Recall: Phase Line for first order (linear) autonomous ODE

$$
y^{\prime}=-2 y
$$

Phase line:


The equilibrium may appear as: stable; semistable; unstable
Phase Portrait: 2-dim generalization
For a homogeneous $2 \times 2$ linear system

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{ll}
a_{11} & a_{12} \\
a_{21} & a_{22}
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

(notice the right-hand-side does not depend on t , aka, autonomous)

$$
\left[\begin{array}{l}
0 \\
0
\end{array}\right] \text { is (always) an equilibrium }
$$

Therefore we can talk about its stability. As one might imagine, there are 9 different types.

Example 1: (Nodal Source)

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{cc}
3 / 2 & -1 / 2 \\
1 / 2 & 3 / 2
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

The eigenvalues are 2 and 1, distinct and both positive. The general solution is

$$
\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]=C_{1} e^{2 t}\left[\begin{array}{l}
1 \\
1
\end{array}\right]+C_{2} e^{t}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]
$$

To draw the phase portrait:
(1) $C_{1}=0, C_{2}=0$. You get the equilibrium $\left[\begin{array}{l}0 \\ 0\end{array}\right]$

(2) $C_{1} \neq 0, C_{2}=0$. Then $\vec{x}(t)=C_{1} e^{2+}\left[\begin{array}{l}1 \\ 1\end{array}\right]$

It moves along the line $x_{1}=x_{2}$
Whenever $C_{1}>0, \vec{x}(t)$ moves in the first quadrant

$$
C_{1}<0 \quad \text { third }
$$

In both cases, as $t \nearrow . \vec{X}(t)$ moves away from $\left[\begin{array}{l}0 \\ 0\end{array}\right]$


## WARNING: We just gave TWO integral curves! None of them ever passes

 the origin!(3) $C_{1}=0, c_{2} \neq 0$. Then $\vec{x}(t)=C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$

It moves along the line $x_{1}=-x_{2}$
Whenever $C_{2}>0, \vec{x}(t)$ moves in the fourth quadrant $C_{2}<0 \quad$ second
In both cases, as $t \uparrow, \vec{x}(t)$ moves away from $\left[\begin{array}{l}0 \\ 0\end{array}\right]$

(4) For generic $C_{1}, C_{2}$, we try to draw using the asymptotic behavior:
(4.1) If $t \rightarrow \infty$
$\vec{x}(t)=C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]+C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$ dominated by $C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$
$\vec{x}(t)=2 C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]+C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$ dominated by $2 C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$
So starting at any point in the plane, as $t$ grows large, the

integral carne tends to become parallel to the line $x_{1}=x_{2}$
As $t \uparrow . \vec{x}(t)$ moves away from the origin.
(4.2) If $t \rightarrow-\infty$
$\vec{x}(t)=C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]+C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$ dominated by $C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$
$\vec{x}^{\prime}(t)=2 C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1 \\ 1\end{array}\right]+C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$ dominated by $C_{2} e^{t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$

So starting at any point in the plane, as $t$ grows large, the integral cure tends to become parallel to the line $x_{1}=-x_{2}$ As $t \nearrow . \vec{x}(t)$ moves away from the origin. As $t \rightarrow-\infty, \vec{x}(t)$ approndes $\left[\begin{array}{l}0 \\ 0\end{array}\right]$.
So basically the integral curve is "tangent" to $x_{1}=-x_{2}$

(4.3) Drawing one means drawing all:


Remark: The equilibrium is obviously "unstable" under perturbation in all directions. The term nodal means it's not "degenerate" (a degenerate example will seen in the case of repeated eigenvalue). The term source is self-evident if you imagine the integral curves as "flows".

Remark: The above process can actually be simplified, as will be seen from the next example.

Example 2: (Nodal Sink)

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{cc}
-3 / 2 & 1 / 2 \\
1 / 2 & -3 / 2
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

The eigenvalues are -2 and -1 , distinct and both negative. The general solution is

$$
\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]=C_{1} e^{-2 t}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2} e^{-t}\left[\begin{array}{l}
1 \\
1
\end{array}\right]
$$

The phase portrait is almost the same as above, except the arrows are reversed and also the integral curves are curved differently.


(2) Look at generic $C_{1}, C_{2}$ :

If $t \rightarrow \infty$
$\vec{x}(t)=C_{1} e^{-2 t}\left[\begin{array}{c}1 \\ -1\end{array}\right]+C_{2} e^{-t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$ dominated by $C_{2} e^{-t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$
$\vec{x}(t)=-2 C_{1} e^{-2+}\left[\begin{array}{c}1 \\ -1\end{array}\right]-C_{2} e^{-t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$ dominated by $-C_{2} e^{-t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$
So starting at any point in the plane, as $t$ grows large, $\vec{x}(t)$ approaches to $\left[\begin{array}{l}0 \\ 0\end{array}\right]$ and will appear tangent to $x_{1} x_{2}$

(3) If $t \rightarrow-\infty$
$\vec{x}(t)=C_{1} e^{-2+}\left[\begin{array}{c}1 \\ -1\end{array}\right]+C_{2} e^{-t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$ dominated by $C e^{-2 t+[ }\left[\begin{array}{c}1 \\ -1\end{array}\right]$
$\vec{x}(t)=-2 C_{1} e^{-2 t}\left[\begin{array}{c}1 \\ -1\end{array}\right]-C_{2} e^{-t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$ dominated by $-2 C_{2} e^{-2 t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$
So starting at any point in the plane, as $t \rightarrow-\infty$ $\vec{x}(t)$ will appear parallel to $x_{1}=-x_{2}$ As $t \boldsymbol{\lambda}, \vec{x}(t)$ will move towards the origin.
(4) Draw all other curves.

Remark: The equilibrium is "stable" under perturbation to all directions. The term sink is self-evident.


## Example 3: (Saddle Point)

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{ll}
1 / 2 & 3 / 2 \\
3 / 2 & 1 / 2
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

The eigenvalues are 2 and -1 , one positive and the other negative. The general solution is

$$
\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]=C_{1} e^{2 t}\left[\begin{array}{l}
1 \\
1
\end{array}\right]+C_{2} e^{-t}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]
$$

The phase would be very different. To draw it:
(1) First mark the equilibrium and draw along the eigenvectors (Coresp. to $C_{1} \neq 0, C_{2}=0$ and $C_{1}=0, C_{2} \neq 0$ )

(2) Look at generic $C_{1}, C_{2}$ :

If $t \rightarrow \infty$
$\vec{x}(t)=C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]+C_{2} e^{-t}\left[\begin{array}{c}1 \\ -1\end{array}\right] \sim C_{1} e^{2 t}\left[\begin{array}{l}1 \\ 1\end{array}\right]$
(the second term dies out)

(3) If $t \rightarrow-\infty$
$\vec{x}(t)=C_{1} e^{2+}\left[\begin{array}{l}1 \\ 1\end{array}\right]+C_{2} e^{-t}\left[\begin{array}{c}1 \\ -1\end{array}\right] \sim C_{2} e^{-t}\left[\begin{array}{c}1 \\ -1\end{array}\right]$ (the first term dies out)

(4) Draw other integral curves


Remark: In this case the equilibrium is "stable" under perturbation from one direction but "unstable" from ALL other directions. This is far from being "semistable" and basically speaking it's "unstable".

Remark: Situation in $2 \times 2$ system is so different that we should use a completely new set of terminologies. The term saddle point is clear if the integral curves are "imagined" as the gradient field (in fact it is the "flow" generated by the gradient field).

## Example 4: (Proper Node)

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{ll}
2 & 0 \\
0 & 2
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

The eigenvalues are 2 and 2, repeated and positive. For this special case, one can find two linearly independent eigenvectors and write the general solution as

$$
\begin{aligned}
{\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right] } & =C_{1} e^{2 t}\left[\begin{array}{l}
1 \\
0
\end{array}\right]+C_{2} e^{2 t}\left[\begin{array}{l}
0 \\
1
\end{array}\right] \\
& =e^{2 t}\left[\begin{array}{l}
C_{1} \\
C_{2}
\end{array}\right]
\end{aligned}
$$

Notice that every vector in the plane is now an eigenvector. So the integral curve is a straight line and the phase portrait looks like


Exercise: Draw the phase portrait for the following system

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{cc}
-3 & 0 \\
0 & -3
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

## Example 5: (Improper Node)

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{cc}
-1 & -1 \\
1 & 1
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

The eigenvalues are 0 and 0 , repeated zero. For this matrix one can obtain only "one" eigenvector and to solve the system one has to get the generalized eigenvector. The general solution is

$$
\begin{aligned}
{\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right] } & =C_{1} e^{0 t}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2} e^{0 t}\left(t\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+\left[\begin{array}{l}
0 \\
1
\end{array}\right]\right) \\
& =C_{1}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2}\left[\begin{array}{l}
0 \\
1
\end{array}\right]+C_{2} t\left[\begin{array}{c}
1 \\
-1
\end{array}\right]
\end{aligned}
$$

The phase portrait for this system is different but easy
(1) Mark the equilibrium $\left[\begin{array}{l}0 \\ 0\end{array}\right]$
$S_{\text {et }} C_{1} \neq 0, C_{2}=0$
Observe that $\vec{x}(t)$ is constantly $\left[\begin{array}{c}c_{1} \\ -c_{L}\end{array}\right]$
So every point along the line is an equilibrium.
In this case we draw the line amy way but with No arrows.
(2) Set $C_{1}=0, C_{2} \neq 0$.
$\vec{X}(t)=C_{2}\left[\begin{array}{c}t \\ 1-t\end{array}\right]$ moves along $x_{2}=C_{2}-x_{1}$
When $C_{2}>0, t \uparrow \Rightarrow x_{1} \lambda$. $C_{2}<0, t \uparrow \Rightarrow x_{1} \downarrow$
Mark the arrows accordingly
(3) $C_{1} \neq 0$ doesn't give anything new!

Translating a line via a vector gives another line that is parallel.

So this is the phase portrait.


Example 5: (Improper Node)

$$
\left[\begin{array}{l}
x_{1}^{\prime} \\
x_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{cc}
2 & -1 \\
1 & 4
\end{array}\right]\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right]
$$

The eigenvalues are 3 and 3 , repeated and positive. For this matrix one can obtain only "one" eigenvector and to solve the system one has to get the generalized eigenvector. The general solution is

$$
\begin{aligned}
{\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right] } & =C_{1} e^{3 t}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2} e^{3 t}\left(t\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+\left[\begin{array}{l}
0 \\
1
\end{array}\right]\right) \\
& =\left(C_{1}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2}\left[\begin{array}{l}
0 \\
1
\end{array}\right]+C_{2} t\left[\begin{array}{c}
1 \\
-1
\end{array}\right]\right) e^{3 t}
\end{aligned}
$$

This case differs to above by simply an exponential factor. We shall use the above example to figure out what we have here.
(1) Along the line defined by the eigenvector, due to the presence of $e^{3 t}$,
$t \xlongequal{ } \uparrow \vec{x}(t)$ moves "away " from $\left[\begin{array}{l}0 \\ 0\end{array}\right]$
Mark the generalized eigenvector $\vec{u}=\left[\begin{array}{l}0 \\ 1\end{array}\right]$

(2) For $C_{1}=0, C_{2} \neq 0$,

$\vec{X}^{\prime}(t)=C_{2}\left[\begin{array}{c}1+3 t \\ -1+3-3 t\end{array}\right] e^{3 t}=C_{2}\left[\begin{array}{c}1 \\ 2\end{array}\right]+3 C_{2}\left[\begin{array}{c}1 \\ -1\end{array}\right] e^{3 t}$
$S_{a y} C_{2}>0$, (comespond to the region the gene e.v. profiting to)
$t \rightarrow \infty, \vec{x}(t)$ explodes, dominated by $C_{2} t\left[\begin{array}{c}1 \\ -1\end{array}\right] e^{3 t}$ $\vec{x}^{\prime}(t)$ explodes, dominated by $3 C_{2}+\left[\begin{array}{c}1 \\ -1 \\ -1\end{array}\right] e^{3 t}$
So the integral arne will tend to parallel| to $\left[\begin{array}{c}1 \\ -1\end{array}\right]$

(3) $t \rightarrow-\infty, \vec{x}(t)$ approaches $t o\left[\begin{array}{l}0 \\ 0\end{array}\right]$

$$
\begin{aligned}
& \vec{x}(t) \text { dominated by } C_{2} t\left[\begin{array}{c}
1 \\
-1
\end{array} e^{3+}\right. \\
& \vec{x}^{\prime}(t) \text { dominated by } S_{2} t\left[\begin{array}{c}
1 \\
-1
\end{array}\right] e^{3 t}
\end{aligned}
$$

Notice $t \leq 0$, so the integral cure tends to $\left[\begin{array}{l}0 \\ 0\end{array}\right]$

$$
\text { tangently } \begin{array}{ll}
\text { to }
\end{array}\left[\begin{array}{c}
-1 \\
1
\end{array}\right]
$$


(4) Draw the other integral ames
(5) For $C_{2}<0$, repeat the analysis above to get the phase poothait as shown.

Remark: This differs to the degenerate case by an exponential factor, which can be thought as "winding" the straight lines into curved lines.


## Example 6: (Improper Node)

The eigenvalues are -3 and -3 , repeated and positive. For this matrix one can obtain only "one" eigenvector and to solve the system one has to get the generalized eigenvector. The general solution is

$$
\begin{aligned}
{\left[\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right] } & =C_{1} e^{-3 t}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2} e^{-3 t}\left(t\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+\left[\begin{array}{c}
0 \\
-1
\end{array}\right]\right) \\
& =\left(C_{1}\left[\begin{array}{c}
1 \\
-1
\end{array}\right]+C_{2}\left[\begin{array}{c}
0 \\
-1
\end{array}\right]+C_{2} t\left[\begin{array}{c}
1 \\
-1
\end{array}\right]\right) e^{-3 t}
\end{aligned}
$$

