} = \wedge^* T^* M \hat{\otimes} \text{End}(E)$ is a bundle of algebras.*

Proof. Let $B = A - \nabla^E$, then B is an odd operator on $\Gamma(\wedge^* T^* M \hat{\otimes} E)$. $\forall f \in C^\infty(M)$, $s \in \Gamma(\wedge^* T^* M \hat{\otimes} E)$,

$$B(fs) = A(fs) - \nabla^E(fs) = df \cdot s + fAs - (df \cdot s + f\nabla^E s) = f(A - \nabla^E)s = fBs$$

so B is $C^\infty(M)$ linear, i.e. tensorial.

Let $Bs_\alpha = \theta_\alpha^\beta s_\beta$, $\theta_\alpha^\beta \in \Gamma(\wedge^* T^* U)$, define $\tilde{B} = \theta_\alpha^\beta s_\alpha^* \otimes s_\beta \in \Gamma(\wedge^* T^* M \hat{\otimes} \text{End}(E))$, where $\{s_\alpha^*\}$ is the dual basis of $\{s_\beta\}$, then \tilde{B} is independent of the chosen frame since B is tensorial.

conversely $\forall \tilde{B} \in \Gamma(\wedge^* T^* M \hat{\otimes} \text{End}(E))$, \tilde{B} acts naturally on $\Gamma(\wedge^* T^* M \hat{\otimes} E) : (\omega \otimes D) \cdot (\eta \otimes s) = (-1)^{\text{deg } D \cdot \text{deg } \eta} \omega \wedge \eta \otimes Ds$, this action is clearly C^∞ linear. So we can identify B and \tilde{B} , and the proposition is proved. \square

Remark 3. if we choose local frame $\{s_\lambda^+\} \cup \{s_\mu^-\}$ of E on U , then $E \cong \mathbb{C}^n \oplus \mathbb{C}^m$. The de Rham operator d is a connection on trivial bundles. By the proposition, every connection is locally of the form:

$$A = d + B, B \in \Gamma(\wedge^* T^* M \hat{\otimes} \text{End}(\mathbb{C}^n \oplus \mathbb{C}^m))_{\text{odd}}$$

1.3 Chern-Weil theory for superconnections

Proposition 1.3.1. *A is a superconnection, then $A^2 \in \Gamma(\mathcal{B}_{\text{even}})$.*

recall that $\mathcal{B} = \wedge^* T^* M \hat{\otimes} \text{End}(E)$.

Proof. Since A is odd, A^2 is an even operator. $\forall f \in C^\infty(M)$,

$$A^2(fs) = A(df \cdot s + f \cdot As) = d(df)s - df \wedge As + df \wedge As + fA^2s = fA^2s$$

so A^2 is a tensor. \square

Definition 1.3.1. The curvature of a superconnection A is A^2 . The Chern character of A is $ch(A) = \varphi Tr_s(e^{-A^2}) \in \Gamma(\wedge^* T^* M)$, where

$$\begin{aligned} \varphi : \wedge_{\mathbb{C}}^* T^* M &\longrightarrow \wedge_{\mathbb{C}}^* T^* M = \wedge^* T^* M \otimes \mathbb{C} \\ \alpha &\mapsto (2\pi i)^{-\frac{\text{deg } \alpha}{2}} \alpha \end{aligned}$$

is the normalization endomorphism, which makes the definition agree with the Chern character of a complex vector bundle.

Theorem 1.3.1 (Quillen). *$ch(A)$ is a closed, even form. $[ch(A)] = ch(E_+) - ch(E_-)$, $[ch(A)] \in H_{dR}^{\text{even}}(M)$ is the cohomological class represented by $ch(A)$.*

Proof. Take a local trivialization $E = E_+ \oplus E_- \cong \mathbb{C}^n \oplus \mathbb{C}^m$, then $A = d + B$, where $B \in \Gamma(\wedge^* T^* M \hat{\otimes} \text{End}(\mathbb{C}^n \oplus \mathbb{C}^m))_{\text{odd}}$

$$\begin{aligned} dTr_s(\exp(-A^2)) &= Tr_s(d \exp(-A^2)) = Tr_s[d, \exp(-A^2)] \\ &= Tr_s[d + B, \exp(-A^2)] = Tr_s[A, \exp(-A^2)] = 0 \end{aligned}$$

In the third equality, we use $Tr_s[B, \exp(-A^2)] = 0$: since $\exp(-A^2)$ is a tensor, the supertrace of supercommutator vanishes. The last equality follows from the Bianchi identity: $[A, A^2] = 0$, which is trivial in this setting.

Remark 4. It's easy to see that, $\forall B \in \Gamma(\wedge^* T^* M \hat{\otimes} \text{End}(E))$, $[A, B] \in \Gamma(\wedge^* T^* M \hat{\otimes} \text{End}(E))$.

Since A^2 is even, $\exp(-A^2)$ is even. It is the sum of the form

$$\text{even form} \otimes \begin{pmatrix} * & 0 \\ 0 & * \end{pmatrix}, \text{ odd form} \otimes \begin{pmatrix} 0 & * \\ * & 0 \end{pmatrix}$$

but $Tr_s(\text{odd form} \otimes \begin{pmatrix} 0 & * \\ * & 0 \end{pmatrix}) = 0$, so $Tr_s(\exp(-A^2))$ contains only even forms.

Next we want to prove $[ch(A)]$ is independent of the superconnection. Let $\mathcal{A} = \{\text{superconnection}\}$, \mathcal{A} is an affine space. Given $A_0, A_1 \in \mathcal{A}$, let $A_t = (1-t)A_0 + tA_1$. Define $\mathbb{A} = dt \frac{\partial}{\partial t} + A_t$, then \mathbb{A} is a superconnection on the vector bundle $\pi_M^* E \rightarrow M \times [0, 1]$, where $\pi_M : M \times [0, 1] \rightarrow M$ is the projection. $\mathbb{A}^2 = A_t^2 + dt \frac{\partial A_t}{\partial t}$,

$$ch(\mathbb{A}) = \varphi Tr_s \exp(-(A_t^2 + dt \frac{\partial A_t}{\partial t})) = ch(A_t) + dt \beta_t$$

β_t is an odd form on $M \times [0, 1]$. We know that $ch(\mathbb{A})$ is closed, so

$$0 = d ch(\mathbb{A}) = d_M ch(A_t) + dt \frac{\partial}{\partial t} ch(A_t) - dt d_M \beta_t = dt \frac{\partial}{\partial t} ch(A_t) - dt d_M \beta_t$$

so $\frac{\partial}{\partial t} ch(A_t) = d \beta_t$ and

$$ch(A_1) - ch(A_0) = \int_0^1 \frac{\partial}{\partial t} ch(A_t) dt = d \int_0^1 \beta_t dt$$

so $[ch(A_0)] = [ch(A_1)]$.

Thus we can choose any superconnection to compute $[ch(A)]$. We let $A = \nabla^E = \begin{pmatrix} \nabla^{E_+} & 0 \\ 0 & \nabla^{E_-} \end{pmatrix}$, then

$$\begin{aligned} ch(A) &= [\varphi Tr_s \exp(-\nabla^{E,+2})] = [\varphi Tr \exp(-\nabla^{E_+,2})] - [\varphi Tr \exp(-\nabla^{E_-,2})] \\ &= [Tr \exp(\frac{\sqrt{-1}}{2\pi} \nabla^{E_+,2})] - [Tr \exp(\frac{\sqrt{-1}}{2\pi} \nabla^{E_-,2})] = ch(E_+) - ch(E_-) \end{aligned}$$

□

Example 1.3.1. $M = \mathbb{C}$, $E = E_+ \oplus E_- = \mathbb{C} \oplus \mathbb{C}$, $D_z = \begin{pmatrix} 0 & \bar{z} \\ z & 0 \end{pmatrix}$, $\nabla^E = d$. $A_t = \nabla^E + \sqrt{t}D = d + \sqrt{t}D$,

$$A_t^2 = d^2 + \sqrt{t}dD + tD^2 = \begin{pmatrix} t|z|^2 & 0 \\ 0 & t|z|^2 \end{pmatrix} + \sqrt{t} \begin{pmatrix} 0 & d\bar{z} \\ dz & 0 \end{pmatrix}$$

Since

$$\frac{1}{n!} Tr_s (t|z|^2 + \sqrt{t} \begin{pmatrix} 0 & d\bar{z} \\ dz & 0 \end{pmatrix})^n = \frac{1}{n!} \frac{n \cdot (n-1)}{2} (t|z|^2)^{n-2} t Tr_s \begin{pmatrix} d\bar{z}dz & 0 \\ 0 & dzd\bar{z} \end{pmatrix} = -t \frac{(t|z|^2)^{n-2}}{(n-2)!} dzd\bar{z}$$

$$Tr_s \exp(-A_t^2) = \sum_{n=0}^{\infty} (-1)^n Tr_s (A_t^2)^n = -t \exp(-t|z|^2) dzd\bar{z}$$

$$ch(A_t) = \varphi Tr_s \exp(-A_t^2) = \frac{1}{2\pi i} \cdot (-t) \exp(-t|z|^2) dzd\bar{z} = \frac{t}{\pi} \exp(-t|z|^2) dx dy = P_{\frac{1}{4t}}(z) dx dy$$

where $p_t(x, y) = \frac{1}{4\pi t} \exp(-\frac{|x-y|^2}{4t})$ is the heat kernel of Laplace $\Delta = -(\frac{d^2}{dx^2} + \frac{d^2}{dy^2}) = -4 \frac{\partial^2}{\partial z \partial \bar{z}}$ on \mathbb{C} . $\alpha_t = ch(A_t)$ has the following properties:

- $\int_{\mathbb{C}} \alpha_t = 1$

- $\alpha_t \xrightarrow{t \rightarrow 0} 0, \alpha_t \xrightarrow{t \rightarrow \infty} \delta_{\{0\}}$, in distributional sense.

We see that: $[\alpha_t] = 0 = ch(E_+) - ch(E_-)$ (for trivial bundle E , $ch(E) = rank(E)$), this gives no information. But for $t > 0$, α_t represents the generator of $H_c^*(\mathbb{C})$, which is the compact supported cohomology (strictly, the forms are fast decaying). By Poincaré's duality,

$$H_c^*(\mathbb{R}^n) = H^{n-*}(\mathbb{R}^n) = \begin{cases} \mathbb{R} & * = n \\ 0 & * \neq n \end{cases}$$

The isomorphism $H_c^n(\mathbb{R}^n) \cong \mathbb{R}$ is given by integration over \mathbb{R}^n .

Example 1.3.2 (Heat kernel method). H_1, H_2 are Hilbert spaces, a closed linear operator with dense domain $P : H_1 \rightarrow H_2$ is Fredholm iff $dim\ ker P < \infty, dim\ coker P = dim H_2 / im P < \infty$. We can prove that $im P$ is closed, so $coker P \cong (im P)^\perp = ker P^*$, and

$$Ind(P) = dim\ ker P - dim\ coker P = dim\ ker P - dim\ ker P^* = dim\ ker P^* P - dim\ ker P P^*$$

$P^* P$ and $P P^*$ have the same nonzero eigenvalues, and the corresponding eigenspaces are isomorphic:

$$\begin{array}{ccc} E_{P^* P}(\lambda) & \longrightarrow & E_{P P^*}(\lambda) \\ x & \mapsto & P x \\ \frac{1}{\lambda} P^* y & \longleftarrow & y \end{array}$$

We can obtain the heat equation method:

$$\begin{aligned} Ind(P) &= dim\ ker P^* P - dim\ ker P P^* = \sum_{\lambda} e^{-t\lambda} dim E_{P^* P}(\lambda) - \sum_{\lambda} e^{-t\lambda} dim E_{P P^*}(\lambda) \\ &= Tr e^{-tP^* P} - Tr e^{-tP P^*} = Tr_s(e^{-tD^2}) = ch(D) \end{aligned}$$

where $D = \begin{pmatrix} 0 & P^* \\ P & 0 \end{pmatrix}$ is a superconnection on the infinite dimensional Hilbert bundle $H = H_1 \oplus H_2 \rightarrow pt$.

We can also use the transgression method to see this:

$$\frac{\partial}{\partial t} Tr_s e^{-tD^2} = -Tr_s(D^2 e^{-tD^2}) = -\frac{1}{2} Tr_s([D, D] e^{-tD^2}) = -\frac{1}{2} Tr_s[D, D e^{-tD^2}] = 0$$

so $Tr_s e^{-tD^2}$ is a constant. Since $Tr_s e^{-tD^2} \xrightarrow{t \rightarrow \infty} dim\ ker P^* P - dim\ ker P P^* = Ind P$, the formula follows.

1.4 Chern-Simon class

E is a complex vector bundle, ∇^E, ∇'^E are two connections, $R^E = \nabla^{E,2}, R'^E = \nabla'^{E,2}$ are their curvatures. P is an invariant polynomial of $GL(n, \mathbb{C})$, i.e. P satisfies:

$$P(TAT^{-1}) = P(A), \forall T \in GL(n, \mathbb{C}), A \in M_n$$

$P(R^E)$ is closed and $[P(R^E)]$ does not depend on ∇^E , so

$$P(R'^E) - P(R^E) = d\alpha, \text{ for some } \alpha \in \Omega^{odd}(M) / d\Omega^{even}(M)$$

We can have a canonical way to construct α , which is called the Chern-Simon form.

As in the proof of Theorem, choose a curve $C_t = \{\nabla_t^E : t \in [0, 1]\}$ connecting ∇^E and ∇'^E in \mathcal{A} which is the affine space of connections. $\nabla^{\mathbb{E}} = \nabla_t^E + dt \frac{\partial}{\partial t}$ is a connection on the bundle $\pi_M^* E$ over $M \times [0, 1]$, where $\pi_M : M \times [0, 1] \rightarrow M$ is the projection. The curvature $R^{\mathbb{E}} = \nabla^{\mathbb{E},2} = \nabla_t^{E,2} + dt \frac{\partial \nabla_t^E}{\partial t}$.

$P : M_n \cong \mathbb{C}^{n^2} \rightarrow \mathbb{C}$ is a polynomial of n^2 variables, we have the Taylor expansion:

$$P(x+y) = \sum_{|\alpha| \leq \deg P} \frac{\partial^\alpha P(x)}{\alpha!} y^\alpha, \forall x, y \in \mathbb{C}^{n^2}$$

Note that each term of $R^{\mathbb{E}}$ is a 2-form, which commutes with other forms. So we use the above formula to have:

$$P(R^{\mathbb{E}}) = P(R_t^E + dt \frac{\partial}{\partial t} \nabla_t^E) = P(R_t^E) + dt \langle P'(R_t^E), \frac{\partial}{\partial t} \nabla_t^E \rangle$$

where $P'(x) : \mathbb{C}^{n^2} \rightarrow \mathbb{C}$ is the total derivative at $x \in \mathbb{C}^{n^2}$. We know that $P(R^{\mathbb{E}})$ is closed, so:

$$0 = d_{M \times [0,1]} P(R^{\mathbb{E}}) = d_M P(R_t^E) + dt \frac{\partial}{\partial t} P(R_t^E) - dt d_M \langle P'(R_t^E), \frac{\partial}{\partial t} \nabla_t^E \rangle$$

so $\frac{\partial}{\partial t} P(R_t^E) = d \langle P'(R_t^E), \frac{\partial}{\partial t} \nabla_t^E \rangle$. Define

$$\tilde{P} = \int_0^1 \langle P'(R_t^E), \frac{\partial}{\partial t} \nabla_t^E \rangle dt$$

then $d\tilde{P} = P(R^E) - P(R^{\mathbb{E}})$.

This is a case of integration along the fibre, see appendix. We write:

$$\tilde{P} = \int_{C_t} P(R_t^E) = \int_{[0,1]} P(R^{\mathbb{E}})$$

Proposition 1.4.1. *The class of \tilde{P} in $\Omega^{odd}/d\Omega^{even}$ does not depend on the path connecting ∇^E and ∇'^E .*

Proof. If $C'_t = \{\nabla_t'^E\}$ is another path, let $\nabla_{s,t}^E = (1-s)\nabla_t^E + s\nabla_t'^E$, define $\nabla^{\mathcal{E}} = ds \frac{\partial}{\partial s} + dt \frac{\partial}{\partial t} + \nabla_{s,t}^E$, then $\nabla^{\mathcal{E}}$ is a connection on $\mathcal{E} = p_1^* E$, $p_1 : M \times [0,1]^2 \rightarrow M$ is the projection. The curvature:

$$R^{\mathcal{E}} = \nabla^{\mathcal{E},2} = \nabla_{s,t}^{E,2} + dt \frac{\partial}{\partial t} \nabla_{s,t}^E + ds \frac{\partial}{\partial s} \nabla_{s,t}^E = \nabla_{s,t}^{E,2} + dt((1-s) \frac{\partial}{\partial t} \nabla_t^E + s \frac{\partial}{\partial t} \nabla_t'^E) + ds(\nabla_t'^E - \nabla_t^E)$$

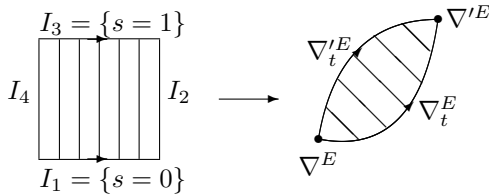
$$R^{\mathcal{E}}|_{s=0} = \nabla_t^{E,2} + dt \frac{\partial}{\partial t} \nabla_t^E = \nabla^{E,2}, \quad R^{\mathcal{E}}|_{s=1} = \nabla_t'^{E,2} + dt \frac{\partial}{\partial t} \nabla_t'^E = \nabla'^{E,2}$$

$$R^{\mathcal{E}}|_{t=0} = \nabla^{E,2}, \quad R^{\mathcal{E}}|_{t=1} = \nabla'^{E,2}$$

Note that $R^{\mathcal{E}}|_{t=0}, R^{\mathcal{E}}|_{t=1}$ have no ds term. Using the Stoke's Formula, we have:

$$\begin{aligned} d \int_{[0,1]^2} P(R^{\mathcal{E}}) &= \int_{[0,1]^2} dP(R^{\mathcal{E}}) - \left(\int_{I_1} P(R^{\mathcal{E}}|_{s=0}) - \int_{I_3} P(R^{\mathcal{E}}|_{s=1}) \right) \\ &= - \int_{[0,1]} P(R^{\mathbb{E}}) + \int_{[0,1]} P(R^{\mathbb{E}'}) = \tilde{P}' - \tilde{P} \end{aligned}$$

□



Given ∇^E, ∇'^E , by the above proposition, we have constructed a class in $\Omega^{odd}/d\Omega^{even}$. We denote it by $\tilde{P}(\nabla^E, \nabla'^E)$, called Chern-Simon form. It satisfies 3 properties:

1. $d\tilde{P}(\nabla^E, \nabla'^E) = P(R'^E) - P(R^E)$
2. $\tilde{P}(\nabla^E, \nabla^E) = 0$ Just choose the constant path: $\nabla_t^E \equiv \nabla^E$.
3. $\tilde{P}(\nabla^E, \nabla'^E)$ is functorial: if $f : N \rightarrow M$ is a differentiable mapping, then

$$\tilde{P}(f^*\nabla^E, f^*\nabla'^E) = f^*\tilde{P}(\nabla^E, \nabla'^E)$$

where $f^*\nabla^E, f^*\nabla'^E$ is the induced connection on $F = f^*E$.

This is because: Choose a curve ∇_t^E connecting ∇^E and ∇'^E , $\nabla_t^F = f^*\nabla_t^E$ is a curve connecting $\nabla^F = f^*\nabla^E$ and $\nabla'^F = f^*\nabla'^E$. Let $\mathbb{E} = \pi_M^*E$, $\nabla^{\mathbb{E}} = \nabla_t^E + dt\frac{\partial}{\partial t}$; $\mathbb{F} = \pi_N^*F$, $\nabla^{\mathbb{F}} = \nabla_t^F + dt\frac{\partial}{\partial t} = (f \times id)^*\nabla^{\mathbb{E}}$.

$$\tilde{P}(f^*\nabla^E, f^*\nabla'^E) = \int_{[0,1]} P(R^{\mathbb{F}}) = \int_{[0,1]} P((f \times id)^*R^{\mathbb{E}}) = f^* \int_{[0,1]} P(R^{\mathbb{E}}) = f^*\tilde{P}(\nabla^E, \nabla'^E)$$

In fact, we have:

Proposition 1.4.2. *The class $\tilde{P}(\nabla^E, \nabla'^E)$ is uniquely defined by properties 1-3*

Proof. We need to check the uniqueness. Assume \tilde{P} satisfies the 3 properties. Choose any path ∇_t^E connecting ∇^E, ∇'^E . Let $\nabla^{\mathbb{E}} = \nabla_t^E + dt\frac{\partial}{\partial t}$, $\nabla_0^{\mathbb{E}} = \nabla^E + dt\frac{\partial}{\partial t}$ be connections on $\mathbb{E} = \pi_M^*E$. By property 1, we have $P(R^{\mathbb{E}}) - P(R_0^{\mathbb{E}}) = d\tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}})$. Integrate this equality along the fibre, and use the Stoke's Formula:

$$\begin{aligned} \int_{[0,1]} P(R^{\mathbb{E}}) &= \int_{[0,1]} P(R_0^{\mathbb{E}}) + \int_{[0,1]} d\tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}}) \\ &= -d \int_{[0,1]} \tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}}) + \tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}})|_{t=1} - \tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}})|_{t=0} \end{aligned}$$

Since $\nabla_0^{\mathbb{E}}|_{t=1} = \nabla^E$, $\nabla^{\mathbb{E}}|_{t=1} = \nabla'^E$, by functoriality, $\tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}})|_{t=1} = \tilde{P}(\nabla^E, \nabla'^E)$. Also $\tilde{P}(\nabla_0^{\mathbb{E}}, \nabla^{\mathbb{E}})|_{t=0} = \tilde{P}(\nabla^E, \nabla^E) = 0$ by property 2. So $\tilde{P}(\nabla^E, \nabla'^E) - \int_{[0,1]} P(R^{\mathbb{E}})$ is exact, that is what we require. \square

Example 1.4.1. λ is a line bundle, $c_1(\lambda)$ is the first Chern class, then $\tilde{c}_1(\nabla^\lambda, \nabla'^\lambda) = -\frac{\nabla'^\lambda - \nabla^\lambda}{2\pi i}$ is the Chern-Simon class. This is because $d(-\frac{\nabla'^\lambda - \nabla^\lambda}{2\pi i}) = \frac{i}{2\pi}r'^\lambda - \frac{i}{2\pi i}r^\lambda = c_1(\nabla'^\lambda) - c_1(\nabla^\lambda)$, and the other two properties are obvious.

Since $\{\text{superconnections}\}$ is an affine space as $\{\text{connections}\}$, everything we've done can be carried out on superconnections, so we have:

Proposition 1.4.3. *The whole theory of Chern-Simon class extends to superconnections*

Proposition 1.4.4. *Let A_t be a curve of superconnections, then*

$$\frac{\partial}{\partial t} Tr_s \exp(-A_t^2) = -d Tr_s \left(\frac{\partial A_t}{\partial t} \exp(-A_t^2) \right)$$

proof 1. Let $\mathbb{A} = A_t + dt\frac{\partial}{\partial t}$ be a superconnection on π_M^*E . By simple computation as example 1.3.1, we see that

$$Tr_s \exp(-\mathbb{A}^2) = Tr_s(\exp(-A_t^2)) - dt Tr_s \left(\frac{\partial A_t}{\partial t} \exp(-A_t^2) \right)$$

We know that $Tr_s \exp(-\mathbb{A}^2)$ is a closed even form on $M \times \mathbb{R}$, so

$$\frac{\partial}{\partial t} Tr_s \exp(-A_t^2) = -d Tr_s \left(\frac{\partial A_t}{\partial t} \exp(-A_t^2) \right)$$

\square

proof 2.

$$\begin{aligned}\frac{\partial}{\partial t} Tr_s \exp(-A_t^2) &= -Tr_s \left(\frac{\partial A_t^2}{\partial t} \exp(-A_t^2) \right) = -Tr_s \left([A_t, \frac{\partial A_t}{\partial t} \exp(-A_t^2)] \right) \\ &= -Tr_s \left([A_t, \frac{\partial A_t}{\partial t} \exp(-A_t^2)] \right) = -d Tr_s \left(\frac{\partial A_t}{\partial t} \exp(-A_t^2) \right)\end{aligned}$$

We used identities:

- $\frac{\partial A_t^2}{\partial t} = \frac{\partial A_t}{\partial t} A_t + A_t \frac{\partial A_t}{\partial t} = [A_t, \frac{\partial A_t}{\partial t}]$ (note both A_t and $\frac{\partial A_t}{\partial t}$ are odd)
- $d Tr_s(B) = Tr_s[A, B]$, for A a superconnection and $B \in \Gamma(\mathcal{B})$ (see the proof of theorem 1.3.1)

□

When $A_t = \nabla^E + \sqrt{t}V$, $\nabla^E = \nabla^{E+} \oplus \nabla^{E-}$, $V = \begin{pmatrix} 0 & V_- \\ V_+ & 0 \end{pmatrix} \in \Gamma(End(E)^{odd})$,

$$\frac{\partial}{\partial t} Tr_s \exp(-A_t^2) = -d Tr_s \left(\frac{V}{2\sqrt{t}} \exp(-A_t^2) \right)$$

We want to integrate this equality on $[0, \infty)$.

Now assume $g^E = g^{E+} \oplus g^{E-}$ is an Hermitian metric on E , ∇^E is a unitary connection, and V is self-adjoint, so $V_- = V_+^*$.

Proposition 1.4.5. $\alpha_t = ch(A_t)$ is a real form.

Proof. We define an adjoint operator $*$ on $\mathcal{B} = \wedge^* T^* M \hat{\otimes} End(E)$:

1. $\forall \omega \in \wedge_{\mathbb{C}}^1 T^* M$, $\omega^* = -\bar{\omega}$
2. $\forall A \in End(E)$, A^* is the adjoint of A
3. $\forall f, g \in \mathcal{B}$, $(f \cdot g)^* = g^* \cdot f^*$.

This is well defined. Now $A_t^2 = \nabla^{E,2} + \sqrt{t}\nabla^E V + tV^2$. $(tV^2)^* = tV^2$, since V is self-adjoint. $\forall X \in TM$, $\nabla_X V$ is also self-adjoint, so $(\nabla^E V)^* = (e^i \nabla_{e_i}^E V)^* = \nabla_{e_i}^E V \cdot (-e^i) = e^i \nabla_{e_i}^E V = \nabla^E V$. Since ∇^E is unitary, $R^E(X, Y) = \nabla^{E,2}(X, Y)$ is skew-Hermitian w.r.t. g^E , $\forall X, Y \in TM$, so

$$(\nabla^{E,2})^* = (e^i \wedge e^j R^E(e_i, e_j))^* = -R^E(e_i, e_j)(-e^j) \cdot (-e^i) = e^i \wedge e^j R^E(e_i, e_j) = \nabla^{E,2}$$

So we see $(A_t^2)^* = A_t^2$, also $(\exp(-A_t^2))^* = \exp(-A_t^2)$. Note that $\forall \omega A \in \mathcal{B}_{even}$,

$$(\omega A)^* = A^* \omega^* = A^* (-1)^{deg \omega} \cdot (-1)^{\frac{deg \omega \cdot (deg \omega - 1)}{2}} \bar{\omega} = (-1)^{\frac{deg \omega \cdot (deg \omega - 1)}{2}} \bar{\omega} A^*$$

so

$$\overline{Tr_s(\omega A)} = Tr_s(\bar{\omega} A^*) = (-1)^{\frac{deg \omega \cdot (deg \omega - 1)}{2}} Tr_s((\omega A)^*) = \begin{cases} Tr_s((\omega A)^*) & deg \omega \equiv 0 \pmod{4} \\ -Tr_s((\omega A)^*) & deg \omega \equiv 2 \pmod{4} \\ 0 & \text{otherwise} \end{cases}$$

So $(Tr_s \exp(-A_t^2))^{(4k)}$ is real, and $(Tr_s \exp(-A_t^2))^{(4k+2)}$ is purely imaginary, then it's clear $\varphi Tr_s \exp(-A_t^2) = ch(A_t)$ is real. □

Remark 5. Similarly, $\forall \omega A \in \mathcal{B}_{odd}$, $(\omega A)^* = (-1)^{\frac{deg \omega \cdot (deg \omega + 1)}{2}} \bar{\omega} A^*$, so $\forall B \in \mathcal{B}_{odd}$ satisfying $B^* = B$, we have $(Tr_s(B))^{(4k+1)}$ is purely imaginary, and $Tr_s(B)^{(4k+3)}$ is real, then $\frac{1}{\sqrt{2\pi i}} \varphi Tr_s(B)$ is real.

Proposition 1.4.6. Assume $ker V$ has locally constant dimensions, then $ker V$ is a smooth sub-bundle of E .

Proof. Since $\ker V$ has locally constant dimensions, $\forall x \in M$, $\exists \varepsilon > 0$, and a neighborhood $U \ni x$, s.t. $\forall y \in U$, $V_y : E_y \rightarrow E_y$ has no nonzero eigenvalues in the disk $B(0, \varepsilon) = \{z \in \mathbb{C} : |z| \leq \varepsilon\}$. Let

$$P^{\ker V} = \frac{1}{2\pi i} \int_{\partial B(0, \varepsilon)} \frac{d\lambda}{\lambda - V}$$

then $P^{\ker V} : E \rightarrow E$ is a smooth projection onto $\ker V$. We have the direct sum:

$$E = P^{\ker V}(E) \oplus (1 - P^{\ker V})(E) = \ker V \oplus (\ker V)^\perp = \ker V \oplus \operatorname{im} V$$

□

Since $\ker V = \ker V_+ \oplus \ker V_-$, and $\dim \ker V_- = \dim(\operatorname{im} V_+)^\perp = \dim E_- - (\dim E_+ - \dim \ker V_+)$,

$$\dim \ker V_+ + \dim \ker V_- = \dim \ker V, \quad \dim \ker V_+ - \dim \ker V = \dim E_+ - \dim E_-$$

So $\dim \ker V_+$, $\dim \ker V_-$ is locally constant, by the proposition $\ker V_\pm$ is subbundles of E_\pm . Let $g^{\ker V}$ be the metric on $\ker V$ induced by g^E , and $\nabla^{\ker V} = P^{\ker V} \nabla^E$ be the induced connection on $\ker V$ by orthogonal projection, then

$$\nabla^{\ker V} = P^{\ker V}(\nabla^{E_+} \oplus \nabla^{E_-}) = P^{\ker V_+} \nabla^{E_+} \oplus P^{\ker V_-} \nabla^{E_-} = \nabla^{\ker V_+} \oplus \nabla^{\ker V_-}$$

Proposition 1.4.7. $ch(E) = ch(\ker V)$ in $H^*(M)$, where $ch(E) = ch(E_+) - ch(E_-)$, $ch(\ker V) = ch(\ker V_+) - ch(\ker V_-)$.

Proof. We have an exact sequence of vector bundles:

$$\begin{array}{ccccccc} 0 & \longrightarrow & \ker V_+ & \longrightarrow & E_+ & \xrightarrow{V_+} & E_- & \longrightarrow & \ker V_- & \longrightarrow & 0 \\ & & & & \parallel & & \parallel & & & & \\ & & & & \ker V_+ \oplus (\ker V_+)^\perp & \xrightarrow{V_+} & (\ker V_-)^\perp \oplus \ker V_- & & & & \end{array}$$

since we can take any connection to compute $ch(E)$, we let

$$\nabla^{E_+} = \nabla^{\ker V_+} \oplus \nabla^{(\ker V_+)^\perp}, \quad \nabla^{E_-} = \nabla^{(\ker V_-)^\perp} \oplus \nabla^{\ker V_-}$$

where $\nabla^{(\ker V_+)^\perp} = V_+^* \nabla^{(\ker V_-)^\perp}$, then

$$\begin{aligned} ch(E) &= ch(E_+) - ch(E_-) = [ch(\nabla^{E_+})] - [ch(\nabla^{E_-})] \\ &= [ch(\nabla^{\ker V_+})] - [ch(\nabla^{\ker V_-})] \end{aligned}$$

□

Theorem 1.4.1. When $t \rightarrow +\infty$, we have asymptotic formula:

$$ch(A_t) = ch(\nabla^{\ker V}) + O\left(\frac{1}{\sqrt{t}}\right)$$

uniformly on compact sets.

Proof. We first compute $\nabla^{\ker V, 2}$. Consider the orthogonal splitting $E = \ker V \oplus (\ker V)^\perp$, $\nabla^E = \nabla^{split} + A$, where $\nabla^{split} = \nabla^{\ker V} \oplus \nabla^{(\ker V)^\perp}$. In blocked matrix with respect to the new splitting, because ∇^E is unitary,

$$\nabla^E = \begin{pmatrix} \nabla^{\ker V} & 0 \\ 0 & \nabla^{(\ker V)^\perp} \end{pmatrix} + \begin{pmatrix} 0 & -B^* \\ B & 0 \end{pmatrix}$$

Because $\forall s \in \Gamma(\ker V)$, $X \in TM$,

$$(\nabla_X^E V)s = \nabla_X^E(Vs) - V(\nabla_X^E s) = -V(\nabla_X^E s) \in \operatorname{im} V$$

$\nabla^E V$ is a 1-form with value in $End(E)$, which maps $kerV$ into $(kerV)^\perp = imV$, and $P^{(kerV)^\perp} \nabla_X^E s = -V^{-1}(\nabla_X^E V)s$, where $V : (kerV)^\perp \rightarrow imV = (kerV)^\perp$ is an isomorphism. So

$$B(X) = P^{(kerV)^\perp} \nabla_X^E \cdot P^{kerV} = -V^{-1}(\nabla_X^E V)P^{kerV}, \quad B(X)^* = -P^{kerV} \cdot \nabla_X^E V \cdot V^{-1}$$

or in our sign rule

$$B = V^{-1} \cdot (\nabla^E V) \cdot P^{kerV}, \quad B^* = -P^{kerV} \nabla^E V \cdot V^{-1}$$

so

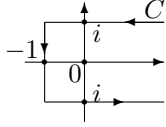
$$A \triangleq \begin{pmatrix} 0 & -B^* \\ B & 0 \end{pmatrix} = \begin{pmatrix} 0 & P^{kerV} \nabla^E V \cdot V^{-1} \\ V^{-1} \nabla^E V & 0 \end{pmatrix}$$

$\nabla^{E,2} = \nabla^{split,2} + \nabla^{split} A + A^2$, since $\nabla^{split} A$ interchanges $kerV$ and $(kerV)^\perp$,

$$\begin{aligned} \nabla^{kerV,2} &= P^{kerV} \nabla^{E,2} P^{kerV} - P^{kerV} A^2 P^{kerV} \\ &= P^{kerV} (\nabla^{E,2} - \nabla^E V \cdot V^{-2} \cdot \nabla^E V) P^{kerV} \end{aligned}$$

we next take the limit. $A_t^2 = tV^2 + \sqrt{t} \nabla^E V + \nabla^{E,2} \in \Gamma(\mathcal{B}_{even})$, $\sqrt{t} \nabla^E V + \nabla^{E,2}$ is nilpotent in \mathcal{B} because they contain Grassmannian variables. So the spectrum of A_t^2 in \mathcal{B} is the same as tV^2 . Since tV^2 is positive, $Sp(A_t^2) = Sp(tV^2) \subset \mathbb{R}_+$. We choose a contour C as in the figure and use the Cauchy integral formula:

$$\exp(-A_t^2) = \frac{1}{2\pi i} \int_C \frac{\exp(-\lambda) d\lambda}{\lambda - A_t^2}$$



We can compute $(\lambda - A_t^2)^{-1}$ by the following

Lemma 1. $M = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$ is a blocked matrix, then

$$(\lambda I - M)^{-1} = \begin{pmatrix} \alpha & \alpha B(\lambda - D)^{-1} \\ (\lambda - D)^{-1} C \alpha & (\lambda - D)^{-1} (1 + C \alpha B(\lambda - D)^{-1}) \end{pmatrix}$$

where

$$\alpha = (\lambda - A - B(\lambda - D)^{-1} C)^{-1}$$

the proof is just computation.

$$\begin{aligned} A_t^2 &= \begin{pmatrix} 0 & 0 \\ 0 & tV^2 \end{pmatrix} + \sqrt{t} \begin{pmatrix} 0 & P^{kerV} \nabla^E V P^{(kerV)^\perp} \\ P^{(kerV)^\perp} \nabla^E V P^{kerV} & P^{(kerV)^\perp} \nabla^E V P^{(kerV)^\perp} \end{pmatrix} + \\ &\quad \begin{pmatrix} P^{kerV} \nabla^{E,2} P^{kerV} & P^{kerV} \nabla^{E,2} P^{(kerV)^\perp} \\ P^{(kerV)^\perp} \nabla^{E,2} P^{kerV} & P^{(kerV)^\perp} \nabla^{E,2} P^{(kerV)^\perp} \end{pmatrix} \\ &= \begin{pmatrix} P^{kerV} \nabla^{E,2} P^{kerV} & \sqrt{t} P^{kerV} \nabla^E V P^{(kerV)^\perp} + O(1) \\ \sqrt{t} P^{(kerV)^\perp} \nabla^E V P^{kerV} + O(1) & tV^2 + O(\sqrt{t}) \end{pmatrix} \end{aligned}$$

Using the lemma, we have

$$\begin{aligned} \alpha_t &= [\lambda - P^{kerV} \nabla^{E,2} P^{kerV} - (\sqrt{t} P^{kerV} \nabla^E V P^{(kerV)^\perp} + O(1))(\lambda - tV^2 - O(\sqrt{t}))^{-1} \\ &\quad (\sqrt{t} P^{(kerV)^\perp} \nabla^E V P^{kerV} + O(1))]^{-1} \\ &= (\lambda - P^{kerV} (\nabla^{E,2} - \nabla^E V V^{-2} \nabla^E V) P^{kerV})^{-1} + O\left(\frac{1}{\sqrt{t}}\right) \\ &= (\lambda - \nabla^{kerV,2})^{-1} + O\left(\frac{1}{\sqrt{t}}\right) \end{aligned}$$

and estimates for other blocks:

$$\begin{array}{ccccccc} \alpha & B & (\lambda - D)^{-1} & = O(\frac{1}{\sqrt{t}}), & (\lambda - D)^{-1} & (1 + C & \alpha & B & (\lambda - D)^{-1} &) = O(\frac{1}{t}) \\ \downarrow & \downarrow & \downarrow & & \downarrow & \downarrow & \downarrow & \downarrow & \downarrow & \\ 1 & \sqrt{t} & t^{-1} & & t^{-1} & \sqrt{t} & 1 & \sqrt{t} & t^{-1} & \end{array}$$

$$(\lambda - D)^{-1}C\alpha = O(\frac{1}{\sqrt{t}})$$

$$\text{So } (\lambda - A_t^2)^{-1} = \begin{pmatrix} (\lambda - \nabla^{kerV,2})^{-1} & 0 \\ 0 & 0 \end{pmatrix} + O(\frac{1}{\sqrt{t}}), \text{ and}$$

$$\exp(-A_t^2) = \frac{1}{2\pi i} \int_C \frac{\exp(-\lambda)}{\lambda - A_t^2} d\lambda = \frac{1}{2\pi i} \int_C \frac{\exp(-\lambda)}{\lambda - \nabla^{kerV,2}} P^{kerV} + O(\frac{1}{\sqrt{t}})$$

$$ch(A_t) = ch(\nabla^{kerV}) + O(\frac{1}{\sqrt{t}})$$

as $t \rightarrow +\infty$, uniformly on compact sets of M . □

Proposition 1.4.8.

$$Tr_s(\sqrt{t}V \exp(-A_t^2)) = O(\frac{1}{t}), \text{ as } t \rightarrow +\infty$$

Proof. Consider the projection $\pi : M \times (0, \infty) \rightarrow M$, $\nabla^{\mathbb{E}} = \nabla^E + da \frac{\partial}{\partial a}$ is a connection on $\mathbb{E} = \pi^* E$. Let $\mathbb{A}_t = \nabla^{\mathbb{E}} + \sqrt{t}aV$ be a superconnection.

$$\mathbb{A}_t^2 = \nabla^{E,2} + ta^2V^2 + \sqrt{t}a\nabla^E V + da\sqrt{t}V$$

So

$$Tr_s(\exp(-\mathbb{A}_t^2)) = Tr_s(\exp(-A_t(a)^2)) - da Tr_s(\sqrt{t}V \exp(-A_t(a)^2))$$

where $A_t(a) = \nabla^E + \sqrt{t}aV$.

Note $ker(aV) = \pi^* kerV$, by the above theorem, $Tr_s(\exp(-A_t(a)^2)) = Tr_s(\exp(-\nabla^{kerV,2})) + O(\frac{1}{\sqrt{t}})$, as $t \rightarrow +\infty$, so

$$Tr_s(\exp(-\mathbb{A}_t^2)) = Tr_s(\exp(-\nabla^{kerV,2})) - da Tr_s(\sqrt{t}V \exp(-A_t(a)^2)) + O(\frac{1}{\sqrt{t}})$$

On the other hand, by the theorem again,

$$Tr_s(\exp(-\mathbb{A}_t^2)) = Tr_s(\exp(-\nabla^{ker(aV),2})) + O(\frac{1}{\sqrt{t}}) = Tr_s(\exp(-\nabla^{kerV,2})) + O(\frac{1}{\sqrt{t}})$$

We used the fact: $\nabla^{ker(aV)} = \pi^* \nabla^{kerV} + da \frac{\partial}{\partial a}$, and so $\nabla^{ker(aV),2} = \pi^* \nabla^{kerV,2}$. Comparing the above two equation, and note that $A_t(1) = A_t$, we see that

$$Tr_s(\sqrt{t}V \exp(-A_t^2)) = O(\frac{1}{\sqrt{t}})$$

□

By the above proposition, we have

$$\frac{\partial}{\partial t} Tr_s \exp(-A_t^2) = -d Tr_s \left(\frac{V}{2\sqrt{t}} \exp(-A_t^2) \right) = \begin{cases} O(\frac{1}{t^{\frac{3}{2}}}) & \text{as } t \rightarrow +\infty \\ O(\frac{1}{\sqrt{t}}) & \text{as } t \rightarrow 0 \end{cases}$$

So we can integrate it from 0 to $+\infty$.

Definition 1.4.1. $\beta = \frac{1}{\sqrt{2\pi i}} \int_0^{+\infty} \varphi Tr_s \left(\frac{V}{2\sqrt{t}} \exp(-A_t^2) \right) dt$

Theorem 1.4.2. β is an odd real form and

$$d\beta = ch(\nabla^E) - ch(\nabla^{kerV})$$

Proof. By the Remark following proposition 1.4.5, β is an odd real form. By theorem 1.4.1 and proposition 1.4.4

$$\begin{aligned} Tr_s(\exp(-\nabla^{kerV,2})) - Tr_s(\exp(-\nabla^{E,2})) &= \int_0^\infty \frac{\partial}{\partial t} Tr_s(\exp(-A_t^2)) dt \\ &= -d \int_0^\infty Tr_s\left(\frac{V}{2\sqrt{t}} \exp(-A_t^2)\right) dt \end{aligned}$$

Since $\varphi \circ d = \frac{1}{\sqrt{2i\pi}} d \circ \varphi$,

$$\varphi Tr_s(\exp(-\nabla^{E,2})) - \varphi Tr_s(\exp(-\nabla^{kerV,2})) = \frac{1}{\sqrt{2i\pi}} d \int_0^\infty \varphi Tr_s\left(\frac{V}{2\sqrt{t}} \exp(-A_t^2)\right) dt$$

i.e. $ch(\nabla^E) - ch(\nabla^{kerV}) = d\beta$. □

We will show β is a Chern-Simon class. First we generalize the definition.

Theorem 1.4.3. Given an exact sequence of vector bundles:

$$E : 0 \longrightarrow E_1 \xrightarrow{v_1} E_2 \xrightarrow{v_2} \dots \xrightarrow{v_{n-1}} E_n \longrightarrow 0$$

and connections $\nabla^E = \{\nabla^{E_i}\}$ on $E = \{E_i, v_i\}$, there is a unique way to associate a class $\widetilde{ch}(\nabla^E)$ in $\Omega^{odd}(M)/d\Omega^{even}$, which satisfies 3 properties:

1. $d\widetilde{ch}(E) = ch(\nabla^{E_1}) - ch(\nabla^{E_2}) + \dots + (-1)^{n-1} ch(\nabla^{E_n}) =: ch(\nabla^E)$

2. If the sequence (E, ∇^E) splits, $\widetilde{ch}(\nabla^E) = 0$. “ (E, ∇^E) splits” means:

$$E : 0 \longrightarrow E_1 \xrightarrow{v_1} E'_2 \oplus E''_2 \xrightarrow{v_2} \dots \xrightarrow{v_{n-1}} E_n \longrightarrow 0$$

and $\nabla^{E_i} = \nabla^{E'_i} \oplus \nabla^{E''_i}$, $\nabla^{E''_i} = v_i^* \nabla^{E'_{i+1}}$, where $E'_i = \text{im} v_{i-1} = \text{ker} v_i$, $E''_i \cong E_i/E'_i$.

3. $\widetilde{ch}(\nabla^E)$ is functorial.

Proof. Fix a splitting: $E_i = E'_i \oplus E''_i$, $\nabla^{E_i, split} = \nabla^{E'_i} \oplus \nabla^{E''_i}$, $\nabla^{E''_i} = v_i^* \nabla^{E'_{i+1}}$. Let

$$\alpha(\nabla^E) = \sum_{i=1}^n (-1)^{i-1} \widetilde{ch}(\nabla^{E_i, split}, \nabla^{E_i})$$

where $\widetilde{ch}(\nabla^E, \nabla^{E'})$ is the Chern-Simon class defined before (see proposition 1.4.2). We'll show $\alpha(\nabla^E)$ satisfies the 3 properties.

1. Since $\sum_{i=1}^n (-1)^{i-1} ch(\nabla^{E_i, split}) = 0$,

$$\begin{aligned} d\alpha(\nabla^E) &= \sum_{i=1}^n (-1)^{i-1} d\widetilde{ch}(\nabla^{E_i, split}, \nabla^{E_i}) = \sum_{i=1}^n (-1)^{i-1} (ch(\nabla^{E_i}) - ch(\nabla^{E_i, split})) \\ &= \sum_{i=1}^n (-1)^{i-1} ch(\nabla^{E_i}) = ch(\nabla^E) \end{aligned}$$

2. First consider another splitting of the kind: $E_i = E'_i \oplus \widetilde{E}''_i$, $\nabla^{E_i} = \nabla^{E'_i} \oplus \nabla^{\widetilde{E}''_i}$, $\nabla^{\widetilde{E}''_i} = v_i^* \nabla^{E'_{i+1}}$, i.e. the bundles have a new splitting $E_i = E'_i \oplus \widetilde{E}''_i$, but we use the fixed $\nabla^{E'_i}$. They together determine the new splitting connections $\nabla^{E_i} = \nabla^{E'_i} \oplus \nabla^{\widetilde{E}''_i}$.

Claim 1. $\widetilde{ch}(\nabla^{E_i, split}, \nabla^{E_i}) = 0$.

proof of the Claim. Let $\tilde{P}^{E'_i}$ and $\tilde{P}^{\tilde{E}''_i}$ be the projection w.r.t this new splitting, we have isomorphisms

$$(E''_i, \nabla^{E''_i}) \begin{array}{c} \xrightarrow{\tilde{P}^{\tilde{E}''_i}|_{E''_i} : E''_i \xrightarrow{\cong} \tilde{E}''_i} \\ \xleftarrow{P^{E''_i}|_{\tilde{E}''_i} : \tilde{E}''_i \xrightarrow{\cong} E''_i} \end{array} (\tilde{E}''_i, \nabla^{\tilde{E}''_i})$$

So $\nabla^{E_i} = \nabla^{E'_i} \tilde{\oplus} (P^{E''_i})^* (\nabla^{E''_i})$. $\forall s \in E$, $s = u + v = \tilde{u} + \tilde{v}$ is decompositions w.r.t the two splittings, we calculate

$$\begin{aligned} P^{E'_i} (\nabla^{E'_i} \tilde{\oplus} (P^{E''_i})^* (\nabla^{E''_i})) P^{E'_i} s &= P^{E'_i} \nabla^{E'_i} \tilde{P}^{E'_i} P^{E'_i} s + P^{E'_i} (P^{E''_i})^{-1} \nabla^{E''_i} P^{E''_i} \tilde{P}^{\tilde{E}''_i} P^{E''_i} s \\ &= \nabla^{E'_i} P^{E'_i} s \end{aligned}$$

$$\begin{aligned} P^{E''_i} (\nabla^{E'_i} \tilde{\oplus} (P^{E''_i})^* (\nabla^{E''_i})) P^{E''_i} s &= P^{E''_i} \nabla^{E'_i} \tilde{P}^{E'_i} P^{E''_i} s + P^{E''_i} ((P^{E''_i})^{-1} \nabla^{E''_i} P^{E''_i}) \tilde{P}^{\tilde{E}''_i} P^{E''_i} s \\ &= \nabla^{E''_i} P^{E''_i} s \end{aligned}$$

$$P^{E''_i} (\nabla^{E'_i} \tilde{\oplus} (P^{E''_i})^* (\nabla^{E''_i})) P^{E''_i} s = P^{E''_i} \nabla^{E'_i} \tilde{P}^{E'_i} P^{E''_i} s + P^{E''_i} (P^{E''_i})^{-1} \nabla^{E''_i} P^{E''_i} \tilde{P}^{\tilde{E}''_i} P^{E''_i} s = 0$$

$$\begin{aligned} P^{E'_i} (\nabla^{E'_i} \tilde{\oplus} (P^{E''_i})^* (\nabla^{E''_i})) P^{E''_i} s &= P^{E'_i} \nabla^{E'_i} \tilde{P}^{E'_i} P^{E''_i} s + P^{E'_i} (P^{E''_i})^{-1} \nabla^{E''_i} P^{E''_i} \tilde{P}^{\tilde{E}''_i} P^{E''_i} s \\ &= \nabla^{E'_i} \tilde{P}^{E'_i} P^{E''_i} s =: B \end{aligned}$$

In the above calculation, we use some identities, like

$$\tilde{P}^{E'_i} P^{E'_i} = P^{E'_i}, \tilde{P}^{E''_i} P^{E''_i} = 0, \tilde{P}^{\tilde{E}''_i} P^{E''_i} = P^{E''_i}, P^{E'_i} (P^{E''_i})^{-1} = P^{E'_i} \tilde{P}^{\tilde{E}''_i} = 0$$

So with respect to the initial splitting $E_i = E'_i \oplus E''_i$, we have

$$\nabla^{E_i} = \begin{pmatrix} \nabla^{E'_i} & 0 \\ 0 & \nabla^{E''_i} \end{pmatrix} + \begin{pmatrix} 0 & B \\ 0 & 0 \end{pmatrix} := \nabla^{E_i, split} + D$$

Let $\nabla_t^{E_i} = \nabla^{E_i, split} + tD$ be a curve connecting $\nabla^{E_i, split}$ and ∇^{E_i} . It's easy to see that

$$ch(\nabla_t^{E_i}) \equiv ch(\nabla^{E_i, split})$$

So $\widetilde{ch}(\nabla^{E_i, split}, \nabla^{E_i}) = 0$. □

Now we consider the general splitting $E_i = E'_i \tilde{\oplus} \tilde{E}''_i$, $\tilde{\nabla}^{E_i} = \tilde{\nabla}^{E'_i} \tilde{\oplus} \tilde{\nabla}^{\tilde{E}''_i}$, $\tilde{\nabla}^{\tilde{E}''_i} = v_i^* \tilde{\nabla}^{E'_{i+1}}$. From the definition, we can see that

$$\widetilde{ch}(\nabla^{E_i, split}, \tilde{\nabla}^{E_i}) = \widetilde{ch}(\nabla^{E_i, split}, \nabla^{E_i}) + \widetilde{ch}(\nabla^{E_i}, \tilde{\nabla}^{E_i})$$

By the above construction, we have

$$\begin{aligned} \widetilde{ch}(\nabla^{E_i}, \tilde{\nabla}^{E_i}) &= \widetilde{ch}(\nabla^{E'_i} \tilde{\oplus} \nabla^{\tilde{E}''_i}, \tilde{\nabla}^{E'_i} \tilde{\oplus} \tilde{\nabla}^{\tilde{E}''_i}) \\ &= \widetilde{ch}(\nabla^{E'_i}, \tilde{\nabla}^{E'_i}) + \widetilde{ch}(v_i^* \nabla^{E'_{i+1}}, v_i^* \tilde{\nabla}^{E'_{i+1}}) \\ &= \widetilde{ch}(\nabla^{E'_i}, \tilde{\nabla}^{E'_i}) + \widetilde{ch}(\nabla^{E'_{i+1}}, \tilde{\nabla}^{E'_{i+1}}) \end{aligned}$$

Using the Claim,

$$\begin{aligned}
\alpha(\tilde{\nabla}^E) &= \sum_{i=1}^n (-1)^{i-1} \widetilde{ch}(\nabla^{E_i, split}, \tilde{\nabla}^{E_i}) \\
&= \sum_{i=1}^n (-1)^{i-1} (\widetilde{ch}(\nabla^{E_i, split}, \nabla^{E_i}) + \widetilde{ch}(\nabla^{E_i}, \tilde{\nabla}^{E_i})) \\
&= \sum_{i=1}^n (-1)^{i-1} \widetilde{ch}(\nabla^{E_i}, \tilde{\nabla}^{E_i}) \\
&= \sum_{i=1}^n (-1)^{i-1} (\widetilde{ch}(\nabla^{E'_i}, \tilde{\nabla}^{E'_i}) + \widetilde{ch}(\nabla^{E'_{i+1}}, \tilde{\nabla}^{E'_{i+1}})) \\
&= 0
\end{aligned}$$

3. By the functoriality of $\widetilde{ch}(\nabla^E, \nabla'^E)$, this is clear.

So we have the existence. We prove the uniqueness. The proof is same as the proof of proposition 1.4.2.

Assume $\alpha(\nabla^E)$ is any class satisfying the 3 properties. Let ∇_t^E be any path connecting $\nabla^{E, split} = \{\nabla^{E_i, split}\}$ and $\nabla^E = \{\nabla^{E_i}\}$. Let $\nabla^{\mathbb{E}} = \nabla_t^E + dt \frac{\partial}{\partial t}$, then $\nabla^{\mathbb{E}}|_{t=0} = \nabla^{E, split}$, $\nabla^{\mathbb{E}}|_{t=1} = \nabla^E$, so

$$\begin{aligned}
\int_{[0,1]} ch(\nabla^{\mathbb{E}}) &= \int_{[0,1]} d\alpha(\nabla^{\mathbb{E}}) \\
&= -d \int_{[0,1]} \alpha(\nabla^{\mathbb{E}}) + \alpha(\nabla^{\mathbb{E}})|_{t=1} - \alpha(\nabla^{\mathbb{E}})|_{t=0} \\
&= -d \int_{[0,1]} \alpha(\nabla^{\mathbb{E}}) + \alpha(\nabla^E)
\end{aligned}$$

Note that $\int_{[0,1]} ch(\nabla^{\mathbb{E}})$ is what we've constructed. □

We can now show $\beta = \frac{1}{\sqrt{2i\pi}} \int_0^\infty \varphi Tr_s(\frac{V}{2\sqrt{t}} \exp(-A_t^2)) dt$ is a Chern-Simon class.

Theorem 1.4.4. *We have an exact sequence of vector bundles*

$$0 \longrightarrow \ker V_+ \longrightarrow E_+ \xrightarrow{V_+} E_- \longrightarrow \text{coker } V_+ \cong \ker V_- \longrightarrow 0$$

and corresponding connections $\nabla^{\ker V_+}, \nabla^{E_+}, \nabla^{E_-}, \nabla^{\ker V_-}$.

$\beta = -\widetilde{ch}$ in $\Omega^{\text{odd}}/d\Omega^{\text{even}}$.

Proof. By proposition 1.4.2, β is a real odd form, and

$$d\beta = ch(\nabla^E) - ch(\nabla^{\ker V}) = -(ch(\nabla^{\ker V_+}) - ch(\nabla^{E_+}) + ch(\nabla^{E_-}) - ch(\nabla^{\ker V_-}))$$

The functorial property is easy. We need to prove β satisfies the 2nd property. If the sequence splits, we assume:

$$\nabla^{E_+} = \nabla^{\ker V_+} \oplus \nabla^{(\ker V_+)^{\perp}}, \nabla^{E_-} = \nabla^{\ker V_-} \oplus \nabla^{(\ker V_-)^{\perp}}, \nabla^{(\ker V_+)^{\perp}} = V_+^* \nabla^{(\ker V_-)^{\perp}}$$

Under the splitting $E = E_+ \oplus E_- = \ker V_+ \oplus (\ker V_+)^{\perp} \oplus \ker V_- \oplus (\ker V_-)^{\perp}$, we have

$$\nabla^E = \begin{pmatrix} \begin{pmatrix} \nabla^{\ker V_+} & 0 \\ 0 & \nabla^{(\ker V_+)^{\perp}} \end{pmatrix} & 0 \\ 0 & \begin{pmatrix} \nabla^{\ker V_-} & 0 \\ 0 & \nabla^{(\ker V_-)^{\perp}} \end{pmatrix} \end{pmatrix}$$

$$V = \begin{pmatrix} 0 & \begin{pmatrix} 0 & 0 \\ 0 & V_- \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ 0 & V_+ \end{pmatrix} & 0 \end{pmatrix}, \nabla^E V = \begin{pmatrix} 0 & \begin{pmatrix} 0 & 0 \\ 0 & \nabla^E V_- \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ 0 & \nabla^E V_+ \end{pmatrix} & 0 \end{pmatrix} = 0$$

$$\begin{aligned} A_t^2 &= \nabla^{E,2} + \sqrt{t} \nabla^E V + tV^2 \\ &= \begin{pmatrix} \begin{pmatrix} \nabla^{ker V_+,2} & 0 \\ 0 & \nabla^{(ker V_+)^{\perp},2} + tV_- V_+ \end{pmatrix} & 0 \\ 0 & \begin{pmatrix} \nabla^{ker V_-,2} & 0 \\ 0 & \nabla^{(ker V_-)^{\perp},2} + tV_+ V_- \end{pmatrix} \end{pmatrix} \end{aligned}$$

So $Tr_s(V \exp(-A_t^2)) = Tr_s \begin{pmatrix} 0 & * \\ * & 0 \end{pmatrix} = 0$, so $\beta = 0$.

So β is indeed the Chern-Simon class of the sequence. □

Chapter 2

Atiyah-Singer Family Index Theorem

2.1 Family Index Theorem

$X \rightarrow M \xrightarrow{\pi} S$ is a fibration with compact fibre X of $\dim 2l$. TX (subbundle of TM) is even dimensional, orientable and spin. This is a global condition over M . g^{TX} is a metric on TX . E is a complex vector bundle on M with Hermitian metric g^E and unitary connection ∇^E . $S^{TX} = S_+^{TX} \oplus S_-^{TX}$ is spinor bundle of TX , $\text{rank}(S^{TX}) = 2^{\frac{\dim X}{2}} = 2^l$. $\forall s \in S$, D_s^X is the Dirac operator along the fibre X_s with twisting bundle $E|_{X_s}$, which acts on $\Gamma(X_s, S^{TX} \otimes E|_{X_s})$. $D_s^X = \begin{pmatrix} 0 & D_{-,s}^X \\ D_{+,s}^X & 0 \end{pmatrix}$ is self-adjoint, it's Fredholm. $\text{Ind}(D_{+,s}^X) = \dim \ker D_{+,s}^X - \dim \ker D_{-,s}^X \in \mathbb{Z}$ does not depend on s by the homotopy invariance of the index. By Atiyah-Singer index theorem,

$$\text{Ind} D_{+,s}^X = \int_{X_s} \hat{A}(TX) \text{ch}(E)$$

The idea of family index is to see $\text{Ind} D_{+,s}^X$ as “vector bundle” on S , s.t. “rank”($\text{Ind} D_{+,s}^X$)=classical index.

Suppose that $\ker D_{+,s}^X$ and $\ker D_{-,s}^X$ have constant dimensions, then one can prove easily that $\ker D_{\pm,s}^X$ are the fibres of smooth vector bundles. Let $\mathcal{V}(M)$ denote the class of vector bundles on M , then we have

$$\begin{aligned} \mathcal{V}(M) &\longrightarrow \mathcal{V}(S) \\ E &\longmapsto \ker D_-^X \oplus \ker D_+^X \end{aligned}$$

One of the purposes of topological K-theory is to make sense such a \oplus sign, the idea is analogous to the idea of extending natural numbers \mathbb{N} to integers \mathbb{Z} . But for vector bundles we must consider stable equivalent relations rather than equality.

Definition 2.1.1. $E, E' \in \mathcal{V}(M)$, $E \stackrel{s}{\sim} E' \Leftrightarrow \exists F \in \mathcal{V}(M)$, s.t. $E \oplus F \cong E' \oplus F$. E and E' are called stable equivalent. Let $\mathcal{W}(M) = \mathcal{V}(M) / \stackrel{s}{\sim}$. Define a equivalent relation on $\mathcal{V}(M) \times \mathcal{V}(M)$:

$$(E, E') \sim (F, F') \Leftrightarrow E \oplus F' = E' \oplus F \text{ in } \mathcal{W}(M) \Leftrightarrow \exists H \in \mathcal{V}(M), \text{ s.t. } E \oplus F' \oplus H = E' \oplus F \oplus H$$

The K -group of M is defined to be $K(M) = \{(E, E') : E, E' \in \mathcal{V}(M)\} / \sim$.

Remark 6. There are several equivalent definitions of $K(M)$. See ref.

If $E \stackrel{s}{\sim} E'$, then $\text{ch}(E) = \text{ch}(E')$. It's clear the map $\text{ch} : \mathcal{V}(M) \rightarrow H^*(M)$ induces

$$\text{ch} : K(M) \rightarrow H^*(M)$$

Assume again that $\ker D_{\pm}^X$ are smooth vector bundles, (it's sufficient to assume constant dimensions) then

$$\text{Ind}(D_+^X) = \ker D_+^X - \ker D_-^X \in K(S)$$

If the assumption is not verified, we can add a finite dimensional space, and find a surjective homomorphism $\tilde{D}_+ : S_+^{TX} \oplus \mathbb{C}^n \rightarrow S_-^{TX}$, for which the assumption is verified, then

$$\text{Ind}(D_+^X) = \text{Ind}(\tilde{D}_+) - \mathbb{C}^n \in K(S)$$

Theorem 2.1.1 (Family Index Theorem).

$$\text{ch}(\text{Ind}(D_+^X)) = \pi_*[\hat{A}(TX)\text{ch}(E)]$$

in $H^{\text{even}}(S, \mathbb{Q})$, i.e. the diagram commutes:

$$\begin{array}{ccc} K(M) & \xrightarrow[\pi_*]{\text{Ind}(D_+^X)} & K(S) \\ \text{ch} \downarrow & & \downarrow \text{ch} \\ H^{\text{even}}(M, \mathbb{Q}) & \xrightarrow{\pi_*[\hat{A}(TX)\cdot]} & H^{\text{even}}(S, \mathbb{Q}) \end{array}$$

Example 2.1.1. If S is a single point, $\mathcal{V}(S)$ = vector spaces, $K(S) = \mathbb{Z}$, $\text{ch}(E \oplus F) = \text{dim}E - \text{dim}F$. We get the index theorem

$$\text{Ind}(D_+^X) = \int_X \hat{A}(TX)\text{ch}(E)$$

2.2 Consequences of the Family Index Theorem

Assume S is even dimensional, compact and spin and we have a given horizontal space $T^H M$ which is a subbundle of TM . So $T^H M \cong \pi^*TS$, $TM = T^H M \oplus TX \cong \pi^*TS \oplus TX$. Let $g^{TM} = g^{TX} \oplus \pi^*g^{TS}$. Then TM is spin, and $S^{TM} = S^{TX} \hat{\otimes} \pi^*S^{TS}$. Since \hat{A} is multiplicative, we have $\hat{A}(TM) = \hat{A}(TX)\pi^*\hat{A}(TS)$. So

$$\begin{aligned} \text{Ind}(D_+^{M,E}) &= \int_M \hat{A}(TM)\text{ch}(E) = \int_M \pi^*(\hat{A}(TS))\hat{A}(TX)\text{ch}(E) \\ &= \int_S \hat{A}(TS)\pi_*[\hat{A}(TX)\text{ch}(E)] \stackrel{F.I.T}{=} \int_S \hat{A}(TS)\text{ch}(\text{Ind} D_+^X) \\ &= \text{Ind}(D_+^{S, \text{Ind} D_+^X}) \end{aligned} \tag{2.1}$$

Claim 2. (2.1) is equivalent to family index theorem:

Proof. Let F be a vector bundle on S , (2.1) is valid when $E \rightsquigarrow E \otimes \pi^*F$. By the same computation of (2.1)

$$\text{Ind}(D_+^{M, E \otimes \pi^*F}) = \int_S \hat{A}(TS)\text{ch}(F)\pi_*[\hat{A}(TX)\text{ch}(E)]$$

Note that $\text{Ind}(D_+^{X, E \otimes \pi^*F}) = F \otimes \text{Ind}(D_+^{X,E})$, so by A-S index theorem,

$$\text{Ind} D_+^{S, \text{Ind}(D_+^{X, E \otimes \pi^*F})} = \int_S \hat{A}(TS)\text{ch}(F)\text{ch}(\text{Ind} D_+^{X,E})$$

Since $\text{ch}(F)$ generate the full $H^{\text{even}}(S, \mathbb{Q})$ as F varies, compare the above two equation, we have

$$\hat{A}(TS)\pi_*[\hat{A}(TX)\text{ch}(E)] = \hat{A}(TS)\text{ch}(\text{Ind} D_+^{X,E})$$

Since $\hat{A} = 1 + \dots$ is invertible in $H^{\text{even}}(S, \mathbb{Q})$, so

$$\pi_*[\hat{A}(TX)\text{ch}(E)] = \text{ch}(\text{Ind} D_+^{X,E})$$

□

Remark 7. Atiyah-Singer proved the family index theorem by proving (2.1).

2.3 Adiabatic Limit

From now on, we assume $g_\varepsilon = g_\varepsilon^{TM} = g^{TX} \oplus \frac{1}{2}\pi^*g^{TS}$, then $T^H M = (TX)^\perp \cong \pi^*TS$, and $\pi : M \rightarrow S$ is a Riemannian submersion. Let $\nabla^{TM,\varepsilon}$ denote the Levi-Civita connection of g_ε^{TM} , $\nabla^{TM,L} = \nabla^{TM,1}$ denote the Levi-Civita connection of $g^{TM} = g^{TX} \oplus \pi^*g^{TS} = g_1^{TM}$.

We want to calculate the Dirac operator D_ε^M using the metric g_ε^{TM} .

Theorem 2.3.1. *As $\varepsilon \rightarrow 0$, $\nabla^{TM,\varepsilon} = \begin{pmatrix} \nabla^{TX} & \\ 0 & \pi^*\nabla^{TS} \end{pmatrix} + O(\varepsilon) =: \nabla^{TM,0} + O(\varepsilon)$.*

Proof. $\forall U \in TS$, U^H denotes the lift of U to $T^H M$. If U is a smooth vector field on S , then U^H is a smooth vector field on M . Let φ_t, ψ_t be one-parameter transformation group generated by U and U^H , then $\varphi_t \circ \pi = \pi \circ \psi_t$, and $\psi_t : X_s \rightarrow X_{\varphi_t(s)}$ is a diffeomorphism. So if V is a smooth section of TX , then $[U^H, V] \in TX$.

By properties of connection, we can assume X is either vertical or $X = U^H \in T^H M$ in the following proof. We have defining equation of the Levi-Civita connection:

$$2\langle \nabla_X Y, Z \rangle = X\langle Y, Z \rangle + Y\langle Z, X \rangle - Z\langle X, Y \rangle + \langle Z, [X, Y] \rangle - \langle Y, [X, Z] \rangle - \langle X, [Y, Z] \rangle \quad (2.2)$$

Using this formula, we calculate all possible cases:

1. Y, Z are vertical, X is vertical

$$\langle \nabla_X^{TM,\varepsilon} Y, Z \rangle_{g^{TX}} = \langle \nabla_X^{TX} Y, Z \rangle_{g^{TX}}$$

2. Y, Z are vertical, $X = U^H$ is horizontal

$$\begin{aligned} 2\langle \nabla_{U^H}^{TM,\varepsilon} Y, Z \rangle_{g^{TX}} &= U^H\langle Y, Z \rangle_{g^{TX}} + \langle Z, [U^H, Y] \rangle_{g^{TX}} - \langle Y, [U^H, Z] \rangle_{g^{TX}} \\ &= 2\langle [U^H, Y], Z \rangle + (L_{U^H} g^{TX})(Y, Z) \\ &= 2\langle [U^H, Y], Z \rangle + \langle (g^{TX})^{-1} L_{U^H} g^{TX}(Y), Z \rangle \end{aligned}$$

$$\text{So } P^{TX} \nabla_{U^H}^{TM,\varepsilon} Y = [U^H, Y] + \frac{1}{2} \langle (g^{TX})^{-1} L_{U^H} g^{TX}(Y), Z \rangle Z.$$

3. $Y = V^H$ is horizontal, Z is vertical, X is vertical

$$2\langle \nabla_X^{TM,\varepsilon} V^H, Z \rangle_{g^{TX}} = V^H\langle X, Z \rangle_{g^{TX}} - \langle [V^H, X], Z \rangle_{g^{TX}} - \langle X, [V^H, Z] \rangle_{g^{TX}} = (L_{V^H} g^{TX})(X, Z)$$

$$\text{So } P^{TX} \nabla_X^{TM,\varepsilon} V^H = \frac{1}{2} \langle (g^{TX})^{-1} (L_{V^H} g^{TX})(X), Z \rangle Z.$$

4. $Y = V^H$ is horizontal, Z is vertical, $X = U^H$ is horizontal

$$2\langle \nabla_{U^H}^{TM,\varepsilon} V^H, Z \rangle_{g^{TX}} = \langle P^{TX}[U^H, V^H], Z \rangle_{g^{TX}}$$

$$\text{So } P^{TX} \nabla_{U^H} V^H = \frac{1}{2} P^{TX}[U^H, V^H].$$

5. Y is vertical, $Z = W^H$ is horizontal, X is vertical

$$2\langle \nabla_X^{TM,\varepsilon} Y, W^H \rangle_{\pi^*g^{TS}} = -\varepsilon \langle L_{W^H} g^{TX}(X), Y \rangle$$

$$\text{So } P^{T^H M} \nabla_X^{TM,\varepsilon} Y = O(\varepsilon).$$

6. Y is vertical, $Z = W^H$ is horizontal, $X = U^H$ is horizontal

$$2\langle \nabla_{U^H}^{TM,\varepsilon} Y, W^H \rangle_{\pi^*g^{TS}} = -\varepsilon \langle Y, P^{TX}[U^H, W^H] \rangle_{g^{TX}}$$

$$\text{So } P^{T^H M} \nabla_{U^H}^{TM,\varepsilon} Y = O(\varepsilon)$$

7. $Y = V^H$, $Z = W^H$ are horizontal, X is vertical

$$2\langle \nabla_X^{TM,\varepsilon} V^H, W^H \rangle_{\pi^*g^{TS}} = -\varepsilon \langle X, P^{TX}[V^H, W^H] \rangle_{g^{TX}}$$

So $P^{T^H M} \nabla_X^{TM,\varepsilon} V^H = O(\varepsilon)$.

8. $Y = V^H$, $Z = W^H$ are horizontal, $X = U^H$ is horizontal

$$\langle \nabla_{U^H}^{TM,\varepsilon} V^H, W^H \rangle_{\pi^*g^{TS}} = \langle \nabla_U^{TS} V, W \rangle_{g^{TS}}$$

So $P^{T^H M} \nabla_{U^H}^{TM,\varepsilon} V^H = (\nabla_U^{TS} V)^H$.

Write in matrix,

$$\nabla^{TM,\varepsilon} = \begin{pmatrix} P^{TX} \nabla^{TM,\varepsilon} P^{TX} & P^{TX} \nabla^{TM,\varepsilon} P^{T^H M} \\ P^{T^H M} \nabla^{TM,\varepsilon} P^{TX} & P^{T^H M} \nabla^{TM,\varepsilon} P^{T^H M} \end{pmatrix} = \begin{pmatrix} \nabla^{TX} & * \\ O(\varepsilon) & \pi^* \nabla^{TS} + O(\varepsilon) \end{pmatrix}$$

where $*$ = $P^{TX} \nabla^{TM,\varepsilon} P^{T^H M}$ is determined by above calculations in cases 3 and 4. The theorem is proved. \square

From above calculations in cases 1 and 2, we see that

Theorem 2.3.2. *The connection $\nabla^{TX} = P^{TX} \nabla^{TM,\varepsilon} P^{TX}$ is characterized by the following two properties:*

1. ∇^{TX} restricts to the Levi-Civita connection of (X, g^{TX}) along the fibre X .

2. $\forall U \in TS$, $Y \in TX$, $\nabla_{U^H}^{TX} Y = [U^H, Y] + \frac{1}{2}((g^{TX})^{-1} L_{U^H} g^{TX}) Y$.

Remark 8. ∇^{TX} preserves the metric g^{TX} . In general, if ∇^E is any connection on an Hermitian vector bundle E , we can modify it to be metric preserving: $\nabla'^E = \nabla^E + \frac{1}{2}(g^E)^{-1}(\nabla^E g^E)$.

$$\begin{aligned} X \langle s, s' \rangle &= (\nabla_X^E g^E)(s, s') + \langle \nabla_X^E s, s' \rangle + \langle s, \nabla_X^E s' \rangle \\ &= \langle \frac{1}{2}(g^E)^{-1}(\nabla_X^E g^E) s, s' \rangle + \langle s, \frac{1}{2}(g^E)^{-1}(\nabla_X^E g^E) s' \rangle + \langle \nabla_X^E s, s' \rangle + \langle s, \nabla_X^E s' \rangle \\ &= \langle \nabla_X'^E s, s' \rangle + \langle s, \nabla_X'^E s' \rangle \end{aligned}$$

Let $\nabla^\oplus = \nabla^{TX} \oplus \pi^* \nabla^{TS}$ denote the splitting connection, we write

$$S^\varepsilon = \nabla^{TM,\varepsilon} - \nabla^\oplus, S = \nabla^{TM,L} - \nabla^\oplus$$

Proposition 2.3.1. *Assume $\nabla^{TM,L}$ is the Levi-Civita connection of g^{TM} , and ∇ is a metric preserving connection with torsion T . Let $\nabla^{TM,L} = \nabla + S$, then*

$$2\langle S(X)Y, Z \rangle = \langle T(X, Z), Y \rangle + \langle T(Y, Z), X \rangle - \langle T(X, Y), Z \rangle \quad (2.3)$$

We first evaluate the torsion T of ∇^\oplus .

Proposition 2.3.2. 1. T takes value in TX ;

2. T vanishes on $TX \times TX$;

3. $T(U^H, V^H) = -P^{TX}[U^H, V^H]$;

4. $T(U^H, A) = \frac{1}{2}(g^{TX})^{-1}(L_{U^H} g^{TX})A$, $U \in TS$, $A \in TX$.

The proof is straightforward computation. By above two proposition, we have the properties of S and S^ε :

Proposition 2.3.3. 1. $P^{TX}S^\varepsilon = P^{TX}S$

2. $P^{T^H M}S^\varepsilon = \varepsilon P^{T^H M}S$

3. $S(\cdot)$ maps TX into $T^H M$

4. $\langle S(U^H)V^H, W^H \rangle = 0$, for $U, V, W \in TS$

Proof. 1. by (2.3)

2. by (2.3) and properties 3,4 in proposition 2.3.2

3. because $\nabla^{TX} = P^{TX}\nabla^{TM,L}P^{TX} = P^{TX}\nabla^\oplus P^{TX}$

4. by (2.3)

□

Remark 9. The proposition also follows from the proof of theorem (2.3.1), and the formula

$$S^\varepsilon = \begin{pmatrix} 0 & * \\ O(\varepsilon) & O(\varepsilon) \end{pmatrix}$$

Theorem 2.3.3.

$$D_\varepsilon^M = D^X + \sqrt{\varepsilon}D^H - \frac{\varepsilon c(T^H)}{4} \quad (2.4)$$

- D^X is the Dirac operator along the fibre: $D^X = \sum_i c(e_i)\nabla_{e_i}^{S^{TX} \otimes E}$, $\{e_i\}$ is a orthonormal basis of TX .
- T^H is the HH component of T : $T^H(U^H, V^H) = -P^{TX}[U^H, V^H]$, and

$$c(T^H) = \frac{1}{2}\langle T^H(f_\alpha^H, f_\beta^H), e_i \rangle c(f_\alpha)c(f_\beta)c(e_i)$$

$\{f_\alpha\}$ is a orthonormal basis of TS .

- $D^H = c(f_\alpha)(\nabla_{f_\alpha^H}^{S^{TX} \otimes S^{TS} \otimes E} + k(f_\alpha))$, $k(U) \triangleq \frac{1}{2}div^X(U^H) = \frac{1}{4}Tr[(g^{TX})^{-1}L_{U^H}g^{TX}]$

Proof.

$$D_\varepsilon^M = c(e_i)\nabla_{e_i}^\varepsilon + c(f_\alpha)\nabla_{\sqrt{\varepsilon}f_\alpha^H}^\varepsilon \quad (2.5)$$

∇^ε is the induced connection on the spinor bundle $S^{TM} = S^{TX} \hat{\otimes} \pi^* S^{TS}$ by the Levi-Civita connection $\nabla^{TM, \varepsilon}$. Note that the Clifford multiplication doesn't change when rescaling the horizontal metric. We have $\nabla^{TM, \varepsilon} = \nabla^{TM} + S^\varepsilon$, and Lie algebra homomorphism:

$$\begin{aligned} so(n) &\longrightarrow spin(n) \subset Cl_n \\ A &\mapsto c(A) = \frac{1}{4}\langle Ae_i, e_j \rangle c(e_i)c(e_j) \end{aligned}$$

so we get

$$\nabla^\varepsilon = \nabla^{S^{TX}} + \nabla^{S^{TS}} + \frac{1}{4}\langle S^\varepsilon(\cdot)e_i, e_j \rangle_{g_\varepsilon} c(e_i)c(e_j) + \frac{1}{4}\langle S^\varepsilon(\cdot)\sqrt{\varepsilon}f_\alpha^H, \sqrt{\varepsilon}f_\beta^H \rangle_{g_\varepsilon} + \frac{1}{2}\langle S^\varepsilon(\cdot)\sqrt{\varepsilon}f_\alpha^H, e_i \rangle_{g_\varepsilon} c(f_\alpha)c(e_i) \quad (2.6)$$

$\nabla^{S^{TX}}$ and $\nabla^{S^{TS}}$ is the connection on spinors induced by ∇^{TX} and ∇^{TS} . By (2.3), proposition 2.3.3, and proposition 2.3.2

- $\langle S^\varepsilon(\cdot)e_i, e_j \rangle_{g_\varepsilon} = 0$
- $\langle S^\varepsilon(\cdot)\sqrt{\varepsilon}f_\alpha^H, \sqrt{\varepsilon}f_\beta^H \rangle_{g_\varepsilon} = \varepsilon \langle S(\cdot)f_\alpha^H, f_\beta^H \rangle$
 $c(e_i) \cdot \frac{1}{4} \langle S(e_i)f_\alpha^H, f_\beta^H \rangle c(f_\alpha)c(f_\beta) = \frac{1}{8} \langle T(f_\alpha^H, f_\beta^H), e_i \rangle c(f_\alpha)c(f_\beta)c(e_i)$
 $\langle S(f_\alpha^H)f_\alpha^H, f_\beta^H \rangle = 0.$
- $\langle S^\varepsilon(\cdot)\sqrt{\varepsilon}f_\alpha^H, e_i \rangle_{g_\varepsilon} = \sqrt{\varepsilon} \langle S(\cdot)f_\alpha^H, e_i \rangle$
 $c(f_\alpha) \cdot \frac{1}{2} \langle S(f_\alpha^H)f_\beta^H, e_i \rangle c(f_\beta)c(e_i) = -\frac{1}{4} \langle T(f_\alpha^H, f_\beta^H), e_i \rangle c(f_\alpha)c(f_\beta)c(e_i)$
 $c(e_i) \cdot \frac{1}{2} \langle S(e_i)f_\alpha^H, e_j \rangle c(f_\alpha)c(e_j) = \frac{1}{4} (\langle T(f_\alpha^H, e_j), e_i \rangle - \langle T(e_i, f_\alpha^H), e_j \rangle) c(e_i)c(f_\alpha)c(e_j)$
 $= \sum_i \frac{1}{2} \langle T(f_\alpha^H, e_i), e_i \rangle c(f_\alpha) + \sum_{i \neq j} \dots$
 $= \frac{1}{2} \langle T(f_\alpha^H, e_i), e_i \rangle c(f_\alpha)$

$$\langle T(f_\alpha^H, e_i), e_i \rangle = \langle \frac{1}{2}(g^{TX})^{-1}(L_{f_\alpha^H}g^{TX})e_i, e_i \rangle = \frac{1}{2} \text{Tr}[(g^{TX})^{-1}L_{f_\alpha^H}g^{TX}]$$

Lemma 2. Let dV_X be the volume element along the fibre, then

$$\text{div}^X(U^H) = \frac{L_{U^H}dV_X}{dV_X} = \frac{1}{2} \text{Tr}[(g^{TX})^{-1}L_{U^H}g^{TX}]$$

proof of the lemma. $\{e_i\}$ is an orthonormal basis of TX , $\{e^i\}$ is the dual basis. We calculate

$$\begin{aligned} L_{U^H}dV_X &= L_{U^H}(e^1 \wedge \dots \wedge e^n) = \sum_i e^1 \wedge \dots \wedge L_{U^H}e^i \wedge \dots \wedge e^n \\ &= \sum_i e^1 \wedge \dots \wedge \left(\sum_j -\langle L_{U^H}e_j, e_i \rangle e^j \right) \wedge \dots \wedge e^n \\ &= -\left(\sum_i \langle L_{U^H}e_i, e_i \rangle \right) dV_X \end{aligned}$$

$$\begin{aligned} \text{Tr}[(g^{TX})^{-1}L_{U^H}g^{TX}] &= \text{Tr}_{g^{TX}}(L_{U^H}g^{TX}) = \sum_i (L_{U^H}g^{TX})(e_i, e_i) \\ &= \sum_i (U^H \langle e_i, e_i \rangle - \langle L_{U^H}e_i, e_i \rangle - \langle e_i, L_{U^H}e_i \rangle) \\ &= -2 \sum_i \langle L_{U^H}e_i, e_i \rangle \end{aligned}$$

□

Put the results together, finally we have

$$\begin{aligned} D_\varepsilon^M &= c(e_i)\nabla_{e_i}^{TX \otimes E} + \sqrt{\varepsilon}c(f_\alpha)(\nabla_{f_\alpha^H}^{S^{TX} \otimes S^{TS} \otimes E} + k(U)) - \frac{\varepsilon}{8} \langle T(f_\alpha^H, f_\beta^H), e_i \rangle c(f_\alpha)c(f_\beta)c(e_i) \\ &= D^X + \sqrt{\varepsilon}D^H - \frac{\varepsilon c(T^H)}{4} \end{aligned}$$

□

2.4 Levi-Civita Superconnection

Definition 2.4.1. $\forall s \in S, H_s = \Gamma(X_s, (S^{TX} \otimes E)|_{X_s})$. $H_s = H_{s,+} \oplus H_{s,-}$.

H is an infinite dimensional \mathbb{Z}_2 -graded vector bundle with L^2 Hermitian metric .

Definition 2.4.2. $\forall U \in TS$, f is a smooth section of H ,

$$\nabla_U^H f = \nabla_{U^H}^{S^{TX} \otimes E} f + \frac{1}{2} \text{div}^X(U^H) f =: \nabla_U'^H f + k(U) f$$

Note that $\Gamma(S, H) = \Gamma(S, \Gamma(X, S^{TX} \otimes E|_X)) = \Gamma(M, S^{TX} \otimes E)$

Proposition 2.4.1. ∇^H is a unitary connection H , whose curvature is a two form on the base with values in 1st order differential operator along the fibre.

Proof. Let $U, U^H, \varphi_t, \psi_t$ be the same as in the beginning of proof of theorem (2.3.1). $\forall f, f' \in \Gamma(H)$,

$$\begin{aligned} U \langle f, f' \rangle_H &= U \int_t \langle f, f' \rangle dV_X = \frac{d}{dt} \int_{\varphi_t(X)} \langle f, f' \rangle dV_{\varphi_t(X)} \\ &= \int_X \frac{d}{dt} \psi_t^* (\langle f, f' \rangle dV_{\varphi_t(X)}) = \int_X (U^H \langle f, f' \rangle + \frac{L_{U^H} dV_X}{dV_X}) dV_X \\ &= \int_X (\langle \nabla_{U^H} f, f' \rangle + \langle f, \nabla_{U^H} f' \rangle + \text{div}^X(U^H) \langle f, f' \rangle) dV_X \\ &= \langle \nabla_U^H f, f' \rangle_H + \langle f, \nabla_U^H f' \rangle \end{aligned}$$

So ∇^H is unitary. Since $\nabla^H = \nabla'^H + k$, $\nabla^{H,2} = \nabla'^{H,2} + dk$. So $\forall U, V \in \Gamma(TS)$,

$$\begin{aligned} R^H(U, V) &= \nabla_{U^H}^{S^{TX} \otimes E} \nabla_{V^H}^{S^{TX} \otimes E} - \nabla_{V^H}^{S^{TX} \otimes E} \nabla_{U^H}^{S^{TX} \otimes E} - \nabla_{[U, V]^H}^{S^{TX} \otimes E} + dk(U, V) \\ &= R^{S^{TX}}(U^H, V^H) + R^E(U^H, V^H) - \nabla_{P^{TX}[U^H, V^H]}^{S^{TX} \otimes E} + dk(U, V) \end{aligned}$$

So R^H is a two form on the base S with value in 1st order differential operator along the fibre. \square

By this proposition and $\Gamma(M, S^{TX} \hat{\otimes} S^{TS} \otimes E) = \Gamma(S, S^{TS} \hat{\otimes} H)$, we see D^H is just the Dirac operator on S with twisting bundle $(H, \langle \cdot, \cdot \rangle_{L^2}, \nabla^H)$.

Definition 2.4.3. The Levi-Civita superconnection on H is given by

$$A = D^X + \nabla^H - \frac{c(T^H)}{4}$$

$$c(T^H) = \frac{1}{2} \langle T^H(f_\alpha^H, f_\beta^H), e_i \rangle f^\alpha f^\beta c(e_i).$$

Remark 10. Formally, this is obtained by changing Clifford variable $c(f_\alpha)$ in D^M to Grassman variable f^α . In fact, it's obtained by Getzler rescaling: in the expression of D_ε^M , we substitute:

$$c(f_\alpha) \rightsquigarrow \frac{f^\alpha \wedge}{\varepsilon} - \varepsilon i_{f_\alpha}$$

and let $\varepsilon \rightarrow 0$, then we get the Levi-Civita superconnection A .

We want to calculate the curvature of A . We use the Lichnerowicz formula for D_ε^M :

$$D_\varepsilon^{M,2} = -\Delta_\varepsilon^M + \frac{K_\varepsilon^M}{4} + \mathcal{R}^E \quad (2.7)$$

where $\Delta_\varepsilon^M = (\nabla_{e_i}^\varepsilon)^2 + \varepsilon (\nabla_{f_\alpha}^\varepsilon)^2$ is the Bochner Laplacian, ∇^ε is given by equation (2.6). K_ε^M is the scalar curvature of g_ε^{TM} , and

$$\mathcal{R}^E = \frac{1}{2} c(e_i) c(e_j) R^E(e_i, e_j) + \frac{\sqrt{\varepsilon}}{2} c(e_i) c(f_\alpha) R^E(e_i, f_\alpha^H) + \frac{\varepsilon}{2} c(f_\alpha) c(f_\beta) R^E(f_\alpha^H, f_\beta^H)$$

The idea is to change Clifford variable $c(f_\alpha)$ to f^α , note that

- $(\sqrt{\varepsilon} c(f_\alpha))^2 = -\varepsilon \rightsquigarrow (f^\alpha)^2 = f^\alpha \wedge f^\alpha = 0$

- $\alpha \neq \beta$, $\sqrt{\varepsilon}c(f_\alpha)\sqrt{\varepsilon}c(f_\beta) = \varepsilon c(f_\alpha)c(f_\beta) \rightsquigarrow f^\alpha \wedge f^\beta$

So if the power of $\sqrt{\varepsilon} > (=)$ the length of Clifford variables, then the term is killed(survives).
 $\nabla^{TM,\varepsilon} = \nabla^{TM} + S^\varepsilon$, $\nabla^{TM,\varepsilon,2} = \nabla^{TM,2} + \nabla^{TM}S^\varepsilon + [S^\varepsilon, S^\varepsilon]$.

Definition 2.4.4.

$$\alpha_t = \varphi Tr_s(\exp(-A_t^2))$$

Theorem 2.4.1. 1. The α_t are real, even, closed forms on S .

$$[\alpha_t] = ch(Ind(D_+^X))$$

2. As $t \rightarrow 0$, $\alpha_t = \pi_*[\hat{A}(\nabla^{TX})ch(\nabla^E)] + O(\sqrt{t})$

3. If $dimker D_\pm^X$ is locally constant, as $t \rightarrow \infty$,

$$\alpha_t = ch(\nabla^{ker D^X}) + O\left(\frac{1}{\sqrt{t}}\right)$$

$$\nabla^{ker D^X} = P^{ker D^X} \nabla^H.$$

Remark 11. 1. The theorem implies the family index theorem:

$$ch(Ind(D_+^X)) = [\alpha_t] = \pi_*[\hat{A}(TX)ch(E)]$$

2. the theorem asserts the following diagram commutes.

Chapter 3

Determinant Bundle

3.1 Finite Dimensional Case

Let \mathcal{C} denote the category of complex line (complex vector space of dim 1) up to canonical isomorphism. For example

- \mathbb{C} denotes the canonical line. If λ has a canonical nonzero element $s \in \lambda$, then λ can be canonically identified with \mathbb{C} : $a \in \mathbb{C} \mapsto as \in \lambda$.
- If λ, μ are complex lines, then canonically $\lambda \otimes \mu \cong \mu \otimes \lambda$, $a \otimes b \mapsto b \otimes a$. The operator \otimes is commutative, associative and has a neutral element:

$$- \lambda \otimes \mathbb{C} \cong \lambda: s \otimes 1 \mapsto s.$$

$$- \lambda^* \otimes \lambda \cong \mathbb{C}: \text{choose any } s \neq 0 \in \lambda, \text{ there is } s^{-1} \in \lambda^*, \text{ s.t. } \langle s^{-1}, s \rangle = 1. s^{-1} \otimes s \text{ does not depend on } s, \text{ so it's a canonical nonzero element in } \lambda^* \otimes \lambda. \text{ We define } \lambda^{-1} = \lambda^*.$$

Definition 3.1.1. E is a finite dimensional complex vector space

$$\det E \triangleq \wedge^{\max} E$$

$\wedge^{\max} E$ means the elements of max degree in $\wedge^* E$.

Remark 12. More correctly, $\mathcal{C} = \{(\lambda, \pm)\}$, $\mathcal{V} = \{\text{finite dimensional complex vector space}\}$.

$$\widehat{\det} : \mathcal{V} \longrightarrow \mathcal{C}, \quad \widehat{\det} E = (\det E, (-1)^{\dim E})$$

then $\widehat{\det}(E \oplus F) = \widehat{\det} E \widehat{\otimes} \widehat{\det} F$, where $\lambda \widehat{\otimes} \mu \cong \mu \widehat{\otimes} \lambda$ is given by $a \widehat{\otimes} b \mapsto (-1)^{\varepsilon_\lambda \varepsilon_\mu} b \widehat{\otimes} a$.

Let $E : 0 \rightarrow E_0 \xrightarrow{V} E_1 \rightarrow 0$, $\lambda = \det E \triangleq (\det E_0)^{-1} \otimes \det E_1$. V induces $\det V : \det E_0 \rightarrow \det E_1$, so $\det V \in \lambda$.

Let E be a \mathbb{Z}_2 -graded vector bundle with Hermitian metric $g^E = g^{E_+} \oplus g^{E_-}$, unitary connection $\nabla^E = \nabla^{E_+} \oplus \nabla^{E_-}$. We have the induced metric g^λ and unitary connection ∇^λ . Assume $V = \begin{pmatrix} 0 & V_- \\ V_+ & 0 \end{pmatrix}$ is self-adjoint and V_+ is invertible. Let $\{e_i\}$ be an orthonormal basis of E_+

$$\begin{aligned} (\nabla_X^\lambda (\det V_+))(e_1 \wedge \cdots \wedge e_n) &= \nabla_X (V_+(e_1) \wedge \cdots \wedge (e_n)) - \det V_+ \left(\sum_i e_1 \wedge \cdots \wedge \nabla_X^E e_i \wedge \cdots \wedge e_n \right) \\ &= \sum_i V_+(e_1) \wedge \cdots \wedge (\nabla_X^{E_-} (V_+(e_i)) - V_+(\nabla_X^{E_+} e_i)) \wedge \cdots \wedge V_+(e_n) \\ &= \sum_i V_+(e_1) \wedge \cdots \wedge V_+(V_+^{-1} \nabla_X V_+)(e_i) \wedge \cdots \wedge V_+(e_n) \\ &= \text{Tr}(V_+^{-1} \nabla_X V_+) \cdot \det V_+(e_1 \wedge \cdots \wedge e_n) \end{aligned}$$

So in anticommutative sign rule

$$\nabla^\lambda(\det V_+) = -(\det V_+) \cdot \text{Tr}(V_+^{-1} \nabla^E V_+) \quad (3.1)$$

Note that

$$V^{-1} \nabla^E V = \begin{pmatrix} 0 & V_+^{-1} \\ V_-^{-1} & 0 \end{pmatrix} \begin{pmatrix} 0 & \nabla^E V_- \\ \nabla^E V_+ & 0 \end{pmatrix} = \begin{pmatrix} V_+^{-1} \nabla^E V_+ & 0 \\ 0 & V_-^{-1} \nabla^E V_- \end{pmatrix}$$

So

$$\text{Tr}(V_+^{-1} \nabla^E V_+) = \frac{1}{2} \text{Tr}(V^{-1} \nabla^E V) + \frac{1}{2} \text{Tr}_s(V_-^{-1} \nabla^E V) \quad (3.2)$$

Since $(V_+^{-1} \nabla_X^E V_+)^* = (\nabla_X^E V_-) \cdot V_-^{-1}$, $\frac{1}{2} \text{Tr}(V^{-1} \nabla_X^E V)$ is real, $\frac{1}{2} \text{Tr}_s(V^{-1} \nabla_X^E V)$ is purely imaginary. Moreover, $\text{Tr}(V^{-1} \nabla^E V)$ is exact: $\det V \in (\det E)^* \otimes \det E \cong \underline{\mathbb{C}}$ is a well defined function, so it's easy to see

$$\text{Tr}(V^{-1} \nabla^E V) = \frac{d(\det V)}{\det V} = d \log |\det V| = 2d \log |\det V_+| \quad (3.3)$$

Let r^λ be the curvature of ∇^λ , by formula (3.1),(3.2),(3.3)

$$r^\lambda = d \frac{\nabla^\lambda \det V_+}{\det V_+} = -d \text{Tr}(V_+^{-1} \nabla^E V_+) = -\frac{1}{2} d \text{Tr}_s(V^{-1} \nabla^E V) \quad (3.4)$$

Lemma 3. $c_1(\det E) = ch(E)^{(2)} = c_1(E)$

in definition 1.4.1, we defined

$$\beta = \frac{1}{\sqrt{2\pi i}} \int_0^\infty \varphi \text{Tr}_s \left(\frac{V}{2\sqrt{t}} \exp(-A_t^2) \right) dt$$

where $A_t^2 = \nabla^E \cdot 2 + \sqrt{t} \nabla^E V + tV^2$.

Note that $\frac{(-1)^n}{n!} \text{Tr}_s(V(\sqrt{t} \nabla^E V + tV^2)^n)^{(1)} = \frac{(-1)^n}{(n-1)!} \sqrt{t} \text{Tr}_s(V \nabla^E V (tV^2)^{n-1})$,

$$\text{Tr}_s \left(\frac{V}{2\sqrt{t}} \exp(-A_t^2) \right)^{(1)} = -\frac{1}{2} \text{Tr}_s(V(\nabla^E V) \exp(-tV^2))$$

$$\beta^{(1)} = -\frac{1}{2\pi i} \cdot \frac{1}{2} \int_0^\infty \text{Tr}_s(V \nabla^E V \exp(-tV^2)) dt = \frac{\sqrt{-1}}{2\pi} \cdot \frac{1}{2} \text{Tr}_s(V^{-1} \nabla^E V)$$

By formula (3.4)

$$d\beta^{(1)} = -\frac{\sqrt{-1}}{2\pi} r^\lambda = -c_1(\det E) = -ch(E)^{(2)} \quad (3.5)$$

3.2 Determinant Bundle

We want to construct the determinant bundle λ of $\{H = H_+ \oplus H_-, D_+^X\}$. If $\ker D_\pm^X$ has constant dimension, $\lambda \triangleq (\det \ker D_+^X)^{-1} \otimes (\det \ker D_-^X)$. In general cases, $\forall a > 0$,

$$U_a \triangleq \{s \in S \mid a \notin Sp(D_s^{X,2})\}$$

U_a is an open set on S . $\bigcup_{a>0} U_a$ is an open covering of S .

Definition 3.2.1. On U_a , $H^{<a} = \bigoplus_{\substack{\lambda \in Sp(D^{X,2}) \\ \lambda < a}} (\text{bundle of eigenvalue } \lambda \text{ of } D^{X,2})$.

We use other similar notations, like $H^{(0,a)}$, $D^{(0,a)} = D^X|_{H^{(0,a)}}$, e.t..

Proposition 3.2.1. $H^{<a}$ is a finite dimensional smooth vector bundle on U_a . $H^{<a} = H_+^{<a} \oplus H_-^{<a}$.

Definition 3.2.2. On U_a , $\lambda^{<a} = \det H^{<a} = (\det H_+^{<a})^{-1} \otimes \det H_-^{<a}$. $\lambda^{<a}$ is a smooth line bundle on U_a .

The idea is that we can't define $\det H = (\det H_+)^{-1} \otimes \det H_-$ directly, but we can patch $\lambda^{<a}$ together to get a well defined line bundle.

Proposition 3.2.2. $\forall s \in U_a$, $\lambda_s^{<a} \cong (\det \ker D_+^{X_s})^{-1} \otimes (\det \ker D_-^{X_s}) =: \det \ker D^{X_s}$. This is a canonical isomorphism.

Proof. We have an exact sequence:

$$\begin{array}{ccccccc} 0 & \longrightarrow & \ker D_+^{X_s} & \longrightarrow & H_+^{<a} & \xrightarrow{D_+} & H_-^{<a} & \longrightarrow & \ker D_-^{X_s} & \longrightarrow & 0 \\ & & & & \parallel & & \parallel & & & & \\ & & & & \ker D_+^{X_s} \oplus H_+^{(0,a)} & \xrightarrow{D_+} & H_-^{(0,a)} \oplus \ker D_-^{X_s} & & & & \end{array}$$

Let $\sigma_{\pm} \neq 0 \in \det \ker D_{\pm}^{X_s}$, $\tau \neq 0 \in \det H_+^{(0,a)}$, define

$$\begin{aligned} \det \ker D^X &\xrightarrow{\cong} \det H^{<a} \\ \sigma_+^{-1} \otimes \sigma_- &\mapsto (\sigma_+ \wedge \tau)^{-1} \otimes (\det D_+^X \tau \wedge \sigma_-) \end{aligned}$$

This isomorphism is independent of τ . In fact, $\det D_+^{(0,a)} = \tau^{-1} \otimes (\det D_+^X \tau)$ is a nonzero canonical section of $\det H^{(0,a)}$. So $\sigma_+^{-1} \otimes \sigma_- \mapsto (\sigma_+^{-1} \otimes \sigma_-) \cdot (\det D_+^{(0,a)})$ is a canonical isomorphism from $\det \ker D^X$ to $\det H^{<a} = (\det \ker D^X) \otimes (\det H^{(0,a)})$. \square

$\forall 0 < a < b$, $a, b \notin Sp(D^{X,2})$, $H^{<b} = H^{<a} \oplus H^{(a,b)}$, $\det H^{<b} = \det H^{<a} \otimes \det H^{(a,b)}$. $\det D_+^{(a,b)}$ is a canonical nonzero section of $\det H^{(a,b)}$, so we have a canonical isomorphism

$$\begin{aligned} \varphi_b^a : \lambda^{<a} &\longrightarrow \lambda^{<b} \\ \sigma &\mapsto \sigma \otimes (\det D_+^{(a,b)}) \end{aligned}$$

Since $D_+^{(a,c)} = D_+^{(a,b)} \oplus D_+^{(b,c)}$, $\det D_+^{(a,c)} = \det D_+^{(a,b)} \otimes \det D_+^{(b,c)} \in \det H^{(a,b)} \otimes \det H^{(b,c)} = \det H^{(a,c)}$. So $\varphi_c^a = \varphi_c^b \circ \varphi_b^a : \lambda^{<a} \rightarrow \lambda^{<c}$. We can define the determinant line bundle:

Definition 3.2.3. The determinant line bundle λ is the complex line bundle obtained by pasting $(\lambda^{<a}, U_a)$ together via the canonical isomorphisms $\{\varphi_b^a : \lambda^{<a} \rightarrow \lambda^{<b}\}$

Remark 13. If $\ker D_{\pm}^X$ have constant dimension, then

$$\lambda \cong \det \ker D^X$$

Proposition 3.2.3. $c_1(\lambda) = \pi_*[\hat{A}(TX)ch(E)]^{(2)}$ in $H^2(S)$.

Proof. By the family index theorem, $ch(Ind D_+^X) = \pi_*[\hat{A}(TX)ch(E)]$ in $H^2(S)$. So $c_1(\lambda) = c_1(\det \ker D^X) = c_1(\det Ind D_+^X) = ch(Ind D_+^X)^{(2)} = \pi_*[\hat{A}(TX)ch(E)]^{(2)}$. \square

We want to construct a metric on λ . First we have a induced metric $|\cdot|_{\lambda^{<a}}$ on $\lambda^{<a}$ by the metric of H . $\forall \sigma \in \lambda^{<a}$,

$$|\varphi_b^a(\sigma)|_{\lambda^{<b}} = |\sigma \otimes (\det D_+^{(a,b)})|_{\lambda^{<b}} = |\sigma|_{\lambda^{<a}} |\det(D_-^{(a,b)} D_+^{(a,b)})|_{H_+^{(a,b)}}^{\frac{1}{2}}$$

We want to define $\|\cdot\|_{\lambda^{<a}} = |\cdot|_{\lambda^{<a}} |\det D_+^{(a,\infty)}|$, then

$$\|\varphi_b^a(\sigma)\|_{\lambda^{<b}} = |\varphi_b^a(\sigma)|_{\lambda^{<b}} |\det D_+^{(b,\infty)}| = |\sigma|_{\lambda^{<a}} |\det D_+^{(a,b)}| |\det D_+^{(b,\infty)}| = |\sigma|_{\lambda^{<a}} |\det D^{(a,\infty)}| = \|\sigma\|_{\lambda^{<a}}$$

so the metric patch together. We have to make sense the meaning of $|\det D_+^{(a,\infty)}|$ in this infinite dimensional case.

$L = D_-^X D_+^X$ is a second order, elliptic, positive differential operator, we define

$$\zeta(s) = \text{Tr}^*(L^{-s}) = \sum_{\lambda \in \text{Sp}^*(D_-^X D_+^X)} \frac{1}{\lambda^s} \quad (\text{eigenvalue } 0 \text{ excluded}) \quad (3.6)$$

We have the result of Seeley:

Proposition 3.2.4. $\zeta(s)$ is holomorphic in $\{s \in \mathbb{C} | \text{Re } s > \frac{n}{2}\}$. It extends to a meromorphic function of $s \in \mathbb{C}$ with simple poles, which is holomorphic at $s = 0$.

Proof. Use the Melin transform:

$$L^{-s} = \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} e^{-tL} dt$$

$$\text{Tr}^*(L^{-s}) = \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} \text{Tr}^*(e^{-tL}) dt = \frac{1}{\Gamma(s)} \left(\int_0^1 + \int_1^\infty \right)$$

$\int_1^{+\infty} t^{s-1} \text{Tr}^*(e^{-tL}) dt$ is uniformly convergent w.r.t s , so is holomorphic in $s \in \mathbb{C}$. When $t \rightarrow 0$, we have asymptotic expansion:

$$\text{Tr}(e^{-tD^2}) = \frac{A_{-\frac{n}{2}}}{t^{\frac{n}{2}}} + \dots + A_0 + A_1 t \dots + O(t^k)$$

So

$$\frac{1}{\Gamma(s)} \int_0^1 t^{s-1} \text{Tr}^*(e^{-tL}) dt = \frac{1}{\Gamma(s)} \int_0^1 t^{s-1} \text{Tr}(e^{-tL}) dt - \frac{1}{\Gamma(s+1)} (\dim \ker L)$$

is holomorphic when $\text{Re } s > \frac{n}{2}$. □

We also define

$$\zeta^a(s) = \text{Tr}(L^{-s} P^{(a,+\infty)}) = \frac{1}{\Gamma(s)} \int_0^{+\infty} t^{s-1} \text{Tr}(e^{-tL} P^{(a,+\infty)}) dt$$

and $\zeta^{(a,b)}(s) = \text{Tr}(L^{-s} P^{(a,b)})$. Then $\forall 0 < a < b$, $\zeta^a(s) = \zeta^{(a,b)}(s) + \zeta^b(s)$. Note that

$$\frac{\partial \zeta^{(a,b)}}{\partial s}(0) = -\log \prod_{\substack{\lambda \in \text{Sp}(L) \\ a < \lambda < b}} \lambda = -\log \det L^{(a,b)} = -2 \log |\det D_+^{(a,b)}|$$

So $|\det D_+^{(a,b)}| = \exp(-\frac{1}{2} \frac{\partial \zeta^{(a,b)}}{\partial s}(0))$. We define

$$|\det D^{(a,\infty)}| = \exp(-\frac{1}{2} \frac{\partial \zeta^a}{\partial s}(0)), \quad |\det^* D_+| = \exp(-\frac{1}{2} \frac{\partial \zeta}{\partial s}(0))$$

then

$$|\det D_+^{(a,\infty)}| = |\det D^{(a,b)}| |\det D^{(b,\infty)}|, \quad |\det^* D_+| = |\det D_+^{(0,a)}| |\det D^{(a,+\infty)}|$$

Definition 3.2.4. $\|\cdot\|_{\lambda < a} = |\cdot|_{\lambda < a} |\det D_+^{(a,\infty)}| = |\cdot|_{\lambda < a} \exp(-\frac{1}{2} \frac{\partial \zeta^a}{\partial s}(0))$.

Theorem 3.2.1. Under the canonical identification of λ with $\lambda <^a$ over U_a , the metric $\|\cdot\|_{\lambda < a}$ patch into a smooth metric $\|\cdot\|_\lambda$ on λ .

the next step is to construct a natural unitary connection ∇^λ on $(\lambda, \|\cdot\|_\lambda)$. First we assume $\ker D_\pm^X$ have constant dimensions. In theorem ?? we see that

$$d\tilde{\eta} = \pi_* [\hat{A}(\nabla^{TX}) ch(\nabla^E)] - ch(\nabla^{\ker D^X})$$

so $d\tilde{\eta}^{(1)} = \pi_* [\hat{A}(\nabla^{TX}) ch(\nabla^E)]^{(2)} - c_1(\nabla^{\det \ker D^X})$.

Definition 3.2.5. $\nabla'^\lambda = \nabla^{\det \ker D^X} - 2\pi i \tilde{\eta}^{(1)}$.

then

$$c_1(\nabla'^\lambda) = c_1(\nabla^{\det \ker D^X}) + d\tilde{\eta}^{(1)} = \pi_*[\hat{A}(\nabla^{TX})ch(\nabla^E)]^{(2)}$$

Since $2\pi i \tilde{\eta}^{(1)}$ is purely imaginary, ∇'^λ is unitary with respect to $|\cdot|_\lambda$. Let g, g' be the metric associated with $\|\cdot\|_\lambda, |\cdot|_\lambda$, and θ the connection form of ∇'^λ , then

$$\begin{aligned} dg &= d\left(\frac{g}{g'}\right) = d\left(\frac{g}{g'}\right) \cdot g' + \frac{g}{g'} dg' = g \cdot d \log \frac{g}{g'} + \frac{g}{g'}(g'\theta + g'\bar{\theta}) \\ &= g(d \log \frac{g}{g'} + \theta + \bar{\theta}) \end{aligned}$$

So if we let

$$\nabla''^\lambda = \nabla'^\lambda + \frac{1}{2}d \log \frac{g}{g'} \quad (3.7)$$

then ∇''^λ is unitary with respect to $\|\cdot\|_\lambda$.

In general case, we want to patch $\nabla^{\lambda^{<a}}$ together. Let $\nabla^{H^{<a}}$ be the orthogonal projection of ∇^H to $H^{<a}$, which is unitary, then $\lambda^{<a}$ has an induced connection $\nabla^{<a}$ which is unitary w.r.t. $|\cdot|_{\lambda^{<a}}$. $\forall \sigma \in \lambda^{<a}$

$$\begin{aligned} \nabla^{\lambda^{<b}}(\varphi_b^a(\sigma)) &= \nabla^{\lambda^{<b}}(\sigma \otimes \det D_+^{(a,b)}) = \nabla^{\lambda^{<a}}\sigma \otimes (\det D_+^{(a,b)}) + \sigma \otimes (\nabla^{\lambda^{(a,b)}} \det D_+^{(a,b)}) \\ &= (\nabla^{\lambda^{<a}}\sigma + \frac{\nabla^{\lambda^{(a,b)}} \det D_+^{(a,b)}}{\det D_+^{(a,b)}}\sigma) \otimes (\det D_+^{(a,b)}) \end{aligned}$$

From formula (3.1), we know that

$$\begin{aligned} \frac{\nabla^{\lambda^{(a,b)}}(\det D_+^{(a,b)})}{\det D_+^{(a,b)}} &= -Tr[(D_+^{(a,b)})^{-1} \nabla^{H^{(a,b)}} D_+^{(a,b)}] \\ &= -\frac{1}{2}Tr[(D^{(a,b)})^{-1} \nabla^{H^{(a,b)}} D^{(a,b)}] - \frac{1}{2}Tr_s[(D^{(a,b)})^{-1} \nabla^{H^{(a,b)}} D^{(a,b)}] \end{aligned}$$

We want to define $\nabla'^{\lambda^{<a}} = \nabla^{\lambda^{<a}} - Tr[(D_+^{(a,\infty)})^{-1} \nabla^{H^{(a,\infty)}} D_+^{(a,\infty)}]$, then

$$\begin{aligned} \nabla'^{\lambda^{<b}}(\varphi_b^a(\sigma)) &= (\nabla^{\lambda^{<a}} - Tr[(D_+^{(a,b)})^{-1} \nabla^{H^{(a,b)}} D_+^{(a,b)}] - Tr[(D_+^{(b,\infty)})^{-1} \nabla^{H^{(b,\infty)}} D_+^{(b,\infty)}])\sigma \otimes \det D_+^{(a,b)} \\ &= (\nabla^{\lambda^{<a}}\sigma - Tr[(D_+^{(a,\infty)})^{-1} \nabla^{H^{(a,\infty)}} D_+^{(a,\infty)}])\sigma \otimes \det D_+^{(a,b)} = \varphi_b^a(\nabla'^{\lambda^{<a}}\sigma) \end{aligned}$$

So the connection $\{\nabla'^{<a}\}$ patch together. First, we introduce

Definition 3.2.6.

$$\begin{aligned} \gamma_t^a &= \int_t^\infty Tr[e^{-sD^2} \nabla^H D \cdot D \cdot P^{(a,\infty)}] ds = -Tr[e^{-tD^2} D^{-1} \nabla^H D \cdot P^{(a,\infty)}] \\ \delta_t^a &= \int_t^\infty Tr_s[e^{-sD^2} \nabla^H D \cdot D \cdot P^{(a,\infty)}] ds = Tr_s[e^{-tD^2} D^{-1} \nabla^H D \cdot P^{(a,\infty)}] \end{aligned}$$

similarly we can define $\gamma_t^{(a,b)}$ and $\delta_t^{(a,b)}$.

As $t \rightarrow 0$, we have asymptotic expansions:

$$\begin{aligned} \gamma_t^a &= \int_t^1 (Tr[e^{-sD^2} \nabla^H D \cdot D \cdot P^{(a,\infty)}] + \sum_{-\frac{\pi}{2}}^0 dA_j \cdot s^{j-1}) ds \\ &\quad + \int_1^\infty Tr[e^{-sD^2} \nabla^H D \cdot D \cdot P^{(a,\infty)}] ds - \sum_{-\frac{\pi}{2}}^{-1} \frac{dA_j}{j} + \sum_{-\frac{\pi}{2}}^{-1} \frac{dA_j}{j} t^j + dA_0 \cdot \log t \\ &= \sum_{-\frac{\pi}{2}}^{-1} dA_j \cdot \frac{t^j}{j} + dA_0 \cdot \log t + \gamma_0^a + O(t) \end{aligned} \quad (3.8)$$

$$\delta_t^a = \delta_0^a + O(t)$$

we have $\gamma_t^b = \gamma_t^a + \gamma_t^{(a,b)}$, $\delta_t^b = \delta_t^a + \delta_t^{(a,b)}$ and

$$\gamma_0^{(a,b)} = -Tr[D^{-1}\nabla^H D \cdot P^{(a,b)}] = -Tr[(D^{(a,b)})^{-1}\nabla^{H^{(a,b)}} D^{(a,b)}]$$

$$\delta_0^{(a,b)} = Tr_s[D^{-1}\nabla^H D \cdot P^{(a,b)}] = Tr_s[(D^{(a,b)})^{-1}\nabla^{H^{(a,b)}} D^{(a,b)}]$$

So $\frac{\nabla^{\lambda^{(a,b)}}(det D_+^{(a,b)})}{det D_+^{(a,b)}} = \frac{1}{2}(\gamma_0^{(a,b)} - \delta_0^{(a,b)})$. If we let

$$\nabla'^{\lambda^{<a}} = \nabla^{\lambda^{<a}} + \frac{1}{2}(\gamma_0^a - \delta_0^a)$$

then $\nabla'^{\lambda^{<a}}$ will patch together. But it may not be unitary. By equation (3.7), we compute $d \log \frac{g}{g'} = d(-\frac{\partial \zeta^a}{\partial s}(0))$. Recall that

$$\zeta^a(s) = Tr(L^{-s}P^{(a,\infty)}) = \frac{1}{2}Tr[(D^2)^{-s}P^{(a,\infty)}] = \frac{1}{2\Gamma(s)} \int_0^\infty t^{s-1}Tr[e^{-tD^2}P^{(a,\infty)}]dt$$

so

$$\begin{aligned} d\zeta^a(s) &= -\frac{1}{\Gamma(s)} \int_0^\infty t^s Tr[e^{-tD^2}\nabla^H D \cdot D \cdot P^{(a,\infty)}]dt \\ &= -\frac{1}{\Gamma(s)} \int_0^1 t^s (Tr[e^{-tD^2}\nabla^H D \cdot D \cdot P^{(a,\infty)}] + \sum_{-\frac{n}{2}}^0 dA_j \cdot t^{j-1})dt \\ &\quad - \frac{1}{\Gamma(s)} \int_1^\infty t^s Tr[e^{-tD^2}\nabla^H D \cdot D \cdot P^{(a,\infty)}]dt + \frac{1}{\Gamma(s)} \sum_{j=-\frac{n}{2}}^{-1} \frac{dA_j}{s+j} + \frac{dA_0}{\Gamma(s+1)} \end{aligned}$$

Note that $\Gamma(s)$ has simple pole at $s = 0$, $Res(\Gamma(s), 0) = 1$, and $(\frac{1}{\Gamma(s)})'(0) = 1$. By equation (3.8), we get

$$\begin{aligned} d[\frac{\partial \zeta^a}{\partial s}(0)] &= -\int_0^1 (Tr[e^{-tD^2}\nabla^H D \cdot D \cdot P^{(a,\infty)}] + \sum_{-\frac{n}{2}}^0 dA_j \cdot t^{j-1})dt \\ &\quad - \int_1^\infty Tr[e^{-tD^2}\nabla^H D \cdot D \cdot P^{(a,\infty)}]dt + \sum_{j=-\frac{n}{2}}^{-1} \frac{dA_j}{j} - \Gamma'(1)dA_0 \\ &= -\gamma_0^a - \Gamma'(1)dA_0 \end{aligned}$$

δ_0^a is purely imaginary, so $\nabla''^{\lambda^{<a}} = \nabla^{\lambda^{<a}} + \frac{1}{2}(\gamma_0^a - \delta_0^a) + \frac{1}{2}\Gamma'(1)dA_0$ is unitary and can be patched together.

We will see whether it satisfies $c_1(\nabla^\lambda) = \pi_*[\hat{A}(\nabla^{TX})ch(\nabla^E)]^{(2)}$, we know that as $t \rightarrow 0$, $ch(A_t) = \varphi Tr_s \exp(-A_t^2) = \pi_*[\hat{A}(\nabla^{TX})ch(\nabla^E)] + O(\sqrt{t})$. So

$$[\varphi Tr_s \exp(-A_t^2)]^{(2)} = \pi_*[\hat{A}(\nabla^{TX})ch(\nabla^E)]^{(2)} + O(\sqrt{t}) \quad (3.9)$$

Proposition 3.2.5.

$$[Tr_s \exp(-A_t^2)]^{(2)} = [Tr_s \exp(-(\sqrt{t}D^X + \nabla^H)^2)]^{(2)} \quad (3.10)$$

Proof. Let $A_t^l = \sqrt{t}D^X + \nabla^H - \frac{lc(T^H)}{4\sqrt{t}}$, then

$$\frac{\partial}{\partial l}[Tr_s \exp(-A_t^{l,2})]^{(2)} = -dTr_s[\frac{\partial A_t^l}{\partial l} \exp(-A_t^{l,2})]^{(1)} = dTr_s[\frac{c(T^H)}{4\sqrt{t}} \exp(-A_t^{l,2})]^{(1)} = 0$$

Since $A_t^l = A_t$, $A_t^0 = \sqrt{t}D^X + \nabla^H$, the proposition follows. \square

Proposition 3.2.6.

$$[Tr_s \exp(-(\sqrt{t}D^X + \nabla^H)^2)]^{(2)} = [Tr_s \exp(-(\sqrt{t}D^{<a} + \nabla^{H^{<a}})^2)]^{(2)} + [Tr_s \exp(-(\sqrt{t}D^{(a,\infty)} + \nabla^{H^{(a,\infty)}})^2)]^{(2)} \quad (3.11)$$

Proof. We use the same transgression trick: Let $M^a = \nabla^H - \nabla^{H,split}$, where $\nabla^{H,split} = \nabla^{H^{<a}} \oplus \nabla^{H^{(a,\infty)}}$, and $\nabla_l^H = \nabla^{H,split} + lM^a$. Then

$$\begin{aligned} \frac{\partial}{\partial l} [Tr_s \exp(-(\sqrt{t}D^X + \nabla_l^H)^2)]^{(2)} &= -dTr_s [M^a \exp(-(\sqrt{t}D^X + \nabla_l^H)^2)]^{(1)} \\ &= -dTr_s [M^a \exp(-tD^{X,2})] \end{aligned}$$

Clearly M^a interchanges $H_{\pm}^{<a}$ and $H_{\pm}^{(a,\infty)}$, while $\exp(-tD^{X,2})$ preserves the splitting $H_{\pm} = H_{\pm}^{<a} \oplus H_{\pm}^{(a,\infty)}$, so $Tr_s [M^a \exp(-tD^{X,2})] = 0$. Since $\nabla^{H,split}$ verifies the proposition, so does ∇^H . \square

Now

$$\begin{aligned} \frac{\partial}{\partial t} [Tr_s \exp(-(\sqrt{t}D^{(a,\infty)} + \nabla^{H^{(a,\infty)}})^2)]^{(2)} &= -dTr_s \left[\frac{D^{(a,\infty)}}{2\sqrt{t}} \exp(-(\sqrt{t}D^{(a,\infty)} + \nabla^{H^{(a,\infty)}})^2) \right]^{(1)} \\ &= \frac{1}{2} dTr_s [D^{(a,\infty)} \nabla^{H^{(a,\infty)}} D^{(a,\infty)} \exp(-tD^{(a,\infty),2})] \\ &= \frac{1}{2} dTr_s [\exp(-tD^2) \nabla^H D \cdot D \cdot P^{(a,\infty)}] \end{aligned}$$

Since $a > 0$, as $t \rightarrow +\infty$, $[Tr_s \exp(-(\sqrt{t}D^{(a,\infty)} + \nabla^{H^{(a,\infty)}})^2)]^{(2)}$ decays exponentially. So we have

$$\begin{aligned} [Tr_s \exp(-(\sqrt{t}D^{(a,\infty)} + \nabla^{H^{(a,\infty)}})^2)]^{(2)} &= -\frac{1}{2} d \int_t^{\infty} Tr_s [\exp(-sD^2) \nabla^H D \cdot D \cdot P^{(a,\infty)}] ds \\ &= -\frac{1}{2} d\delta_t^a \end{aligned} \quad (3.12)$$

By formula (3.9),(3.10),(3.11),(3.12), we have

$$[Tr_s \exp(-A_t^2)]^{(2)} = [Tr_s \exp(-(\sqrt{t}D^{<a} + \nabla^{H^{<a}})^2)]^{(2)} - \frac{1}{2} d\delta_t^a$$

Let $t \rightarrow 0$, we get

$$\pi_* [\hat{A}(\nabla^{TX}) ch(\nabla^E)]^{(2)} = c_1(\nabla^{\lambda^{<a}}) - \frac{1}{2\pi i} \frac{1}{2} d\delta_0^a \quad (3.13)$$

Now we have $\nabla''^{\lambda^{<a}} = \nabla^{\lambda^{<a}} + \frac{1}{2}(\gamma_0^a - \delta_0^a) + \frac{1}{2}\Gamma'(1)dA_0$. Note that $\gamma_0^a = -d[\frac{\partial \zeta^a}{\partial s}(0)] - \Gamma'(1)dA_0$ is exact, so

$$c_1(\nabla''^{\lambda^{<a}}) = c_1(\nabla^{\lambda^{<a}}) - \frac{1}{2\pi i} \cdot \frac{1}{2} d\delta_0^a \quad (3.14)$$

So $\nabla''^{\lambda^{<a}}$ is the connection we want:

Theorem 3.2.2. *Let ${}^1\nabla^{\lambda^{<a}} = \nabla^{\lambda^{<a}} + \frac{1}{2}(\gamma_0^a - \delta_0^a) + \frac{1}{2}\Gamma'(1)dA_0$, then identifying λ with $\lambda^{<a}$ over U_a , the connection ${}^1\nabla^{\lambda^{<a}}$ patch together into a connection ${}^1\nabla^{\lambda}$ on λ , which is unitary w.r.t. the metric $\|\cdot\|$, and*

$$c_1({}^1\nabla^{\lambda}) = \pi_* [\hat{A}(\nabla^{TX}) ch(\nabla^E)]^{(2)}$$