

INTEGRAL GROUP ACTIONS ON SYMMETRIC SPACES AND DISCRETE DUALITY SYMMETRIES OF SUPERGRAVITY THEORIES

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ABSTRACT. For $G = G(\mathbb{R})$ a split, simply connected, semisimple Lie group of rank n and K the maximal compact subgroup of G , we give a method for computing Iwasawa coordinates of $K \backslash G$ using the Chevalley generators and the Steinberg presentation. When $K \backslash G$ is a scalar coset for a supergravity theory in dimensions ≥ 3 , we determine the action of the integral form $G(\mathbb{Z})$ on $K \backslash G$. We give explicit results for the action of the discrete U-duality groups $SL_2(\mathbb{Z})$ and $E_7(\mathbb{Z})$ on the scalar cosets $SO(2) \backslash SL_2(\mathbb{R})$ and $[SU(8)/\{\pm \text{Id}\}] \backslash E_{7(+7)}(\mathbb{R})$ for type IIB supergravity in ten dimensions and 11-dimensional supergravity reduced to $D = 4$ dimensions, respectively. For the former, we use this to determine the discrete U-duality transformations on the scalar sector in the Borel gauge and we describe the discrete symmetries of the dyonic charge lattice. We determine the spectrum-generating symmetry group for fundamental BPS solitons of type IIB supergravity in $D = 10$ dimensions at the classical level and we propose an analog of this symmetry at the quantum level. We indicate how our methods can be used to study the orbits of discrete U-duality groups in general.

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1. INTRODUCTION

In classical supergravity theories, one is interested in classifying solutions of the equations of motion that enjoy certain desirable properties. These often take the form of extremal black p -brane solutions, and are usually related by the underlying algebraic structure of classical U-duality. In the quantum theory, that is in string theory or in M-theory, the corresponding solutions are believed to be related by discrete U-duality ([HT]), leading to a subtle and complex orbit structure under the discrete U-duality group.

Thus knowing the orbits of the U-duality group on the space of solutions allows one, at least in principle, to know all solutions of a given class of equations in the quantum theory. Such solutions are usually charged under (higher analogs of) electric and magnetic fields. The electric and magnetic charge vectors of the asymptotically flat p -brane solutions form irreducible representations of the U-duality group ([LPS2]).

The question of determining the distinct charge vector orbits under U-duality is of central importance. The U-duality orbits for real valued charges are classified and have been understood for some time ([FG], [LPS2]). However, a classification of orbits under the discrete U-duality group is known only in certain cases (see [BDDFMR] for an excellent overview of the status of this question and the open issues).

We recall that in dimensionally reduced supergravity theories, there is a Lie group G and subgroup K such that the scalar fields take values in the space $K \backslash G$ of right cosets. This coset space has a K -action on the left and a right action by the discrete U-duality group $G(\mathbb{Z})$. In most cases, $G = G(\mathbb{R})$ is a split, simply connected, semisimple algebraic group of rank n and K is its maximal compact subgroup.

In this work, we give a method using the Steinberg presentation for computing Iwasawa coordinates of the quotient $K \backslash G$ and the action of the \mathbb{Z} -form $G(\mathbb{Z})$ on $K \backslash G$ by means of Iwasawa decomposition $G = KH^+U$ (see Sections 3 and 4). These computations encode the discrete U-duality symmetries on the scalar coset $K \backslash G$. Similar computations are common in the theory of automorphic forms on finite dimensional Lie groups (see [H-C]).

Our computations hence involve finding the standard Iwasawa form of a generic element $g \in G$ and then allowing $z \in G(\mathbb{Z})$ to act on the coset Kg . The task is then to determine a compensating element $k \in K$ so that kgz is again in standard Iwasawa form. This requires completely determining the constants given in the Steinberg relations for the group G ([St]) and explicitly describing K and $G(\mathbb{Z})$ in terms of the Steinberg generators. The computer algebra system

Magma ([BCFS]) is essential here, since it automatically computes these constants without infeasible hand calculations.

The hardest part of constructing the action of $G(\mathbb{Z})$ seems to be the action of the Weyl group representatives in G . Tables 1 through 8 give the Iwasawa coordinates and compensating elements for the action of these elements, for all finite-dimensional U-duality groups G . As examples, we compute the action of the generators of the discrete groups $\mathrm{SL}_2(\mathbb{Z})$ and $E_{7(+7)}(\mathbb{Z})$ on the cosets $\mathrm{SO}(2)\backslash\mathrm{SL}_2(\mathbb{R})$ and $[\mathrm{SU}(8)/\{\pm\mathrm{Id}\}]\backslash E_{7(+7)}(\mathbb{R})$. These cosets are known to encode the scalar fields of type IIB supergravity in $D = 10$ dimensions and 11-dimensional supergravity reduced to $D = 4$ dimensions, respectively.

These techniques can also be extended to determine the orbits of $G(\mathbb{Z})$ on the charge lattice. Our methods are general and do not depend on the choice of group G . Our techniques could hence be applied to all U-duality groups, including the hyperbolic Kac–Moody group E_{10} , which is conjectured to be the U-duality group of 11-dimensional supergravity reduced to 1 dimension (see Section 7).

Earlier fundamental works, particularly [BDDFMR], make use of Jordan algebras and Freudenthal systems (as described in [Bh] and [Kr]) and of integral forms of groups (as described in [BDDR], [EG] and [Gr]) to study extremal black p -brane solutions of supergravity in 6, 5 and 4-dimensional classical theories and their relationship under U-duality. After preparation of this manuscript, we also learned of the work of Cacciatori, Cerchiai and Marrani ([CC], [CCM1]–[CCM5]) who studied various properties of the Iwasawa parametrization and gave a general method for constructing it explicitly at the group level. In particular, in [CCM1] and [CCM4], they gave Iwasawa coordinates in the symplectic frame for the symmetric space $[\mathrm{SU}(8)/\{\pm\mathrm{Id}\}]\backslash E_{7(+7)}(\mathbb{R})$ and described some of the implications for the corresponding supergravity theory.

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2. PRELIMINARIES ON U-DUALITY

A rank d U-duality group arises in the reduction of M-theory on a $(d + 1)$ -torus T^{d+1} . It is generated by the T-duality group $\mathrm{SO}(d, d)$ and the S-duality group $\mathrm{SL}_{d+1}(\mathbb{Z})$, which acts geometrically as the modular group of the torus T^{d+1} . The U-duality group is then

$$\mathrm{SL}_{d+1}(\mathbb{Z}) \bowtie \mathrm{SO}_{d,d}(\mathbb{Z}),$$

the group generated by the two non-commuting actions (see [OP] for a survey).

For example, dimensional reduction of M-theory on the torus T^7 leads to $D = 4$, $N = 8$ maximal supergravity which has non-compact $E_{7(+7)}(\mathbb{R})$ global symmetry action on the 56-dimensional fundamental representation of the Lie algebra E_7 .

In $D = 4$ spacetime dimensions, 11-dimensional supergravity has the maximal number $N = 8$ of supersymmetries. The equations of motion and the Bianchi identities of the $N = 8$ supergravity theory in four dimensions are invariant under the non-compact U-duality group $E_{7(+7)}(\mathbb{R})$ ([CJ2]).

In $D = 4$ dimensions, the electric and magnetic charges are subject to the Dirac–Zwanziger–Schwinger quantization condition. The electric and magnetic charges live on a lattice \mathcal{Q} in a 56-dimensional vector space V , with 28 electric and 28 magnetic fundamental charges. The representation V gives rise to a faithful representation of $E_{7(+7)}(\mathbb{R})$ in $\text{Sp}_{56}(\mathbb{R})$.

The discrete group $E_{7(+7)}(\mathbb{Z})$ acts on ([CJ2]):

- the set of magnitudes of electric and magnetic charges on the lattice \mathcal{Q} given by the Dirac–Zwanziger–Schwinger quantization condition;
- the abelian gauge fields as generalized electromagnetic duality; and
- the 70 scalar fields of the theory, which take values in the coset space

$$[\text{SU}(8)/\{\pm \text{Id}\}] \backslash E_{7(+7)}(\mathbb{R}).$$

The \mathbb{Z} -form $G(\mathbb{Z})$, for any simple and simply connected Chevalley group or Kac–Moody group G , may be defined as the stabilizer

$$G(\mathbb{Z}) = \{g \in G(\mathbb{R}) \mid g \cdot V_{\mathbb{Z}} \subseteq V_{\mathbb{Z}}\}$$

of a lattice $V_{\mathbb{Z}}$ in a fundamental representation V for G ([St], [BC]).

So for $G = E_7$ we get

$$E_{7(+7)}(\mathbb{Z}) = E_{7(+7)}(\mathbb{R}) \cap \text{Sp}_{56}(\mathbb{Z}),$$

as discovered in [HT], following [CJ] in the framework of type II string theory.

Soulé ([S]) gave a rigorous mathematical proof that the $E_{7(+7)}(\mathbb{Z})$ of Hull and Townsend coincides with the Chevalley \mathbb{Z} -form of $G = E_7$. Here $E_{7(+7)}(\mathbb{R}) \cap \text{Sp}_{56}(\mathbb{Z})$ is the stabilizer of the standard lattice in the 56-dimensional fundamental representation of E_7 . (See also [MS] for a discussion of the role of E_8).

2.1. $\text{SL}_2(\mathbb{Z})$ -symmetry in supergravity. The U-duality group $G(\mathbb{Z}) = E_{n(+n)}(\mathbb{Z})$ has many subgroups isomorphic to $\text{SL}_2(\mathbb{Z})$. We highlight two particularly distinguished ones. The first interchanges the NS-NS (Neveu–Schwarz) fields with the R-R (Ramond–Ramond) fields (called X-duality in [LPS1]) and is conjectured to be a non-perturbative symmetry of type IIB superstring theory in $D = 10$ ([Sc]) and type IIA superstring theory in $D = 9$ ([BHO]).

The second $\text{SL}_2(\mathbb{Z})$ subgroup implements electric-magnetic duality ([DL]), which exists only in $D = 4$ space-time dimensions, and is again non-perturbative.

2.2. The role of the Weyl group. The role of the Weyl group W associated to the discrete U-duality group $G(\mathbb{Z})$ is analogous to that of the \mathbb{Z}_2 subgroup of the $\text{U}(1)$ electric-magnetic duality group in Maxwell theory, which describes the discrete interchange of electric and magnetic fields

$$E \mapsto B \text{ and } B \mapsto -E.$$

The Weyl group acts as rotations by integral multiples of $\pi/2$ on the space of field strengths. In contrast, the full U-duality group G (over \mathbb{R} or \mathbb{Z}) includes more general rotations in the space of field strengths ([LPS1]). The Weyl group preserves the total number of electric and magnetic charges, whereas the U-duality group does not, so that W gives a characterization of the independent p -brane solutions of a given type ([LPS1]).

2.3. The classical coset. The dimensional reduction of 11-dimensional supergravity on a d -torus gives rise to a theory in D dimensions with the scalar fields having the following symmetry pattern given by Cremmer and Julia [CJ1] (see [HPS], [Ni]):

D	d	E_d	$G = E_{d(+d)}$	K	#scalars = $\dim K \backslash G$
9	2	A_1	$\mathrm{SL}_2(\mathbb{R})$	$\mathrm{SO}(2)$	$4 - 1 = 3$
8	3	$A_2 A_1$	$\mathrm{SL}_3(\mathbb{R}) \times \mathrm{SL}_2(\mathbb{R})$	$\mathrm{SO}(3) \times \mathrm{SO}(2)$	$11 - 4 = 7$
7	4	A_4	$\mathrm{SL}_5(\mathbb{R})$	$\mathrm{SO}(5)$	$24 - 10 = 14$
6	5	D_5	$\mathrm{Spin}(5, 5)$	$(\mathrm{Spin}(5) \times \mathrm{Spin}(5)) / \{\pm \mathrm{Id}\}$	$45 - 20 = 25$
5	6	E_6	$E_{6(+6)}(\mathbb{R})$	$\mathrm{Sp}(8) / \{\pm \mathrm{Id}\}$	$78 - 36 = 42$
4	7	E_7	$E_{7(+7)}(\mathbb{R})$	$\mathrm{SU}(8) / \{\pm \mathrm{Id}\}$	$133 - 63 = 70$
3	8	E_8	$E_{8(+8)}(\mathbb{R})$	$\mathrm{Spin}(16) / \{\pm \mathrm{Id}\}$	$248 - 120 = 128$

We recall that the notation $E_{d(+d)}$ refers to the non-compact split, simply connected form of E_d .

A systematic way to study the orbits of the global symmetry groups $G = E_{11-D,+(11-D)}$ in D dimensions, on the charge vector space, is to start with the simplest solution and then apply the global symmetry group G . Since the global symmetry commutes with supersymmetry, the generated solutions will also be supersymmetric. This study was carried out in [LPS2] for $4 \leq D \leq 9$ where a classification of orbits was also given.

To give the spectrum of BPS solitons, one needs to classify the sets of solutions at fixed values of the scalar moduli, that is, the asymptotic values of all the diatonic and axionic scalars.

We recall the case of $D = 4$, following [CJ2], [HT], [ADFFT] and [BFT]. The equations of motion and the Bianchi identities of the $N = 8$ classical supergravity theory in four dimensions are invariant under the classical U-duality group $E_{7(+7)}(\mathbb{R})$. The group $E_{7(+7)}(\mathbb{R})$ acts simultaneously on both the space spanned by the 70 scalar fields ϕ^α and on the vector space \mathbf{Q} generated by the 28 electric and 28 magnetic quantized charges in the 56-dimensional fundamental representation.

A static, spherically symmetric BPS black hole solution is characterized in general by a vector $\vec{Q} \in \mathbf{Q}$ and a particular point ϕ_∞ on the moduli space of the theory whose 70 coordinates ϕ_∞^α are the values of the scalar fields at spatial infinity ($r \rightarrow \infty$).

Acting on a black hole solution (ϕ_∞, \vec{Q}) by means of a U-duality transformation $g \in E_{7(+7)}(\mathbb{R})$ one generates a new black hole solution $(\phi_\infty^g, \vec{Q}^g)$. The BPS black hole solutions therefore fill the U-duality orbits.

The orbits of the classical $E_{7(+7)}(\mathbb{R})$ group are studied in [FM], [FG] and [LPS2]. The orbits can be viewed as similar, in a sense, to the orbits of time-like, light-like and space-like vectors in Minkowski space, except that one uses a quartic invariant I_4 instead of a quadratic form.

The different orbits with various supersymmetries can also be related to intersecting branes and are characterized by certain group invariant polynomials ([FM]), namely the analogous quartic invariant I_4 in four dimensions ([KK]).

2.4. The quantum coset. In the full quantum theory, it is conjectured that charges are quantized and the duality symmetry is broken to the discrete U-duality subgroup $G(\mathbb{Z})$ ([HT]) as a consequence of the Dirac–Zwanziger–Schwinger charge quantization condition. Classical supergravity solutions correspond to the limit of large values of integer quantized charges.

The orbits of the charge vectors under the integral group $G(\mathbb{Z})$ are subtle and a complete characterization is not yet known. Partial results on discrete orbit classification are made in certain cases by introducing new arithmetic U-duality invariants not appearing in the real form. These are given by the ‘greatest common divisor’ of U-duality representations built out of the basic charge vector representations ([BDDFMR], [BDDR], [DGN], [Se]).

However, a classification of orbits under the discrete U-duality group is known in certain cases. For example, for the subclass of ‘projective black holes’ satisfying an arithmetic condition (Section 4.4 of [BDDFMR]), a complete classification was given in [BDDFMR] in terms of an $E_{7(+7)}(\mathbb{Z})$ -orbit of a certain canonical form. Here the techniques involved the use of integral Jordan algebras, the integral Freudenthal triple system and the work of Krutelevich ([Kr]). In [BDDFMR], it was shown that all black holes with the same quartic norm are U-duality related, while the situation for non-projective black holes remains unclear.

2.5. The relationship between $G(\mathbb{R})$ and $G(\mathbb{Z})$ orbits. For completeness, here we state two observations, the first from [BDDFMR] and the second from [FMMS].

First, the $G(\mathbb{Z})$ orbits fall into disjoint sets corresponding to the orbits of the classical theory under the continuous group $G(\mathbb{R})$. This follows from the fact that the conditions separating the continuous orbits are manifestly invariant under the corresponding discrete U-dualities and hence states that are unrelated in the continuous case remain unrelated in the discrete case.

Second, in order for the greatest common divisor of a collection of U-duality representations to be well defined, such representations must be non-vanishing ([BDDFMR]). One must then identify the class of orbits to which a given state belongs to in the continuous case. This in turn yields an identification of the subset of arithmetic invariants that are well defined for this particular state.

2.6. The impact on M-theory of identifying the $E_{7(+7)}(\mathbb{Z})$ invariants. Manifestly $E_{7(+7)}(\mathbb{Z})$ invariant partition functions would tell us about the full microscopic physics of M-theory (on a seven-torus).

As discussed above, the integral form $E_{7(+7)}(\mathbb{Z})$ has discrete invariants that are not seen in the real form and are given by the greatest common divisor of certain sets of numbers which correspond to covariant tensors of $E_{7(+7)}(\mathbb{R})$ ([BDDR], [Se]).

As shown in [BFK], some of the orbits of $E_{7(+7)}(\mathbb{Z})$ should play an important role in counting the micro-states of $D = 4$, $N = 8$ supergravity.

A relation between time-like, light-like and space-like orbits of the $E_{7(+7)}(\mathbb{R})$ symmetry and discrete $E_{7(+7)}(\mathbb{Z})$ invariants is established in [BFK]. The time-like, light-like, and space-like orbits in $E_{7(+7)}(\mathbb{R})$ corresponds to $I_4 > 0$, $I_4 = 0$, and $I_4 < 0$, respectively.

3. ACTION OF $SL_2(\mathbb{Z})$ ON $SO(2) \backslash SL_2(\mathbb{R})$

In this section we work out the case for $G = SL_2(\mathbb{R})$ in detail. We begin with the details of the action of $SL_2(\mathbb{Z})$ on $SO(2) \backslash SL_2(\mathbb{R})$ in matrix form and then reduce it to a discussion that uses only Chevalley generators of the group $SL_2(\mathbb{R})$. This allows us to generalize the discussion to higher rank in the next section.

Recall that $\mathrm{SO}(2) = \{A \in \mathrm{SL}_2(\mathbb{R}) \mid AA^T = A^T A = \mathrm{Id}\}$. The group $\mathrm{SL}_2(\mathbb{R})$ acts on the Poincaré upper half plane $\mathbb{H} = \{z \in \mathbb{C} \mid \mathrm{Im}(z) > 0\}$ by fractional linear transformations

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} : z \mapsto \frac{az + b}{cz + d}.$$

As is well known (see [Mi]), the action of $\mathrm{SL}_2(\mathbb{R})$ on \mathbb{H} is transitive and the stabiliser of $i \in \mathbb{H}$ is $\mathrm{SO}(2)$. Hence the coset space $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R})$ is homeomorphic to \mathbb{H} via the map

$$\mathrm{SO}(2)A \mapsto \frac{ai + b}{ci + d} \quad \text{where } A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

The \mathbb{Z} -form $G(\mathbb{Z}) = \mathrm{SL}_2(\mathbb{Z})$ acts on \mathbb{H} by integral fractional linear transformations with fundamental domain $\mathcal{F} = \{z \in \mathbb{H} \mid |z| \geq 1 \text{ and } |\mathrm{Re}(z)| \leq 1/2\}$. It is well known that $\mathrm{SL}_2(\mathbb{Z})$ is generated by

$$T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} : z \mapsto z + 1 \quad \text{and} \quad S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} : z \mapsto -1/z.$$

Writing $a = TS$, $b = S$ we obtain group presentations

$$\begin{aligned} \mathrm{SL}_2(\mathbb{Z}) &= \langle a, b \mid a^6 = b^4 = 1, a^3 = b^2 \rangle, \\ \mathrm{PSL}_2(\mathbb{Z}) &= \langle a, b \mid a^3 = b^2 = 1 \rangle. \end{aligned}$$

We also note that S sends $z \in \mathbb{H}$ to $-1/z$ and rotates \mathcal{F} about the point i , fixing i as shown in Figure 1. This is the only fixed point for the action of S on \mathbb{H} . Furthermore, S inverts the arc from $e^{2\pi i/3}$ to $e^{\pi i/3}$, while TS fixes $e^{\pi i/3}$.

Since we will use the Steinberg presentation to generalize our results in the next section, we briefly describe the construction of $\mathrm{SL}_2(\mathbb{R})$ as a Chevalley group. Recall that the matrices

$$e = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \quad h = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad f = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

form a basis for the Lie algebra $\mathfrak{sl}_2(\mathbb{R})$ of the Lie group $\mathrm{SL}_2(\mathbb{R})$. The multiplication table for this Lie algebra is completely determined by the relations

$$[e, f] = h, \quad [h, e] = 2e, \quad [h, f] = -2f.$$

Let $\mathfrak{h} = \mathbb{R}h$, let \mathfrak{h}^* be the dual vector space, and let $\alpha \in \mathfrak{h}^*$ be defined by $\alpha(h) = 2$. Then $\mathfrak{sl}_2(\mathbb{R})$ has a *root space decomposition*, that is the eigenspace decomposition of $\mathfrak{sl}_2(\mathbb{R})$ with respect to the adjoint $\mathrm{ad} h$, where $(\mathrm{ad} h)(X) = [h, X]$. The roots are $\alpha, -\alpha \in \mathfrak{h}^*$ with corresponding root vectors $e_\alpha = e$ and $e_{-\alpha} = f$.

The map

$$\exp : \mathfrak{sl}_2(\mathbb{R}) \longrightarrow \mathrm{SL}_2(\mathbb{R})$$

is the matrix exponential $\exp(X) = \sum_{k=0}^{\infty} \frac{1}{k!} X^k$. Now $\mathrm{SL}_2(\mathbb{R})$ has Steinberg generators

$$\chi_\alpha(x) = \exp(xe_\alpha) = \mathrm{Id} + xe = \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad \chi_{-\alpha}(x) = \exp(xe_{-\alpha}) = \mathrm{Id} + xf = \begin{pmatrix} 1 & 0 \\ x & 1 \end{pmatrix},$$

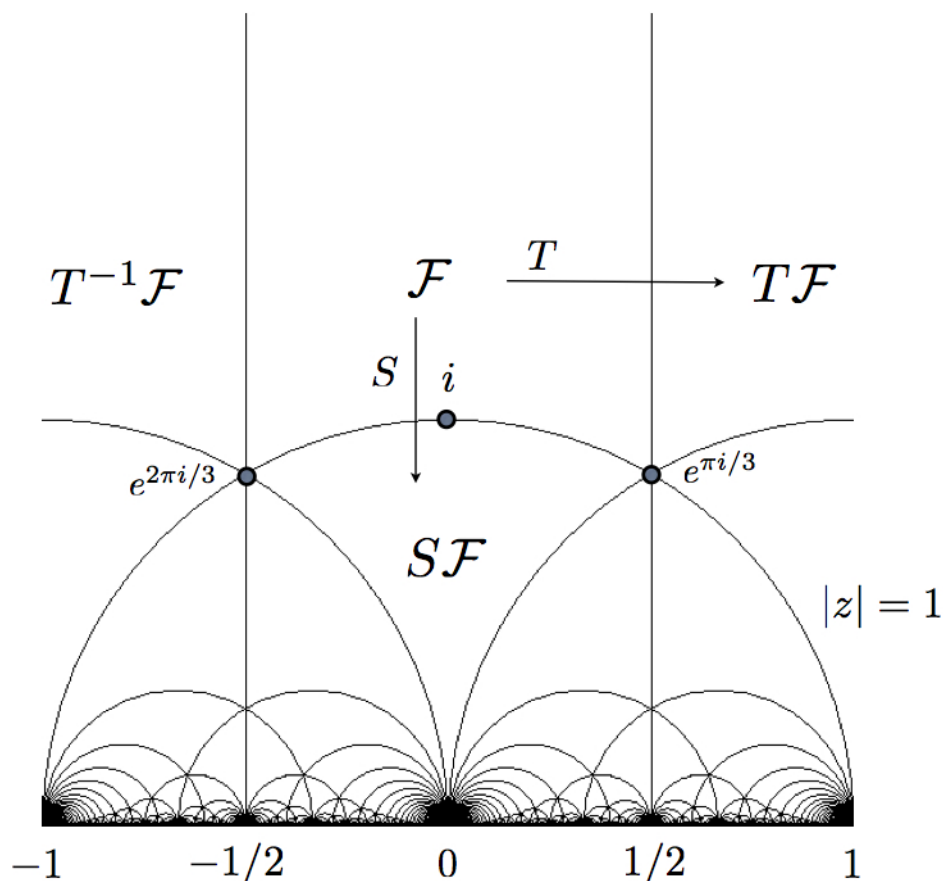


FIGURE 1. The fundamental domain for $\mathrm{SL}_2(\mathbb{Z})$ on \mathbb{H}

for $x \in \mathbb{R}$. We also define supplementary generators for $t \in \mathbb{R}^*$

$$\tilde{w}_\alpha(t) = \chi_\alpha(t)\chi_{-\alpha}(-t^{-1})\chi_\alpha(t) = \begin{pmatrix} 0 & t \\ -t^{-1} & 0 \end{pmatrix},$$

$$\tilde{w}_\alpha = \tilde{w}_\alpha(1) = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad \text{and}$$

$$h_\alpha(t) = \tilde{w}_\alpha^{-1}\tilde{w}_\alpha(t) = \begin{pmatrix} t^{-1} & 0 \\ 0 & t \end{pmatrix}.$$

Define the *toral subgroup*

$$H = \{h_\alpha(u) \mid u \in \mathbb{R}^*\},$$

which is isomorphic to \mathbb{R}^* . Let $X_{\pm\alpha} = \{\chi_{\pm\alpha}(u) \mid u \in \mathbb{R}\}$. Then $X_{\pm\alpha}$ is a subgroup isomorphic to the additive group of \mathbb{R} , called the *root group* associated with $\pm\alpha$. We note that \tilde{w}_α has order 4 and we set $\tilde{W} = \langle \tilde{w}_\alpha \mid \tilde{w}_\alpha^4 = 1 \rangle$. Then \tilde{W} is called the *extended Weyl group*. The *Weyl group*

$$W = \langle w_\alpha \mid w_\alpha^2 = 1 \rangle,$$

which is a symmetry group of the root system, is the quotient $\tilde{W}/\tilde{W} \cap H$.

Lemma 3.1. $\mathrm{SL}_2(\mathbb{R})$ is generated by $\chi_\alpha(u)$, $\chi_{-\alpha}(u)$ for $u \in \mathbb{R}$; with supplementary generators $\tilde{w}_\alpha(t)$, \tilde{w}_α and $h_\alpha(t)$ defined as above for $t \in \mathbb{R}^*$. The Steinberg relations give defining relation for $\mathrm{SL}_2(\mathbb{R})$ with respect to these generators:

$$\begin{aligned} (1) \quad & h_\alpha(uv) = h_\alpha(u)h_\alpha(v), \\ (2) \quad & \tilde{w}_\alpha h_\alpha(t) \tilde{w}_\alpha^{-1} = h_\alpha(t^{-1}), \\ (3) \quad & \chi_\alpha(u+v) = \chi_\alpha(u)\chi_\alpha(v), \\ (4) \quad & \tilde{w}_\alpha(t)\chi_\alpha(u)\tilde{w}_\alpha(t)^{-1} = \chi_{-\alpha}(-t^2u). \end{aligned}$$

Expanding the last relation in terms of χ 's, we get

$$(5) \quad \chi_{-\alpha}(-t^{-1})\chi_\alpha(u)\chi_{-\alpha}(t^{-1}) = \chi_\alpha(-t)\chi_{-\alpha}(-t^2u)\chi_\alpha(t).$$

We note that this relation is required for $\mathrm{SL}_2(\mathbb{R})$, but for higher rank groups it is a consequence of the other Steinberg relations ([St2]).

We can now give an analogous the presentation for $\mathrm{SL}_2(\mathbb{Z})$. The main difference in the defining relations is that $\mathbb{Z}^* = \{\pm 1\}$, which causes the \tilde{w}_α action on the torus (relation (2) above) to be trivial.

Proposition 3.2. $\mathrm{SL}_2(\mathbb{Z})$ is generated by $h_\alpha(-1)$ and $\chi_{\pm\alpha}(n)$ for $n \in \mathbb{Z}$. Defining relations are

$$\begin{aligned} (6) \quad & \chi_{\pm\alpha}(n) = \chi_{\pm\alpha}(1)^n, \\ (7) \quad & \tilde{w}_\alpha \chi_\alpha(\pm 1) \tilde{w}_\alpha^{-1} = \chi_{-\alpha}(\mp 1), \\ (8) \quad & h_\alpha(-1)^2 = 1. \end{aligned}$$

It follows that $h_\alpha(-1) = \tilde{w}_\alpha^2$. We note that for higher rank groups with more than one root, there are additional relations between the root groups, which will not always commute.

We now obtain Iwasawa coordinates for the coset space $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R})$ in terms of these Chevalley generators. We use the Iwasawa decomposition: every $g \in \mathrm{SL}_2(\mathbb{R})$ has a unique representation as $g = k h u$, where

$$k \in K = \mathrm{SO}(2), \quad h \in H^+ = \{h_\alpha(t) \mid t \in \mathbb{R}, t > 0\}, \quad u \in U = X_\alpha.$$

Now $Kg = Kkhu = Khu$, and so the coset $Kg \in K \backslash G$ is uniquely determined by $h = h_\alpha(t^{-1}) \in H^+$ and $u = \chi_\alpha(x) \in U$. So we may represent an element of $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R})$ by its Iwasawa coordinates

$$(9) \quad hu = \begin{pmatrix} t & tx \\ 0 & t^{-1} \end{pmatrix}.$$

We next consider the right action of $\mathrm{SL}_2(\mathbb{Z})$ on $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R})$ in Iwasawa coordinates. For $z \in \mathrm{SL}_2(\mathbb{Z})$, we compute $Kg \cdot z = K h u \cdot z$ and then try to use the relations from the Steinberg presentation to ensure that the result is again in Iwasawa coordinates. It suffices to consider z a generator of $\mathrm{SL}_2(\mathbb{Z})$. It is easy to see that the action of S on $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R}) \cong \mathbb{H}$ has order 2. The element S fixes the coset $\mathrm{SO}(2) \cdot \mathrm{Id}$ corresponding to the point $i \in \mathbb{H}$ and has no other fixed points in $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R})$. The element T acts as a translation on \mathbb{H} and hence has infinite order and acts without fixed points on $\mathrm{SO}(2) \backslash \mathrm{SL}_2(\mathbb{R})$.

We note that $T = \chi_\alpha(1) \in N$ and $S = \tilde{w}_\alpha \in K$, since $\det(S) = 1$ and $SS^T = S^T S = \text{Id}$. The right action of T is

$$Khu \cdot T = K \begin{pmatrix} t & t(x+1) \\ 0 & t^{-1} \end{pmatrix} = Kh_\alpha(t^{-1})\chi_\alpha(x+1),$$

which is in Iwasawa form. The right action of S is

$$Khu \cdot S = K \begin{pmatrix} -tx & t \\ -t^{-1} & 0 \end{pmatrix},$$

which is *not* in Iwasawa form. So we need a *compensating element* $\kappa \in \text{SO}(2)$ such that $\kappa hu \cdot S$ is upper triangular.

Proposition 3.3. *A compensating element for the action of S on $\text{SO}(2) \backslash \text{SL}_2(\mathbb{R})$ is*

$$\kappa = \frac{1}{\sqrt{t^2 x^2 + t^{-2}}} \begin{pmatrix} -tx & -t^{-1} \\ t^{-1} & -tx \end{pmatrix}.$$

In Iwasawa coordinates the action is

$$Kh_\alpha(t^{-1})\chi_\alpha(x) \cdot S = Kh_\alpha\left(1/\sqrt{t^2 x^2 + t^{-2}}\right)\chi_\alpha\left(\frac{-t^2 x}{t^2 x^2 + t^{-2}}\right).$$

Proof. Note that $\det(\kappa) = 1$ and $\kappa\kappa^T = \kappa^T\kappa = \text{Id}$, thus $\kappa \in \text{SO}(2)$. We have

$$\kappa hu \cdot S = \frac{1}{\sqrt{t^2 x^2 + t^{-2}}} \begin{pmatrix} -tx & -t^{-1} \\ t^{-1} & -tx \end{pmatrix} \begin{pmatrix} -tx & t \\ -t^{-1} & 0 \end{pmatrix} = \frac{1}{\sqrt{t^2 x^2 + t^{-2}}} \begin{pmatrix} t^2 x^2 + t^{-2} & -t^2 x \\ 0 & 1 \end{pmatrix},$$

which is in upper triangular form. \square

An example of the right action of S in Iwasawa coordinates is

$$\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \mapsto \begin{pmatrix} \sqrt{2} & 0 \\ 0 & \frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} 1 & -\frac{1}{2} \\ 0 & 1 \end{pmatrix}.$$

The physical significance of the compensating element κ arises from its role as the element that restores the Borel gauge in the dimensional reduction of supergravity on a 2-torus T^2 .

We now write κ in terms of the Steinberg generators. The computer algebra system Magma ([BCFS]) was used to find this formula, which can easily be verified by direct matrix multiplication.

Proposition 3.4. $\kappa = \chi_\alpha(-t^2 x) h_\alpha\left(-1/\sqrt{t^4 x^2 + 1}\right) \tilde{w}_\alpha \chi_\alpha(-t^2 x)$.

We finally remark that the general linear group can be similarly analyzed. In this case, the Steinberg generators for $\text{GL}_2(\mathbb{R})$ are ([CMT])

$$\chi_\alpha(x) = \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad \chi_{-\alpha}(x) = \begin{pmatrix} 1 & 0 \\ x & 1 \end{pmatrix},$$

for $x \in \mathbb{R}$, and also two types of toral generator

$$h_1(t) = \begin{pmatrix} t & 0 \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad h_2(t) = \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix},$$

for $t \in \mathbb{R}^*$.

4. IWASAWA COORDINATES IN A FINITE-DIMENSIONAL LIE GROUP

Let $G = G(\mathbb{R})$ be a split, simply connected, semisimple Lie group of rank n . Let K be the maximal compact subgroup of G . In this section we present our method, using the Steinberg presentation, for computing Iwasawa coordinates of $K \backslash G$. We also compute the action of the generators of the extended Weyl group on $K \backslash G$, since these are the most interesting case as we saw in the last section. When $K \backslash G$ is a scalar coset for a supergravity theory, we apply this method in later sections to find the action of the integral form $G(\mathbb{Z})$ of G on $K \backslash G$.

Let Φ be the root system of G and let Δ be a system of simple roots. The Chevalley generators are $\chi_\alpha(x)$ for $\alpha \in \Phi$ and $x \in \mathbb{R}$. We also define supplementary generators

$$\tilde{w}_\beta(t) = \chi_\beta(t)\chi_{-\beta}(-t^{-1})\chi_\beta(t), \quad \tilde{w}_\beta = \tilde{w}_\beta(1), \quad h_\beta(t) = \tilde{w}_\beta^{-1}\tilde{w}_\beta(t)$$

for $\beta \in \Delta$ a simple root and $t \in \mathbb{R}^\times$. The precise Steinberg relations we are using are from [CMT] based on the method of Carter ([Ca]):

$$(10) \quad h_\beta(t)h_\beta(u) = h_\beta(tu),$$

$$(11) \quad h_\alpha(t)h_\beta(u) = h_\beta(u)h_\alpha(t),$$

$$(12) \quad \tilde{w}_\alpha^{-1}h_\beta(t)\tilde{w}_\alpha = \prod_{\gamma \in \Delta} h_\gamma(t^{a_\gamma}) \quad \text{where } w_\alpha\beta = \sum_{\gamma \in \Delta} a_\gamma\gamma,$$

$$(13) \quad \chi_\alpha(a)\chi_\alpha(b) = \chi_\alpha(a+b),$$

$$(14) \quad \chi_\beta(-b)\chi_\alpha(a)\chi_\beta(b) = \chi_\alpha(a) \prod_{i,j>0} \chi_{i\alpha+j\beta}(C_{ij\alpha\beta}a^i b^j),$$

$$(15) \quad \chi_{-\alpha}(-t)\chi_\alpha(a)\chi_{-\alpha}(t) = \chi_{-\alpha}(-t^2 a)\chi_\alpha(t^{-1}),$$

$$(16) \quad \tilde{w}_\alpha^{-1}\chi_\beta(a)\tilde{w}_\alpha = \chi_{w_\alpha\beta}(\eta_{\alpha\beta}a).$$

The structure constants $C_{ij\alpha\beta}$ and $\eta_{\alpha\beta}$ are as defined in [CMT], with all extraspecial structure constants equal to $(+1)$. Note that equation (15) does not quite agree with the relation (5) given in the last section – this is a consequence of the fact that groups act on the right in Magma.

Let $\Phi^+ = \{\alpha_1, \dots, \alpha_N\}$ denote the positive roots. We list the elements of Φ^+ an order consistent with height, that is, $\text{ht}(\alpha_i) < \text{ht}(\alpha_j)$ implies $i < j$. In particular, the simple roots are $\alpha_1, \dots, \alpha_n$. Write κ_i for κ_{α_i} , h_i for h_{α_i} , χ_r for χ_{α_r} , and \tilde{w}_i for \tilde{w}_{α_i} . Define the root subgroups $X_r = X_{\alpha_r} = \{\chi_r(a) \mid a \in \mathbb{R}\}$. We define the unipotent subgroup $U = \langle X_\alpha \mid \alpha \in \Phi^+ \rangle$, toral subgroup $H = \langle h_\alpha(t) \mid \alpha \in \Delta, t \in \mathbb{R}^* \rangle$, and $H^+ = \langle h_\alpha(t) \mid \alpha \in \Delta, t \in \mathbb{R}, t > 0 \rangle$. Let θ be a Cartan involution on G defined by $\theta(\chi_\alpha(x)) = \chi_{-\alpha}(-1/x)$ for all roots $\alpha \in \Phi$. Then the maximal compact subgroup K of G is the stabilizer of θ . Then we get an Iwasawa decomposition $G = KH^+U$ as in the previous section, and every coset in $K \backslash G$ has Iwasawa coordinates Khu for unique $h \in H^+$ and $u \in U$. We now have

Lemma 4.1. *The Iwasawa coordinates for a coset in $K \backslash G$ are Khu where*

$$h = \prod_{j=1}^n h_j(t_j) \quad \text{and} \quad u = \prod_{r=1}^N \chi_r(x_r),$$

for some $t_j, x_r \in \mathbb{R}$, with $t_j > 0$.

Given a root $\beta \in \Phi$, the reflection w_β has a reduced expression as a product of simple reflections $w_{\beta_1} \cdots w_{\beta_\ell}$, and we define $\tilde{w}_\beta := \tilde{w}_{\beta_1} \cdots \tilde{w}_{\beta_\ell}$. Note that this agrees with the previous definition

for β simple, and that \tilde{w}_β was shown to be independent of the particular reduced expression used in [CMT]. Since G is simply connected, for each $\beta \in \Phi$, the group $\langle \chi_\beta(t), h_\beta(t), \tilde{w}_\beta \mid t \in \mathbb{R}^\times \rangle$ is isomorphic to $\mathrm{SL}_2(\mathbb{R})$. We can therefore use the results of Section 3 for this special case to finding the element $\kappa_\beta(x) \in K$ that makes $\kappa_\beta(x)\chi_\beta(x)\tilde{w}_\beta$ an element of H^+U , for a given root $\beta \in \Phi$.

Lemma 4.2. *If $\kappa_\beta(x) := \chi_\beta(-x)h_\beta(-1/\sqrt{x^2+1})\tilde{w}_\beta\chi_\beta(-x)$, then $\kappa_\beta(x)$ is in the maximal compact subgroup K , and $\kappa_\beta(x)\chi_\beta(x)\tilde{w}_\beta$ is in H^+U .*

Given Iwasawa coordinates hu as in Lemma 4.1 and $\alpha_i \in \Delta$, we now find $\kappa \in K$ such that $\kappa hu\tilde{w}_i \in H^+U$. First we can rewrite $u = \chi_i(x_i)\tilde{u}$, where

$$\tilde{u} = \prod_{r=1}^{i-1} \chi_r(x_r)^{\chi_i(x_i)} \cdot \prod_{r=i+1}^N \chi_r(x_r).$$

Applying the Steinberg commutator relations we get

$$\chi_r(x_r)^{\chi_i(x_i)} = \chi_r(x_r) \prod_{\substack{a,b>0 \text{ s.t.} \\ a\alpha_r+b\alpha_i \in \Phi}} \chi_{a\alpha_r+b\alpha_i}(C_{abri}x_r^a x_i^b),$$

and hence $\tilde{u} \in \prod_{r \neq i} X_r$. We now rearrange

$$hu\tilde{w}_i = h\chi_i(x_i)\tilde{u}\tilde{w}_i = \chi_i(x_i)^{h^{-1}} h\tilde{u}\tilde{w}_i = \chi_i(x_i)^{h^{-1}} \tilde{w}_i(h\tilde{u})^{\tilde{w}_i}.$$

Notice that $(h\tilde{u})^{\tilde{w}_i} \in H^+U$, since $h^{\tilde{w}_i} \in H^+$, and $\tilde{u}^{\tilde{w}_i} \in (\prod_{r \neq i} X_r)^{\tilde{w}_i} = \prod_{r \neq i} X_{w_i\alpha_r} \subseteq U$ because $w_i\alpha_r \in \Phi^+$ when $r \neq i$. Also $\chi_i(x_i)^{h^{-1}} = \chi_i(c_i x_i)$, where $c_i := \prod_{j=1}^n t_j^{-C_{ij}}$ for $C = (C_{ij})$ the Cartan matrix. We can now choose $\kappa = \kappa_i(c_i x_i)$ to ensure that $\kappa hu\tilde{w}_i = \kappa \chi_i(c_i x_i) \tilde{w}_i \cdot (h\tilde{u})^{\tilde{w}_i}$ is indeed in H^+U .

We now give an explicit formula for $\kappa hu\tilde{w}_i$.

Proposition 4.3. *Given $h = \prod_{j=1}^n h_j(t_j) \in H^+$ and $u = \prod_{r=1}^N \chi_r(x_r) \in U$, we have*

$$\kappa hu \cdot \tilde{w}_i = h^{(i)} u^{(i)},$$

where

$$\begin{aligned} h^{(i)} &= h_i(c_i/r_i)h \in H^+, \quad \text{and} \\ u^{(i)} &= \chi_i\left(- (c_i/r_i)^2 x_i\right) \cdot \prod_{r=1}^{i-1} \chi_{w_i\alpha_r}(\eta_{ri}x_i)\tilde{\chi}_{ri} \cdot \prod_{r=i+1}^N \chi_{w_i\alpha_r}(\eta_{ri}x_i) \in U, \end{aligned}$$

for $c_i = \prod_{j=1}^n t_j^{-C_{ij}}$, $r_i = \sqrt{(c_i x_i)^2 + 1}$, and

$$\tilde{\chi}_{ri} = \prod_{\substack{a,b>0 \text{ s.t.} \\ a\alpha_r+b\alpha_i \in \Phi}} \chi_{w_i(a\alpha_r+b\alpha_i)}(\eta_{a\alpha_r+b\alpha_i, \alpha_i} C_{abri}x_r^a x_i^b).$$

Proof. We have

$$\begin{aligned} \kappa hu\tilde{w}_i = \kappa_i(c_i x_i)hu\tilde{w}_i &= \chi_i(-c_i x_i)h_i(-1/r_i)\tilde{w}_i\chi_i(-c_i x_i) \cdot \chi_i(c_i x_i)\tilde{w}_i(h\tilde{u})^{\tilde{w}_i} \\ &= \chi_i(-c_i x_i)h_i(-1/r_i)\tilde{w}_i^2(h\tilde{u})^{\tilde{w}_i} \\ &= \chi_i(-c_i x_i)\left(h_i(-1/r_i)h_i(-1)h^{\tilde{w}_i}\right)\tilde{u}^{\tilde{w}_i} \\ &= h^{(i)}u^{(i)} \end{aligned}$$

where

$$h^{(i)} := h_i(-1/r_i)h_i(-1)h^{\tilde{w}_i} = h_i(-1/r_i)h_i(-1)h_i(c_i)h = h_i(c_i/r_i)h$$

and

$$\begin{aligned} u^{(i)} &:= \chi_i(-c_i x_i) h^{(i)} \tilde{u}^{\tilde{w}_i} \\ &= \chi_i \left((c_i/r_i)^2 \prod_{j=1}^n t_j^{C_{ij}} \cdot (-c_i x_i) \right) \left(\prod_{r=1}^{i-1} \chi_r(x_r)^{\chi_i(x_i)} \cdot \prod_{r=i+1}^N \chi_r(x_r) \right)^{\tilde{w}_i} \\ &= \chi_i \left(-(c_i/r_i)^2 x_i \right) \cdot \prod_{r=1}^{i-1} \chi_r(x_r)^{\chi_i(x_i)\tilde{w}_i} \cdot \prod_{r=i+1}^N \chi_r(x_r)^{\tilde{w}_i} \\ &= \chi_i \left(-(c_i/r_i)^2 x_i \right) \cdot \prod_{r=1}^{i-1} \chi_{w_i \alpha_r}(\eta_{ri} x_i) \tilde{\chi}_{ri} \cdot \prod_{r=i+1}^N \chi_{w_i \alpha_r}(\eta_{ri} x_i). \end{aligned}$$

□

We can find the Iwasawa coordinates for the expression given in this proposition using the symbolic collection algorithm in [CMH]. These coordinates are easy to store in Magma, but too large to list here. Instead we give the uncollected expression for the Iwasawa coordinates for the action of generators of the extended Weyl group on supergravity cosets for $n = 2, \dots, 8$. Since each root system is contained in the previous one, we can fix a single root ordering for all of them. The root indices for each value of n are given in Table 1. Tables 2 to 9 give expressions for $\kappa_i(c_i x_i) h u \tilde{w}_i$ for $i = 1, \dots, n$ in the supergravity coset space of rank n .

i	c_i	$\kappa_i(c_i x_i) h u \tilde{w}_i$
1	$t_2 x_1 / t_1^2$	$h_1(c_1/r_1)h \cdot \chi_i(-c_1/r_1)^2 x_1 \cdot \chi_6(-x_2)\chi_3(x_3)\chi_4(x_4)\chi_5(x_5)\chi_2(x_6)\chi_{10}(-x_7)\chi_8(x_8)\chi_9(x_9)\chi_7(x_{10})$ $\chi_{14}(-x_{11})\chi_{15}(-x_{12})\chi_{13}(x_{13})\chi_{11}(x_{14})\chi_{12}(x_{15})\chi_{17}(-x_{16})\chi_{16}(x_{17})\chi_{19}(-x_{18})\chi_{18}(x_{19})\chi_{20}(x_{20})$
2	$t_1 t_3 x_2 / t_2^2$	$h_2(c_2/r_2)h \cdot \chi_i(-c_2/r_2)^2 x_2 \cdot \chi_1(-x_1 x_2)\chi_6(x_1) \cdot \chi_7(-x_3)\chi_4(x_4)\chi_5(x_5)\chi_1(-x_6)\chi_3(x_7)\chi_{11}(-x_8)\chi_{12}(-x_9)\chi_{10}(x_{10})$ $\chi_8(x_{11})\chi_9(x_{12})\chi_{16}(-x_{13})\chi_{14}(x_{14})\chi_{15}(x_{15})\chi_{13}(x_{16})\chi_{17}(x_{17})\chi_{18}(x_{18})\chi_{20}(-x_{19})\chi_{19}(x_{20})$
3	$t_2 t_4 t_5 x_3 / t_3^2$	$h_3(c_3/r_3)h \cdot \chi_i(-c_3/r_3)^2 x_3 \cdot \chi_1(x_1)\chi_2(-x_2 x_3)\chi_7(x_2) \cdot \chi_8(-x_4)\chi_9(-x_5)\chi_{10}(x_6)\chi_2(-x_7)\chi_4(x_8)\chi_5(x_9)\chi_6(-x_{10})$ $\chi_{11}(x_{11})\chi_{12}(x_{12})\chi_{13}(x_{13})\chi_{14}(x_{14})\chi_{15}(x_{15})\chi_{18}(-x_{16})\chi_{19}(-x_{17})\chi_{16}(x_{18})\chi_{17}(x_{19})\chi_{20}(x_{20})$
4	$t_3 x_4 / t_4^2$	$h_4(c_4/r_4)h \cdot \chi_i(-c_4/r_4)^2 x_4 \cdot \chi_1(x_1)\chi_2(x_2)\chi_3(-x_3 x_4)\chi_8(x_3) \cdot \chi_5(x_5)\chi_6(x_6)\chi_{11}(x_7)\chi_3(-x_8)\chi_{13}(-x_9)\chi_{14}(x_{10})$ $\chi_7(-x_{11})\chi_{16}(-x_{12})\chi_9(x_{13})\chi_{10}(-x_{14})\chi_{17}(-x_{15})\chi_{12}(x_{16})\chi_{15}(x_{17})\chi_{18}(x_{18})\chi_{19}(x_{19})\chi_{20}(x_{20})$
5	$t_3 x_5 / t_5^2$	$h_5(c_5/r_5)h \cdot \chi_i(-c_5/r_5)^2 x_5 \cdot \chi_1(x_1)\chi_2(x_2)\chi_3(-x_3 x_5)\chi_9(x_3)\chi_4(x_4) \cdot \chi_6(x_6)\chi_{12}(x_7)\chi_{13}(-x_8)\chi_3(-x_9)\chi_{15}(x_{10})$ $\chi_{16}(-x_{11})\chi_7(-x_{12})\chi_8(x_{13})\chi_{17}(-x_{14})\chi_{10}(-x_{15})\chi_{11}(x_{16})\chi_{14}(x_{17})\chi_{18}(x_{18})\chi_{19}(x_{19})\chi_{20}(x_{20})$

TABLE 5. Iwasawa coordinates for D_5 with compensating element
 $\kappa_i(c_i x_i) = \chi_i(-c_i x_i) h_\beta(-1/\sqrt{(c_i x_i)^2 + 1}) \tilde{w}_i \chi_i(-c_i x_i)$

i	c_i	$\kappa_i(c_i x_i) h u \tilde{w}_i$
1	$t_3 x_1 / t_1^2$	$h_1(c_1/r_1)h \cdot \chi_i(-c_1/r_1)^2 x_1 \cdot \chi_2(x_2)\chi_7(-x_3)\chi_4(x_4)\chi_5(x_5)\chi_6(x_6)\chi_3(x_7)\chi_8(x_8)\chi_{12}(-x_9)\chi_{10}(x_{10})$ $\chi_{11}(x_{11})\chi_9(x_{12})\chi_{17}(-x_{13})\chi_{14}(x_{14})\chi_{18}(-x_{15})\chi_{16}(x_{16})\chi_{13}(x_{17})\chi_{15}(x_{18})\chi_{22}(-x_{19})\chi_{20}(x_{20})$ $\chi_{23}(-x_{21})\chi_{19}(x_{22})\chi_{21}(x_{23})\chi_{26}(-x_{24})\chi_{27}(-x_{25})\chi_{24}(x_{26})\chi_{25}(x_{27})\chi_{30}(-x_{28})\chi_{29}(x_{29})\chi_{28}(x_{30})$
2	$t_4 x_2 / t_2^2$	$h_2(c_2/r_2)h \cdot \chi_i(-c_2/r_2)^2 x_2 \cdot \chi_1(x_1) \cdot \chi_3(x_3)\chi_8(-x_4)\chi_5(x_5)\chi_6(x_6)\chi_7(x_7)\chi_4(x_8)\chi_{13}(-x_9)\chi_{14}(-x_{10})$ $\chi_{11}(x_{11})\chi_{17}(-x_{12})\chi_9(x_{13})\chi_{10}(x_{14})\chi_{19}(-x_{15})\chi_{20}(-x_{16})\chi_{12}(x_{17})\chi_{22}(-x_{18})\chi_{15}(x_{19})\chi_{16}(x_{20})$ $\chi_{25}(-x_{21})\chi_{18}(x_{22})\chi_{27}(-x_{23})\chi_{24}(x_{24})\chi_{21}(x_{25})\chi_{26}(x_{26})\chi_{23}(x_{27})\chi_{28}(x_{28})\chi_{29}(x_{29})\chi_{30}(x_{30})$
3	$t_1 t_4 x_3 / t_3^2$	$h_3(c_3/r_3)h \cdot \chi_i(-c_3/r_3)^2 x_3 \cdot \chi_1(-x_1 x_3)\chi_7(x_1)\chi_2(x_2) \cdot \chi_9(-x_4)\chi_5(x_5)\chi_6(x_6)\chi_1(-x_7)\chi_{13}(-x_8)\chi_4(x_9)\chi_{15}(-x_{10})$ $\chi_{11}(x_{11})\chi_{12}(x_{12})\chi_8(x_{13})\chi_{19}(-x_{14})\chi_{10}(x_{15})\chi_{21}(-x_{16})\chi_{17}(x_{17})\chi_{18}(x_{18})\chi_{14}(x_{19})\chi_{25}(-x_{20})$ $\chi_{16}(x_{21})\chi_{22}(x_{22})\chi_{23}(x_{23})\chi_{24}(x_{24})\chi_{20}(x_{25})\chi_{29}(-x_{26})\chi_{27}(x_{27})\chi_{28}(x_{28})\chi_{26}(x_{29})\chi_{32}(-x_{30})$
4	$t_2 t_3 t_5 x_4 / t_4^2$	$h_4(c_4/r_4)h \cdot \chi_i(-c_4/r_4)^2 x_4 \cdot \chi_1(x_1)\chi_2(-x_2 x_4)\chi_8(x_2)\chi_3(-x_3 x_4)\chi_9(x_3) \cdot \chi_{10}(-x_5)\chi_6(x_6)\chi_{12}(x_7)\chi_2(-x_8)\chi_3(-x_9)\chi_5(x_{10})$ $\chi_{16}(-x_{11})\chi_7(-x_{12})\chi_{13}(x_{13})\chi_{14}(x_{14})\chi_{15}(x_{15})\chi_{11}(x_{16})\chi_{17}(x_{17})\chi_{18}(x_{18})\chi_{24}(-x_{19})\chi_{20}(x_{20})$ $\chi_{21}(x_{21})\chi_{26}(-x_{22})\chi_{23}(x_{23})\chi_{19}(x_{24})\chi_{28}(-x_{25})\chi_{22}(x_{26})\chi_{30}(-x_{27})\chi_{25}(x_{28})\chi_{29}(x_{29})\chi_{27}(x_{30})$
5	$t_4 t_6 x_5 / t_5^2$	$h_5(c_5/r_5)h \cdot \chi_i(-c_5/r_5)^2 x_5 \cdot \chi_1(x_1)\chi_2(x_2)\chi_3(x_3)\chi_4(-x_4 x_5)\chi_{10}(x_4) \cdot \chi_{11}(-x_6)\chi_7(x_7)\chi_{14}(x_8)\chi_{15}(x_9)\chi_4(-x_{10})$ $\chi_6(x_{11})\chi_{18}(x_{12})\chi_{19}(x_{13})\chi_8(-x_{14})\chi_9(-x_{15})\chi_{16}(x_{16})\chi_{22}(x_{17})\chi_{12}(-x_{18})\chi_{13}(-x_{19})\chi_{20}(x_{20})$ $\chi_{21}(x_{21})\chi_{17}(-x_{22})\chi_{23}(x_{23})\chi_{24}(x_{24})\chi_{25}(x_{25})\chi_{26}(x_{26})\chi_{27}(x_{27})\chi_{31}(-x_{28})\chi_{29}(x_{29})\chi_{33}(-x_{30})$
6	$t_5 x_6 / t_6^2$	$h_6(c_6/r_6)h \cdot \chi_i(-c_6/r_6)^2 x_6 \cdot \chi_1(x_1)\chi_2(x_2)\chi_3(x_3)\chi_4(x_4)\chi_5(-x_5 x_6)\chi_{11}(x_5) \cdot \chi_7(x_7)\chi_8(x_8)\chi_9(x_9)\chi_{16}(x_{10})$ $\chi_5(-x_{11})\chi_{12}(x_{12})\chi_{13}(x_{13})\chi_{20}(x_{14})\chi_{21}(x_{15})\chi_{10}(-x_{16})\chi_{17}(x_{17})\chi_{23}(x_{18})\chi_{25}(x_{19})\chi_{14}(-x_{20})$ $\chi_{15}(-x_{21})\chi_{27}(x_{22})\chi_{18}(-x_{23})\chi_{28}(x_{24})\chi_{19}(-x_{25})\chi_{30}(x_{26})\chi_{22}(-x_{27})\chi_{24}(-x_{28})\chi_{32}(x_{29})\chi_{26}(-x_{30})$

TABLE 6. Iwasawa coordinates for E_6 with compensating element
 $\kappa_i(c_i x_i) = \chi_i(-c_i x_i) h_\beta(-1/\sqrt{(c_i x_i)^2 + 1}) \tilde{w}_i \chi_i(-c_i x_i)$

i	c_i	$\kappa_i(c_i; x_i) hu \tilde{w}_i$
1	t_2/t_1^2	$h_1(c_1/r_1)h \cdot \chi_1(-c_1/r_1)^2 x_1 \cdot \chi_8(-x_2) \chi_3(x_3) \chi_4(x_4) \chi_5(x_5) \chi_6(x_6) \chi_7(x_7) \chi_2(x_8) \chi_9(x_9) \chi_{14}(-x_{10}) \chi_{11}(x_{11}) \chi_{12}(x_{12}) \chi_{13}(x_{13}) \chi_{10}(x_{14}) \chi_{20}(-x_{15}) \chi_{16}(x_{16}) \chi_{21}(-x_{17})$ $\chi_{18}(x_{18}) \chi_{19}(x_{19}) \chi_{15}(x_{20}) \chi_{17}(x_{21}) \chi_{26}(-x_{22}) \chi_{23}(x_{23}) \chi_{27}(-x_{24}) \chi_{35}(x_{35}) \chi_{22}(x_{26}) \chi_{24}(x_{27}) \chi_{32}(-x_{28}) \chi_{33}(-x_{29}) \chi_{30}(x_{30})$ $\chi_{34}(-x_{31}) \chi_{28}(x_{32}) \chi_{29}(x_{33}) \chi_{31}(x_{34}) \chi_{38}(-x_{35}) \chi_{39}(-x_{36}) \chi_{37}(x_{37}) \chi_{35}(x_{38}) \chi_{36}(x_{39}) \chi_{43}(-x_{41}) \chi_{42}(x_{42}) \chi_{40}(x_{43})$ $\chi_{41}(x_{44}) \chi_{48}(-x_{45}) \chi_{46}(x_{46}) \chi_{47}(x_{47}) \chi_{45}(x_{48}) \chi_{52}(-x_{49}) \chi_{50}(x_{50}) \chi_{51}(x_{51}) \chi_{49}(x_{52}) \chi_{53}(x_{53}) \chi_{54}(x_{54}) \chi_{55}(x_{55}) \chi_{56}(x_{56})$ $\chi_{57}(x_{57}) \chi_{58}(x_{58}) \chi_{59}(x_{59}) \chi_{60}(x_{60}) \chi_{61}(x_{61}) \chi_{63}(-x_{62}) \chi_{62}(x_{63})$
2	$t_1 t_4/t_2^2$	$h_2(c_2/r_2)h \cdot \chi_2(-x_1) \chi_2(x_2) \chi_1 \cdot \chi_3(x_3) \chi_{10}(-x_4) \chi_5(x_5) \chi_6(x_6) \chi_7(x_7) \chi_{11}(-x_8) \chi_{15}(-x_9) \chi_4(x_{10}) \chi_{17}(-x_{11}) \chi_{12}(x_{12}) \chi_{13}(x_{13}) \chi_{14}(x_{14}) \chi_9(x_{15})$ $x_{22}(x_{22}) \chi_{46}(-x_{43}) \chi_{47}(-x_{44}) \chi_{45}(x_{45}) \chi_{43}(x_{46}) \chi_{44}(x_{47}) \chi_{51}(-x_{48}) \chi_{49}(x_{49}) \chi_{50}(x_{50}) \chi_{48}(x_{51}) \chi_{55}(-x_{52}) \chi_{53}(x_{53}) \chi_{54}(x_{54}) \chi_{52}(x_{55}) \chi_{56}(x_{56}) \chi_{57}(x_{57})$ $\chi_{58}(x_{58}) \chi_{59}(x_{59}) \chi_{60}(x_{60}) \chi_{62}(-x_{61}) \chi_{61}(x_{62}) \chi_{63}(x_{63})$
3	t_4/t_3^2	$h_3(c_3/r_3)h \cdot \chi_3(-c_3/r_3)^2 x_3 \cdot \chi_1(x_1) \chi_2(x_2) \cdot \chi_9(-x_4) \chi_5(x_5) \chi_6(x_6) \chi_7(x_7) \chi_8(x_8) \chi_4(x_9) \chi_{15}(-x_{10}) \chi_{16}(-x_{11}) \chi_{12}(x_{12}) \chi_{13}(x_{13}) \chi_{20}(-x_{14}) \chi_{10}(x_{15}) \chi_{11}(x_{16}) \chi_{22}(-x_{17})$ $\chi_{23}(-x_{18}) \chi_{19}(x_{19}) \chi_{14}(x_{20}) \chi_{26}(-x_{21}) \chi_{17}(x_{22}) \chi_{18}(x_{23}) \chi_{29}(-x_{24}) \chi_{30}(-x_{25}) \chi_{21}(x_{26}) \chi_{33}(-x_{27}) \chi_{28}(x_{28}) \chi_{24}(x_{29}) \chi_{25}(x_{30}) \chi_{36}(-x_{31}) \chi_{32}(x_{32}) \chi_{27}(x_{33}) \chi_{39}(-x_{34})$ $\chi_{35}(x_{35}) \chi_{31}(x_{36}) \chi_{37}(x_{37}) \chi_{38}(x_{38}) \chi_{34}(x_{39}) \chi_{40}(x_{40}) \chi_{41}(x_{41}) \chi_{42}(x_{42}) \chi_{43}(x_{43}) \chi_{44}(x_{44}) \chi_{45}(x_{45}) \chi_{46}(x_{46}) \chi_{47}(x_{47}) \chi_{48}(x_{48}) \chi_{49}(x_{49}) \chi_{55}(-x_{50})$ $\chi_{51}(x_{51}) \chi_{52}(x_{52}) \chi_{50}(x_{53}) \chi_{56}(-x_{54}) \chi_{55}(x_{55}) \chi_{54}(x_{56}) \chi_{58}(-x_{57}) \chi_{57}(x_{58}) \chi_{60}(-x_{59}) \chi_{59}(x_{60}) \chi_{61}(x_{61}) \chi_{62}(x_{62}) \chi_{63}(x_{63})$ $h_4(c_4/r_4)h \cdot \chi_4(-c_4/r_4)^2 x_4 \cdot \chi_1(x_1) \chi_2(x_2) \chi_3(-x_3) \chi_4 \chi_9(x_5) \cdot \chi_{11}(-x_5) \chi_6(x_6) \chi_7(x_7) \chi_{14}(x_8) \chi_3(-x_9) \chi_2(-x_{10}) \chi_5(x_{11}) \chi_{18}(x_{13}) \chi_8(x_{14})$ $\chi_{15}(x_{15}) \chi_{16}(x_{16}) \chi_{17}(x_{17}) \chi_{12}(x_{18}) \chi_{25}(-x_{19}) \chi_{20}(x_{20}) \chi_{21}(x_{21}) \chi_{28}(-x_{22}) \chi_{23}(x_{23}) \chi_{24}(x_{24}) \chi_{19}(x_{25}) \chi_{32}(-x_{26}) \chi_{27}(x_{27}) \chi_{22}(x_{28}) \chi_{35}(-x_{29}) \chi_{30}(x_{30})$ $\chi_{31}(x_{31}) \chi_{26}(x_{32}) \chi_{38}(-x_{33}) \chi_{34}(x_{34}) \chi_{29}(x_{35}) \chi_{41}(-x_{36}) \chi_{37}(x_{37}) \chi_{33}(x_{38}) \chi_{44}(-x_{39}) \chi_{40}(x_{40}) \chi_{36}(x_{41}) \chi_{42}(x_{43}) \chi_{39}(x_{44}) \chi_{45}(x_{45}) \chi_{46}(x_{46}) \chi_{47}(x_{47}) \chi_{48}(x_{48}) \chi_{49}(x_{49}) \chi_{50}(-x_{46})$ $\chi_{47}(x_{47}) \chi_{48}(x_{48}) \chi_{49}(x_{49}) \chi_{46}(x_{50}) \chi_{54}(-x_{51}) \chi_{52}(x_{52}) \chi_{53}(x_{53}) \chi_{51}(x_{54}) \chi_{57}(-x_{55}) \chi_{56}(x_{56})$ $\chi_{55}(x_{57}) \chi_{58}(x_{58}) \chi_{59}(x_{59}) \chi_{61}(-x_{60}) \chi_{60}(x_{61}) \chi_{62}(x_{62}) \chi_{63}(x_{63})$
4	$t_2 t_3 t_5/t_4^2$	$h_5(c_5/r_5)h \cdot \chi_5(-c_5/r_5)^2 x_5 \cdot \chi_1(x_1) \chi_2(x_2) \chi_3(x_3) \chi_4(-x_4) \chi_5 \chi_{11}(x_{11}) \cdot \chi_{12}(-x_6) \chi_7(x_7) \chi_8(x_8) \chi_{16}(x_9) \chi_{17}(x_{10}) \chi_4(-x_{11}) \chi_6(x_{12}) \chi_{19}(-x_{13}) \chi_{21}(x_{14}) \chi_{22}(x_{15}) \chi_9(-x_{16})$ $\chi_{10}(-x_{17}) \chi_{18}(x_{18}) \chi_{13}(x_{19}) \chi_{26}(x_{20}) \chi_{14}(-x_{21}) \chi_{15}(-x_{22}) \chi_{23}(x_{23}) \chi_{24}(x_{24}) \chi_{25}(x_{25}) \chi_{20}(-x_{26}) \chi_{27}(x_{27}) \chi_{28}(x_{28}) \chi_{29}(x_{29}) \chi_{30}(x_{30}) \chi_{31}(x_{31}) \chi_{32}(x_{32})$ $\chi_{33}(x_{33}) \chi_{34}(x_{34}) \chi_{40}(-x_{35}) \chi_{36}(x_{36}) \chi_{37}(x_{37}) \chi_{43}(-x_{38}) \chi_{39}(x_{39}) \chi_{35}(x_{40}) \chi_{45}(-x_{41}) \chi_{46}(-x_{42}) \chi_{38}(x_{43}) \chi_{48}(-x_{44}) \chi_{41}(x_{45}) \chi_{42}(x_{46}) \chi_{51}(-x_{47}) \chi_{44}(x_{48})$ $h_6(c_6/r_6)h \cdot \chi_6(-c_6/r_6)^2 x_6 \cdot \chi_1(x_1) \chi_2(x_2) \chi_3(x_3) \chi_4(x_4) \chi_5(-x_5) \chi_6(x_6) \chi_{12}(x_{12}) \cdot \chi_{13}(-x_7) \chi_8(x_8) \chi_9(x_9) \chi_{10}(x_{10}) \chi_{18}(x_{11}) \chi_5(-x_{12}) \chi_7(x_{13}) \chi_{14}(x_{14}) \chi_{15}(x_{15}) \chi_{23}(x_{16})$ $\chi_{24}(x_{17}) \chi_{11}(-x_{18}) \chi_{19}(x_{19}) \chi_{20}(x_{20}) \chi_{27}(x_{21}) \chi_{29}(x_{22}) \chi_{16}(-x_{23}) \chi_{17}(-x_{24}) \chi_{25}(x_{25}) \chi_{33}(x_{26}) \chi_{21}(-x_{27}) \chi_{35}(x_{28}) \chi_{22}(-x_{29}) \chi_{30}(x_{30}) \chi_{31}(x_{31}) \chi_{38}(x_{32})$ $\chi_{26}(-x_{33}) \chi_{34}(x_{34}) \chi_{28}(-x_{35}) \chi_{36}(x_{36}) \chi_{42}(x_{37}) \chi_{32}(-x_{38}) \chi_{39}(x_{39}) \chi_{40}(x_{40}) \chi_{41}(x_{41}) \chi_{37}(-x_{42}) \chi_{43}(x_{43}) \chi_{44}(x_{44}) \chi_{49}(-x_{45}) \chi_{46}(x_{46}) \chi_{47}(x_{47}) \chi_{48}(x_{48})$ $\chi_{45}(x_{49}) \chi_{50}(x_{50}) \chi_{55}(-x_{51}) \chi_{48}(x_{52}) \chi_{53}(x_{53}) \chi_{57}(-x_{54}) \chi_{51}(x_{55}) \chi_{58}(-x_{56}) \chi_{54}(x_{57}) \chi_{56}(x_{58}) \chi_{59}(x_{59}) \chi_{60}(x_{60}) \chi_{61}(x_{61}) \chi_{62}(x_{62}) \chi_{63}(x_{63})$ $h_7(c_7/r_7)h \cdot \chi_7(-c_7/r_7)^2 x_7 \cdot \chi_1(x_1) \chi_2(x_2) \chi_3(x_3) \chi_4(x_4) \chi_5(x_5) \chi_6(-x_6) \chi_7(x_7) \cdot \chi_8(x_8) \chi_9(x_9) \chi_{10}(x_{10}) \chi_{11}(x_{11}) \chi_{19}(x_{12}) \chi_6(-x_{13}) \chi_{14}(x_{14}) \chi_{15}(x_{15}) \chi_{16}(x_{16})$ $\chi_{17}(x_{17}) \chi_{25}(x_{18}) \chi_{12}(-x_{19}) \chi_{20}(x_{20}) \chi_{21}(x_{21}) \chi_{22}(x_{22}) \chi_{30}(x_{23}) \chi_{31}(x_{24}) \chi_{18}(-x_{25}) \chi_{26}(x_{26}) \chi_{34}(x_{27}) \chi_{28}(x_{28}) \chi_{36}(x_{29}) \chi_{23}(-x_{30}) \chi_{24}(-x_{31}) \chi_{32}(x_{32})$ $\chi_{39}(x_{33}) \chi_{27}(-x_{34}) \chi_{41}(x_{35}) \chi_{29}(-x_{36}) \chi_{37}(x_{37}) \chi_{44}(x_{38}) \chi_{33}(-x_{39}) \chi_{45}(x_{40}) \chi_{35}(-x_{41}) \chi_{47}(x_{42}) \chi_{48}(x_{43}) \chi_{38}(-x_{44}) \chi_{40}(-x_{45}) \chi_{51}(x_{46}) \chi_{42}(-x_{47}) \chi_{43}(-x_{48})$ $\chi_{49}(x_{49}) \chi_{54}(x_{50}) \chi_{46}(-x_{51}) \chi_{52}(x_{52}) \chi_{56}(x_{53}) \chi_{60}(-x_{54}) \chi_{55}(x_{55}) \chi_{63}(-x_{56}) \chi_{57}(x_{57}) \chi_{58}(x_{58}) \chi_{59}(x_{59}) \chi_{60}(x_{60}) \chi_{61}(x_{61}) \chi_{62}(x_{62}) \chi_{63}(x_{63})$

TABLE 7. Iwasawa coordinates for E_7 with compensating element $\kappa_i(c_i; x_i) = \chi_i(-c_i; x_i) h_i(-1/\sqrt{c_i; x_i})^2 + 1 \tilde{w}_i \chi_i(-c_i; x_i)$

i	c_i	$\kappa_i(c_i; x_i) hu \tilde{u}_i$
1	t_2/t_1^2	$h_1(c_1/r_1)h \cdot X_1 \cdot X_9(-x_2)X_3(x_3)X_4(x_4)X_5(x_5)X_6(x_6)X_7(x_7)X_8(x_8)X_2(x_9)X_{10}(x_{10})X_{16}(-x_{11})X_{12}(x_{12})X_{13}(x_{13})X_{14}(x_{14})X_{15}(x_{15})X_{11}(x_{16})X_{23}(-x_{17})$ $X_{18}(x_{18})X_{24}(-x_{19})X_{20}(x_{20})X_{21}(x_{21})X_{22}(x_{22})X_{17}(x_{23})X_{19}(x_{24})X_{30}(-x_{25})X_{26}(x_{26})X_{31}(-x_{27})X_{28}(x_{28})X_{29}(x_{29})X_{25}(x_{30})X_{27}(x_{31})X_{37}(-x_{32})X_{38}(-x_{33})$ $X_{34}(x_{34})X_{39}(-x_{35})X_{36}(x_{36})X_{32}(x_{37})X_{33}(x_{38})X_{35}(x_{39})X_{45}(-x_{40})X_{42}(x_{42})X_{47}(-x_{43})X_{44}(x_{44})X_{40}(x_{45})X_{41}(x_{46})X_{43}(x_{47})X_{52}(-x_{48})X_{53}(-x_{49})$ $X_{54}(-x_{50})X_{51}(x_{51})X_{48}(x_{52})X_{49}(x_{53})X_{50}(x_{54})X_{59}(-x_{55})X_{57}(x_{57})X_{58}(x_{58})X_{55}(x_{59})X_{56}(x_{60})X_{66}(-x_{61})X_{67}(-x_{62})X_{63}(x_{63})X_{64}(x_{64})X_{65}(x_{65})$ $X_{61}(x_{66})X_{62}(x_{67})X_{73}(-x_{68})X_{69}(x_{69})X_{70}(x_{70})X_{71}(x_{71})X_{72}(x_{72})X_{68}(x_{73})X_{79}(-x_{74})X_{75}(x_{75})X_{76}(x_{76})X_{77}(x_{77})X_{78}(x_{78})X_{74}(x_{79})X_{80}(x_{80})X_{81}(x_{81})$ $X_{82}(x_{82})X_{83}(x_{83})X_{84}(x_{84})X_{85}(x_{85})X_{86}(x_{86})X_{87}(x_{87})X_{88}(x_{88})X_{89}(x_{89})X_{90}(x_{90})X_{92}(x_{92})X_{97}(-x_{93})X_{94}(x_{94})X_{95}(x_{95})X_{96}(x_{96})X_{93}(x_{97})$ $X_{101}(-x_{98})X_{99}(x_{99})X_{100}(x_{100})X_{98}(x_{101})X_{104}(-x_{102})X_{103}(x_{103})X_{102}(x_{104})X_{107}(-x_{105})X_{105}(x_{105})X_{106}(x_{106})X_{108}(x_{108})X_{111}(-x_{110})X_{110}(x_{111})X_{113}(-x_{112})$ $X_{112}(x_{113})X_{114}(x_{114})X_{115}(x_{115})X_{116}(x_{116})X_{117}(x_{117})X_{118}(x_{118})X_{119}(x_{119})X_{120}(x_{120})$
2	$t_1 t_4/t_2^2$	$h_2(c_2/r_2)h \cdot X_2(-x_2/r_2)^2 X_2 \cdot X_1(-x_1, x_2)X_9(x_1) \cdot X_3(x_3)X_{11}(-x_4)X_5(x_5)X_6(x_6)X_7(x_7)X_8(x_8)X_1(-x_9)X_{17}(x_{17})X_{18}(x_{18})X_{19}(x_{19})X_{13}(x_{13})X_{14}(x_{14})X_{15}(x_{15})$ $X_{16}(x_{16})X_{10}(x_{17})X_{25}(-x_{18})X_{12}(x_{19})X_{27}(-x_{20})X_{21}(x_{21})X_{22}(x_{22})X_{23}(x_{23})X_{24}(x_{24})X_{18}(x_{25})X_{33}(-x_{26})X_{20}(x_{27})X_{35}(-x_{28})X_{29}(x_{29})X_{30}(x_{30})X_{31}(x_{31})$ $X_{32}(x_{32})X_{26}(x_{33})X_{41}(-x_{34})X_{28}(x_{35})X_{43}(-x_{36})X_{44}(-x_{37})X_{38}(x_{38})X_{39}(x_{39})X_{40}(x_{40})X_{34}(x_{41})X_{50}(-x_{42})X_{36}(x_{43})X_{37}(x_{44})X_{51}(-x_{45})X_{46}(x_{46})X_{47}(x_{47})$ $X_{48}(x_{48})X_{49}(x_{49})X_{42}(x_{50})X_{45}(x_{51})X_{57}(-x_{52})X_{58}(-x_{53})X_{54}(x_{54})X_{55}(x_{55})X_{56}(x_{56})X_{52}(x_{57})X_{63}(x_{58})X_{61}(x_{61})X_{62}(x_{62})X_{63}(x_{63})$ $X_{59}(x_{64})X_{60}(x_{65})X_{71}(-x_{66})X_{72}(-x_{67})X_{68}(x_{68})X_{69}(x_{69})X_{70}(x_{70})X_{66}(x_{71})X_{67}(x_{72})X_{78}(-x_{73})X_{74}(x_{74})X_{75}(x_{75})X_{76}(x_{76})X_{77}(x_{77})X_{73}(x_{78})X_{84}(-x_{79})$ $X_{80}(x_{80})X_{81}(x_{81})X_{82}(x_{82})X_{83}(x_{83})X_{79}(x_{84})X_{85}(x_{85})X_{86}(x_{86})X_{87}(x_{87})X_{88}(x_{88})X_{93}(-x_{89})X_{90}(x_{90})X_{91}(x_{91})X_{92}(x_{92})X_{89}(x_{93})X_{98}(-x_{94})X_{95}(x_{95})$ $X_{96}(x_{96})X_{97}(x_{97})X_{94}(x_{98})X_{102}(-x_{99})X_{100}(x_{100})X_{101}(x_{101})X_{99}(x_{102})X_{105}(-x_{103})X_{104}(x_{104})X_{103}(x_{105})X_{108}(x_{108})X_{109}(x_{109})X_{110}(x_{110})X_{114}(-x_{111})$ $X_{112}(x_{112})X_{115}(-x_{113})X_{111}(x_{114})X_{113}(x_{115})X_{116}(x_{116})X_{117}(x_{117})X_{118}(x_{118})X_{119}(x_{119})X_{120}(x_{120})$
3	t_4/t_3^2	$h_3(c_3/r_3)h \cdot X_3(-x_3/r_3)^2 X_3 \cdot X_1(x_1)X_2(x_2) \cdot X_{10}(-x_4)X_5(x_5)X_6(x_6)X_7(x_7)X_8(x_8)X_9(x_9)X_4(x_{10})X_{17}(-x_{11})X_{18}(-x_{12})X_{13}(x_{13})X_{14}(x_{14})X_{15}(x_{15})X_{23}(-x_{16})X_{11}(x_{17})$ $X_{12}(x_{18})X_{25}(-x_{19})X_{26}(-x_{20})X_{21}(x_{21})X_{22}(x_{22})X_{16}(x_{23})X_{30}(-x_{24})X_{19}(x_{25})X_{20}(x_{26})X_{33}(-x_{27})X_{34}(-x_{28})X_{29}(x_{29})X_{32}(x_{32})X_{27}(x_{33})$ $X_{28}(x_{34})X_{41}(-x_{35})X_{42}(-x_{36})X_{37}(x_{37})X_{31}(x_{38})X_{46}(-x_{39})X_{40}(x_{40})X_{35}(x_{41})X_{36}(x_{42})X_{50}(-x_{43})X_{44}(x_{44})X_{45}(x_{45})X_{39}(x_{46})X_{54}(-x_{47})X_{48}(x_{48})X_{49}(x_{49})$ $X_{43}(x_{50})X_{51}(x_{51})X_{52}(x_{52})X_{53}(x_{53})X_{47}(x_{54})X_{55}(x_{55})X_{56}(x_{56})X_{57}(x_{57})X_{58}(x_{58})X_{59}(x_{59})X_{60}(x_{60})X_{61}(x_{61})X_{62}(x_{62})X_{69}(-x_{63})X_{64}(x_{64})X_{65}(x_{65})$ $X_{66}(x_{66})X_{67}(x_{67})X_{68}(x_{68})X_{63}(x_{69})X_{75}(-x_{70})X_{71}(x_{71})X_{72}(x_{72})X_{73}(x_{73})X_{74}(x_{74})X_{70}(x_{75})X_{80}(-x_{77})X_{78}(x_{78})X_{79}(x_{79})X_{76}(x_{80})X_{77}(x_{81})$ $X_{85}(-x_{82})X_{86}(-x_{83})X_{84}(x_{84})X_{82}(x_{85})X_{83}(x_{86})X_{90}(-x_{87})X_{91}(-x_{88})X_{89}(x_{89})X_{87}(x_{90})X_{88}(x_{91})X_{95}(-x_{92})X_{93}(x_{93})X_{94}(x_{94})X_{92}(x_{95})X_{100}(-x_{96})X_{97}(x_{97})$ $X_{98}(x_{98})X_{99}(x_{99})X_{96}(x_{100})X_{101}(x_{101})X_{102}(x_{102})X_{103}(x_{103})X_{104}(x_{104})X_{105}(x_{105})X_{106}(x_{106})X_{107}(x_{107})X_{108}(x_{108})X_{109}(x_{109})X_{112}(-x_{110})X_{111}(x_{111})X_{110}(x_{112})$ $X_{111}(x_{113})X_{115}(-x_{114})X_{114}(x_{116})X_{116}(x_{116})X_{117}(x_{117})X_{118}(x_{118})X_{119}(x_{119})X_{120}(x_{120})$
4	$t_2 t_3 t_5/t_4^2$	$h_4(c_4/r_4)h \cdot X_4(-x_4/r_4)^2 X_4 \cdot X_1(x_1)X_2(x_2)X_3(-x_3, x_4)X_{11}(x_2)X_3 \cdot X_{12}(-x_5)X_6(x_6)X_7(x_7)X_8(x_8)X_{16}(x_9)X_3(-x_{10})X_2(-x_{11})X_5(x_{12})X_{20}(-x_{13})X_{14}(x_{14})$ $X_{15}(x_{15})X_9(-x_{16})X_{17}(x_{17})X_{18}(x_{18})X_{19}(x_{19})X_{13}(x_{20})X_{28}(-x_{21})X_{22}(x_{22})X_{23}(x_{23})X_{24}(x_{24})X_{32}(-x_{25})X_{26}(x_{26})X_{27}(x_{27})X_{21}(x_{28})X_{36}(-x_{29})$ $X_{37}(-x_{30})X_{31}(x_{31})X_{25}(x_{32})X_{40}(-x_{33})X_{34}(x_{34})X_{35}(x_{35})X_{29}(x_{36})X_{30}(x_{37})X_{45}(-x_{38})X_{39}(x_{39})X_{33}(x_{40})X_{49}(-x_{41})X_{42}(x_{42})X_{43}(x_{43})X_{44}(x_{44})$ $X_{38}(x_{45})X_{53}(-x_{46})X_{47}(x_{47})X_{48}(x_{48})X_{41}(x_{49})X_{56}(-x_{50})X_{51}(x_{51})X_{52}(x_{52})X_{46}(x_{53})X_{60}(-x_{54})X_{55}(x_{55})X_{50}(x_{56})X_{63}(-x_{57})X_{58}(x_{58})X_{59}(x_{59})X_{54}(x_{60})$ $X_{61}(x_{61})X_{62}(x_{62})X_{57}(x_{63})X_{70}(-x_{64})X_{65}(x_{65})X_{66}(x_{66})X_{67}(x_{67})X_{68}(x_{68})X_{69}(x_{69})X_{76}(-x_{71})X_{77}(-x_{72})X_{73}(x_{73})X_{74}(x_{74})$ $X_{75}(x_{75})X_{71}(x_{76})X_{72}(x_{77})X_{83}(-x_{78})X_{79}(x_{79})X_{80}(x_{80})X_{81}(x_{81})X_{82}(x_{82})X_{78}(x_{83})X_{87}(x_{87})X_{84}(x_{88})X_{85}(x_{89})X_{94}(-x_{90})X_{91}(x_{91})X_{92}(x_{92})X_{93}(x_{93})X_{90}(x_{94})X_{99}(-x_{95})$ $X_{96}(x_{96})X_{97}(x_{97})X_{98}(x_{98})X_{95}(x_{99})X_{103}(-x_{100})X_{101}(x_{101})X_{102}(x_{102})X_{100}(x_{103})X_{104}(x_{104})X_{105}(x_{105})X_{106}(x_{106})X_{107}(x_{107})X_{108}(x_{108})X_{109}(x_{109})X_{112}(-x_{110})X_{111}(x_{111})X_{110}(x_{112})$ $X_{107}(x_{107})X_{110}(-x_{108})X_{111}(-x_{109})X_{108}(x_{110})X_{110}(x_{110})X_{109}(x_{111})X_{114}(x_{114})X_{113}(x_{113})X_{112}(x_{112})X_{115}(x_{115})X_{116}(x_{116})X_{117}(x_{117})X_{118}(x_{118})X_{119}(x_{119})X_{120}(x_{120})$

TABLE 8. Iwasawa coordinates for E_8 with compensating element $\kappa_i(c_i; x_i) = X_i(-c_i; x_i)h_i(-1/\sqrt{(c_i; x_i)^2 + 1})\tilde{w}_i X_i(-c_i; x_i)$: Part 1

i	c_i	$\kappa_i(c_i, x_i) h u \tilde{w}_i$
5	$t_4 t_6 / t_5^2$	$h_5 (c_5 / r_5) h \cdot X_5 (- (c_5 / r_5)^2 x_5 \cdot X_1 (x_1) X_2 (x_2) X_3 (x_3) X_4 (-x_4 x_5^2) X_{12} (x_{12}) \cdot X_{13} (-x_6) X_7 (x_7) X_8 (x_8) X_9 (x_9) X_{18} (x_{18}) X_{19} (x_{19}) X_{21} (x_{21}) X_{24} (x_{24}) X_{25} (x_{25}) X_{26} (x_{26}) X_{27} (x_{27}) X_{28} (x_{28}) X_{29} (x_{29}) X_{30} (x_{30}) X_{31} (x_{31}) X_{32} (x_{32}) X_{33} (x_{33}) X_{34} (x_{34}) X_{35} (x_{35}) X_{36} (x_{36}) X_{37} (x_{37}) X_{38} (x_{38}) X_{39} (x_{39}) X_{48} (-x_{40}) X_{41} (x_{41}) X_{42} (x_{42}) X_{43} (x_{43}) X_{44} (x_{44}) X_{45} (-x_{46}) X_{46} (x_{46}) X_{47} (x_{47}) X_{48} (x_{48}) X_{55} (-x_{49}) X_{50} (x_{50}) X_{57} (-x_{51}) X_{45} (x_{52}) X_{59} (-x_{53}) X_{54} (x_{54}) X_{49} (x_{55}) X_{62} (-x_{56}) X_{51} (x_{56}) X_{64} (-x_{58}) X_{53} (x_{59}) X_{67} (-x_{60}) X_{61} (x_{61}) X_{56} (x_{62}) X_{63} (x_{63}) X_{58} (x_{64}) X_{72} (-x_{65}) X_{66} (x_{66}) X_{60} (x_{67}) X_{68} (x_{68}) X_{69} (x_{69}) X_{70} (x_{70}) X_{71} (x_{71}) X_{65} (x_{72}) X_{73} (x_{73}) X_{74} (x_{74}) X_{75} (x_{75}) X_{82} (-x_{76}) X_{77} (x_{77}) X_{78} (x_{78}) X_{79} (x_{79}) X_{85} (-x_{80}) X_{81} (x_{81}) X_{76} (x_{82}) X_{87} (-x_{83}) X_{84} (x_{84}) X_{80} (x_{85}) X_{90} (-x_{86}) X_{83} (x_{87}) X_{92} (-x_{88}) X_{89} (x_{89}) X_{86} (x_{90}) X_{95} (-x_{91}) X_{88} (x_{92}) X_{93} (x_{93}) X_{94} (x_{94}) X_{91} (x_{95}) X_{96} (x_{96}) X_{97} (x_{97}) X_{98} (x_{98}) X_{99} (x_{99}) X_{100} (x_{100}) X_{101} (x_{101}) X_{102} (x_{102}) X_{103} (x_{103}) X_{104} (x_{104}) X_{108} (-x_{105}) X_{103} (x_{106}) X_{109} (-x_{107}) X_{105} (x_{108}) X_{107} (x_{109}) X_{110} (x_{110}) X_{111} (x_{111}) X_{112} (x_{112}) X_{113} (x_{113}) X_{114} (x_{114}) X_{115} (x_{115}) X_{117} (-x_{116}) X_{116} (x_{117}) X_{118} (x_{118}) X_{119} (x_{119}) X_{120} (x_{120})$
6	$t_5 t_7 / t_6^2$	$h_6 (c_6 / r_6) h \cdot X_6 (- (c_6 / r_6)^2 x_6 \cdot X_1 (x_1) X_2 (x_2) X_3 (x_3) X_4 (x_4) X_5 (-x_5 x_6^2) X_{13} (x_{13}) \cdot X_{14} (-x_7) X_8 (x_8) X_9 (x_9) X_{10} (x_{10}) X_{11} (x_{11}) X_{12} (x_{12}) X_{21} (x_{21}) X_{25} (-x_{13}) X_{27} (x_{27}) X_{31} (x_{31}) X_{32} (x_{32}) X_{33} (x_{33}) X_{34} (x_{34}) X_{35} (x_{35}) X_{36} (x_{36}) X_{37} (x_{37}) X_{38} (x_{38}) X_{39} (x_{39}) X_{48} (-x_{40}) X_{41} (x_{41}) X_{42} (x_{42}) X_{43} (x_{43}) X_{44} (x_{44}) X_{45} (-x_{46}) X_{46} (x_{46}) X_{47} (x_{47}) X_{48} (x_{48}) X_{55} (-x_{49}) X_{50} (x_{50}) X_{57} (-x_{51}) X_{45} (x_{52}) X_{59} (-x_{53}) X_{54} (x_{54}) X_{49} (x_{55}) X_{62} (-x_{56}) X_{51} (x_{56}) X_{64} (-x_{58}) X_{53} (x_{59}) X_{67} (-x_{60}) X_{61} (x_{61}) X_{56} (x_{62}) X_{63} (x_{63}) X_{58} (x_{64}) X_{72} (-x_{65}) X_{66} (x_{66}) X_{60} (x_{67}) X_{68} (x_{68}) X_{69} (x_{69}) X_{70} (x_{70}) X_{71} (x_{71}) X_{65} (x_{72}) X_{73} (x_{73}) X_{74} (x_{74}) X_{75} (x_{75}) X_{82} (-x_{76}) X_{77} (x_{77}) X_{78} (x_{78}) X_{79} (x_{79}) X_{85} (-x_{80}) X_{81} (x_{81}) X_{76} (x_{82}) X_{87} (-x_{83}) X_{84} (x_{84}) X_{80} (x_{85}) X_{90} (-x_{86}) X_{83} (x_{87}) X_{92} (-x_{88}) X_{89} (x_{89}) X_{86} (x_{90}) X_{95} (-x_{91}) X_{88} (x_{92}) X_{93} (x_{93}) X_{94} (x_{94}) X_{91} (x_{95}) X_{96} (x_{96}) X_{97} (x_{97}) X_{98} (x_{98}) X_{99} (x_{99}) X_{100} (x_{100}) X_{101} (x_{101}) X_{102} (x_{102}) X_{103} (x_{103}) X_{104} (x_{104}) X_{108} (-x_{105}) X_{103} (x_{106}) X_{109} (-x_{107}) X_{105} (x_{108}) X_{107} (x_{109}) X_{110} (x_{110}) X_{111} (x_{111}) X_{112} (x_{112}) X_{113} (x_{113}) X_{114} (x_{114}) X_{115} (x_{115}) X_{117} (-x_{116}) X_{116} (x_{117}) X_{118} (x_{118}) X_{119} (x_{119}) X_{120} (x_{120})$
7	$t_6 t_8 / t_7^2$	$h_7 (c_7 / r_7) h \cdot X_7 (- (c_7 / r_7)^2 x_7 \cdot X_1 (x_1) X_2 (x_2) X_3 (x_3) X_4 (x_4) X_5 (x_5) X_6 (-x_6 x_7^2) X_{14} (x_{14}) \cdot X_{15} (-x_8) X_9 (x_9) X_{10} (x_{10}) X_{11} (x_{11}) X_{12} (x_{12}) X_{21} (x_{21}) X_{25} (x_{25}) X_{27} (x_{27}) X_{31} (x_{31}) X_{32} (x_{32}) X_{33} (x_{33}) X_{34} (x_{34}) X_{35} (x_{35}) X_{36} (x_{36}) X_{37} (x_{37}) X_{38} (x_{38}) X_{39} (x_{39}) X_{48} (-x_{40}) X_{41} (x_{41}) X_{42} (x_{42}) X_{43} (x_{43}) X_{44} (x_{44}) X_{45} (-x_{46}) X_{46} (x_{46}) X_{47} (x_{47}) X_{48} (x_{48}) X_{55} (-x_{49}) X_{50} (x_{50}) X_{57} (-x_{51}) X_{45} (x_{52}) X_{59} (-x_{53}) X_{54} (x_{54}) X_{49} (x_{55}) X_{62} (-x_{56}) X_{51} (x_{56}) X_{64} (x_{57}) X_{51} (-x_{58}) X_{52} (-x_{59}) X_{60} (x_{60}) X_{61} (x_{61}) X_{62} (x_{62}) X_{70} (x_{63}) X_{57} (-x_{64}) X_{65} (x_{65}) X_{59} (x_{66}) X_{73} (-x_{67}) X_{62} (x_{68}) X_{69} (x_{69}) X_{76} (-x_{70}) X_{64} (x_{71}) X_{78} (-x_{72}) X_{67} (x_{73}) X_{74} (x_{74}) X_{80} (-x_{75}) X_{70} (x_{76}) X_{83} (-x_{77}) X_{72} (x_{78}) X_{79} (x_{79}) X_{75} (x_{80}) X_{86} (-x_{81}) X_{82} (x_{82}) X_{77} (x_{83}) X_{84} (x_{84}) X_{85} (x_{85}) X_{81} (x_{86}) X_{87} (x_{87}) X_{88} (x_{88}) X_{89} (x_{89}) X_{90} (x_{90}) X_{91} (x_{91}) X_{96} (-x_{92}) X_{93} (x_{93}) X_{94} (x_{94}) X_{100} (-x_{95}) X_{92} (x_{96}) X_{97} (x_{97}) X_{98} (x_{98}) X_{99} (x_{99}) X_{100} (x_{100}) X_{101} (x_{101}) X_{102} (x_{102}) X_{103} (x_{103}) X_{104} (x_{104}) X_{108} (-x_{105}) X_{106} (x_{106}) X_{109} (x_{109}) X_{110} (x_{110}) X_{111} (x_{111}) X_{112} (x_{112}) X_{113} (x_{113}) X_{114} (x_{114}) X_{115} (x_{115}) X_{116} (x_{116}) X_{118} (-x_{117}) X_{117} (x_{118}) X_{119} (x_{119}) X_{120} (x_{120})$
8	t_7 / t_8^2	$h_8 (c_8 / r_8) h \cdot X_8 (- (c_8 / r_8)^2 x_8 \cdot X_1 (x_1) X_2 (x_2) X_3 (x_3) X_4 (x_4) X_5 (x_5) X_6 (x_6) X_7 (-x_7 x_8^2) X_{15} (x_{15}) \cdot X_9 (x_9) X_{10} (x_{10}) X_{11} (x_{11}) X_{12} (x_{12}) X_{13} (x_{13}) X_{22} (x_{22}) X_{24} (x_{24}) X_{25} (x_{25}) X_{26} (x_{26}) X_{27} (x_{27}) X_{28} (x_{28}) X_{29} (x_{29}) X_{30} (x_{30}) X_{31} (x_{31}) X_{32} (x_{32}) X_{33} (x_{33}) X_{34} (x_{34}) X_{35} (x_{35}) X_{36} (x_{36}) X_{37} (x_{37}) X_{38} (x_{38}) X_{39} (x_{39}) X_{48} (-x_{40}) X_{41} (x_{41}) X_{42} (x_{42}) X_{43} (x_{43}) X_{44} (x_{44}) X_{45} (-x_{46}) X_{46} (x_{46}) X_{47} (x_{47}) X_{48} (x_{48}) X_{55} (-x_{49}) X_{50} (x_{50}) X_{57} (-x_{51}) X_{45} (x_{52}) X_{59} (-x_{53}) X_{54} (x_{54}) X_{49} (x_{55}) X_{62} (-x_{56}) X_{51} (x_{56}) X_{64} (x_{57}) X_{51} (-x_{58}) X_{52} (-x_{59}) X_{60} (x_{60}) X_{61} (x_{61}) X_{62} (x_{62}) X_{70} (x_{63}) X_{57} (-x_{64}) X_{65} (x_{65}) X_{59} (x_{66}) X_{73} (-x_{67}) X_{62} (x_{68}) X_{69} (x_{69}) X_{76} (-x_{70}) X_{64} (x_{71}) X_{78} (-x_{72}) X_{67} (x_{73}) X_{74} (x_{74}) X_{80} (-x_{75}) X_{70} (x_{76}) X_{83} (-x_{77}) X_{72} (x_{78}) X_{79} (x_{79}) X_{75} (x_{80}) X_{86} (-x_{81}) X_{82} (x_{82}) X_{77} (x_{83}) X_{84} (x_{84}) X_{85} (x_{85}) X_{81} (x_{86}) X_{87} (x_{87}) X_{88} (x_{88}) X_{89} (x_{89}) X_{90} (x_{90}) X_{91} (x_{91}) X_{96} (-x_{92}) X_{93} (x_{93}) X_{94} (x_{94}) X_{90} (x_{95}) X_{96} (x_{96}) X_{97} (x_{97}) X_{98} (x_{98}) X_{99} (x_{99}) X_{100} (x_{100}) X_{101} (x_{101}) X_{102} (x_{102}) X_{103} (x_{103}) X_{104} (x_{104}) X_{108} (-x_{105}) X_{106} (x_{106}) X_{109} (x_{109}) X_{110} (x_{110}) X_{111} (x_{111}) X_{112} (x_{112}) X_{113} (x_{113}) X_{114} (x_{114}) X_{115} (x_{115}) X_{116} (x_{116}) X_{118} (-x_{117}) X_{117} (x_{118}) X_{119} (x_{119}) X_{120} (x_{120})$

TABLE 9. Iwasawa coordinates for E_8 with compensating element $\kappa_i(c_i, x_i) = \chi_i(-c_i, x_i) h_i(-1/\sqrt{c_i, x_i})^2 + 1/\sqrt{c_i, x_i} \tilde{w}_i \chi_i(-c_i, x_i)$: Part 2

5. THE SL_2 AND E_7 COSETS

5.1. SL_2 symmetry in type IIB supergravity in $D = 10$ dimensions. Type IIB supergravity in $D = 10$ spacetime dimensions has two real scalar fields, the dilation ϕ and the axion χ . These can be combined to define a complex scalar field

$$\tau = \chi + ie^\phi,$$

which parametrizes the upper half-plane $\text{SO}(2) \backslash \text{SL}_2(\mathbb{R}) \cong \mathbb{H}$ and transforms under $\text{SL}_2(\mathbb{R})$ as the fractional linear transformation

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} : \tau \mapsto \frac{a\tau + b}{c\tau + d}.$$

The group $G = \text{SL}_2(\mathbb{R})$ is a global symmetry of the theory and G acts on the asymptotic values of the scalar fields at spatial infinity. The orbits of the global symmetry group G yield families of p -brane solutions ([CLPS]).

5.2. Action of $\text{SL}_2(\mathbb{Z})$ on the scalar sector. We will use the results of Section 3 but change some of the notation to follow the conventions of [HPS]. The coset element $V(x)$ of the (ϕ, χ) scalar sector is given in Borel gauge by ([HPS])

$$V(x) = \exp \left[\frac{\phi(x)}{2} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \right] \exp \left[\chi(x) \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \right] = \begin{pmatrix} e^{-\frac{\phi(x)}{2}} & \chi(x)e^{\frac{\phi(x)}{2}} \\ 0 & e^{\frac{\phi(x)}{2}} \end{pmatrix},$$

where $h = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$ and $e = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$ are $\mathfrak{sl}_2(\mathbb{R})$ generators and x is a spacetime coordinate. To determine the action of $\text{SL}_2(\mathbb{Z})$ on the coset element $V(x)$, it is enough to determine the action of the generators T and S on $V(x)$.

Theorem 5.1. *The action of the generators T and S of $\text{SL}_2(\mathbb{Z})$ on $V(x) \in \text{SO}(2) \backslash \text{SL}_2(\mathbb{R})$ is given in Iwasawa coordinates by*

$$V(x) \cdot T = \begin{pmatrix} e^{-\phi(x)/2} & 0 \\ 0 & e^{\phi(x)/2} \end{pmatrix} \begin{pmatrix} 1 & e^{\phi(x)}[e^{-\phi(x)} + \chi(x)] \\ 0 & e^{\phi(x)/2} \end{pmatrix}$$

and

$$V(x) \cdot S = \begin{pmatrix} e^{\phi(x)/2} \sqrt{1 + \chi(x)^2} & 0 \\ 0 & \frac{e^{-\phi(x)/2}}{\sqrt{1 + \chi(x)^2}} \end{pmatrix} \begin{pmatrix} 1 & -\frac{e^{-\phi(x)}\chi(x)}{1 + \chi(x)^2} \\ 0 & 1 \end{pmatrix}.$$

Proof: We start with the right action of T :

$$V(x) \cdot T = \begin{pmatrix} e^{-\phi(x)/2} & \chi(x)e^{\phi(x)/2} \\ 0 & e^{\phi(x)/2} \end{pmatrix} \cdot \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} e^{-\phi(x)/2} & e^{\phi(x)/2}[e^{-\phi(x)} + \chi(x)] \\ 0 & e^{\phi(x)/2} \end{pmatrix},$$

which may be written in Iwasawa coordinates as

$$V(x) \cdot T = \begin{pmatrix} e^{-\phi(x)/2} & 0 \\ 0 & e^{\phi(x)/2} \end{pmatrix} \begin{pmatrix} 1 & e^{\phi(x)}[e^{-\phi(x)} + \chi(x)] \\ 0 & e^{\phi(x)/2} \end{pmatrix}.$$

Similarly, we have for the action of S :

$$V(x) \cdot S = \begin{pmatrix} e^{-\phi(x)/2} & \chi(x)e^{\phi(x)/2} \\ 0 & e^{\phi(x)/2} \end{pmatrix} \cdot \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = \begin{pmatrix} -tu & t \\ -t^{-1} & 0 \end{pmatrix},$$

where $t = e^{-\phi(x)/2}$ and $u = \chi(x)e^{\phi(x)}$. However, this last expression is not in Iwasawa form, hence we multiply on the left by the compensating element (as in Proposition 3.3)

$$\kappa = \frac{1}{\sqrt{t^2 u^2 + t^{-2}}} \begin{pmatrix} -tu & -t^{-1} \\ t^{-1} & -tu \end{pmatrix}$$

to obtain the desired form

$$\kappa \cdot V(x) \cdot S = \begin{pmatrix} e^{\phi(x)/2} \sqrt{1 + \chi(x)^2} & 0 \\ 0 & \frac{e^{-\phi(x)/2}}{\sqrt{1 + \chi(x)^2}} \end{pmatrix} \begin{pmatrix} 1 & -\frac{e^{-\phi(x)} \chi(x)}{1 + \chi(x)^2} \\ 0 & 1 \end{pmatrix}.$$

□

For example, if $\phi(x) = \chi(x) = 0$, that is, at the point i , the action reduces to

$$V \cdot T = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix},$$

$$V \cdot S = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

Note that the latter result is expected since $S \in K = \text{SO}(2)$, so S fixes V at the point i .

5.3. The action of $G(\mathbb{Z})$ on $K \backslash G(\mathbb{R})$. Let $G = G(\mathbb{R})$ be a split, simply connected Lie group with maximal compact subgroup K as in Section 4. Let \mathfrak{g} be the Lie algebra of G . Let V be a highest-weight representation of \mathfrak{g} . Then $G \leq \text{Aut}(V)$. Provided V contains an admissible lattice $V_{\mathbb{Z}}$, we can define a \mathbb{Z} -form $G(\mathbb{Z})$ of G as the stabilizer of $V_{\mathbb{Z}}$ ([St]). In order to describe the action of $G(\mathbb{Z})$ on $K \backslash G(\mathbb{R})$, we need a generating set for $G(\mathbb{Z})$. The following result provides generators for a large class of integral forms:

Theorem 5.2 ([CCa],[BC]). *Let V be a highest-weight representation of \mathfrak{g} and suppose that the lattice generated by the weights of V contains all the fundamental weights. Then $G(\mathbb{Z})$ is generated by $\chi_{\alpha_i}(1)$ and $\tilde{w}_{\alpha_i} = \chi_{\alpha_i}(1)\chi_{-\alpha_i}(-1)\chi_{\alpha_i}(1)$, for $i = 1, \dots, n$.*

The Iwasawa coordinates are

$$Khu = K \prod_{j=1}^n h_j(t_j) \prod_{r=1}^N \chi_r(x_r)$$

where $t_j, x_r \in \mathbb{R}$, and $t_j > 0$.

The right action of $\chi_i(1)$ for $i = 1, \dots, n$, is given by

$$\prod_{j=1}^n h_j(t_j) \prod_{r=1}^N \chi_r(x_r) \cdot \chi_i(1),$$

which can be rewritten as an expression in Iwasawa form using the collection algorithms of [CMH]. However the right action of \tilde{w}_i for $i = 1, \dots, n$, is given by

$$\prod_{j=1}^n h_j(t_j) \prod_{r=1}^N \chi_r(x_r) \cdot \tilde{w}_i,$$

which is not in Iwasawa form. To remedy this, we use the compensating element $\kappa \in K$ of Section 4 to left multiply the expression by $\kappa_i(c_i x_i)$ which gives an expression in Iwasawa form.

This result applies to the noncompact split real form $E_{7(+7)}(\mathbb{R})$ of the exceptional Lie group E_7 with the following form of $E_7(\mathbb{Z})$:

$$E_{7(+7)}(\mathbb{Z}) = E_{7(+7)}(\mathbb{R}) \cap \mathrm{Sp}_{56}(\mathbb{Z}) ,$$

discovered in [HT] following [CJ2] in the framework of type II string theory. Soulé gave a rigorous mathematical proof that the $E_{7(+7)}(\mathbb{Z})$ of Hull and Townsend coincides with the Chevalley \mathbb{Z} -form $G(\mathbb{Z})$ of $G = E_7$ ([S]). Here $E_{7(+7)}(\mathbb{R}) \cap \mathrm{Sp}_{56}(\mathbb{Z})$ is the stabilizer of the standard lattice in the fundamental representation of E_7 which has dimension 56. The charge lattice of [HT] can be normalized to coincide with the lattice $V_{\mathbb{Z}}$. Once a basis for $V_{\mathbb{Z}}$ has been chosen, the $E_7(\mathbb{Z})$ orbits can be computed explicitly in terms of this basis. The non-compact split real form $E_{7(+7)}(\mathbb{R})$ has maximal compact subgroup $K = [\mathrm{SU}(8)/\{\pm \mathrm{Id}\}]$. The coset $[\mathrm{SU}(8)/\{\pm \mathrm{Id}\}] \backslash E_{7(+7)}(\mathbb{R})$ occurs as a scalar coset for dimensional reduction of $N = 1$ supergravity in 11 dimensions to $N = 8$ supergravity in four dimensions ([CJ2]).

6. THE ACTION OF $\mathrm{SL}_2(\mathbb{Z})$ ON THE CHARGE LATTICE

For now, we assume the existence of dyons, that is, particles that carry both electric and magnetic charge, see [O1] and [O2]. The Dirac quantization condition was generalized by Zwanziger and Schwinger to a pair of dyons with

$$\begin{aligned} &\text{electric charges } q_1 \text{ and } q_2 \\ &\text{magnetic charges } g_1 \text{ and } g_2 \end{aligned}$$

such that

$$(17) \quad q_1 g_2 - q_2 g_1 = 2\pi n \hbar, \quad n = 0, \pm 1, \pm 2, \dots .$$

Witten solved this quantization condition giving a general family of solutions $\{(q, g)\}$ of electric and magnetic charges respectively, with

$$q + ig = q_0(m\tau + n)$$

where $m, n \in \mathbb{Z}$ are coprime, q_0 is the electron charge, and τ is the parameter

$$\tau = \frac{\theta}{2\pi} + \frac{2\pi i n \hbar}{q_0^2} ,$$

where θ is the ‘‘vacuum angle’’ ([Wi]). Here τ is a complex variable with positive imaginary part and which depends on dimensionless parameters corresponding to the particular theory.

If we equate real and imaginary parts of $q + ig = q_0(m\tau + n)$ and set $m = 0$, then we obtain $q = q_0 n$, $n \in \mathbb{Z}$. This means that the electric charge is an integer multiple of the basic electron charge q_0 .

The collection of possible dyonic charges is encoded in the *charge lattice*

$$\mathcal{Q} = \{q + ig \mid q, g \text{ satisfy condition (17)}\} \subset \mathbb{C} ,$$

depicted in Figure 2.

Here electric particles lie on the real axis, monopoles lie on the imaginary axis, and a general point in the (q, g) plane off the axes is a dyon. The origin represents a state with no electric or magnetic charge, hence could be taken to represent a photon or a charge-less (flat) background. For magnetically neutral states, electric charges are integral multiples of q_0 . The open circles represent ‘primitive vectors’ on the lattice. These are vectors that can be connected to the origin by a line that does not intersect any other lattice points. Equivalently, (m, n) are relatively prime

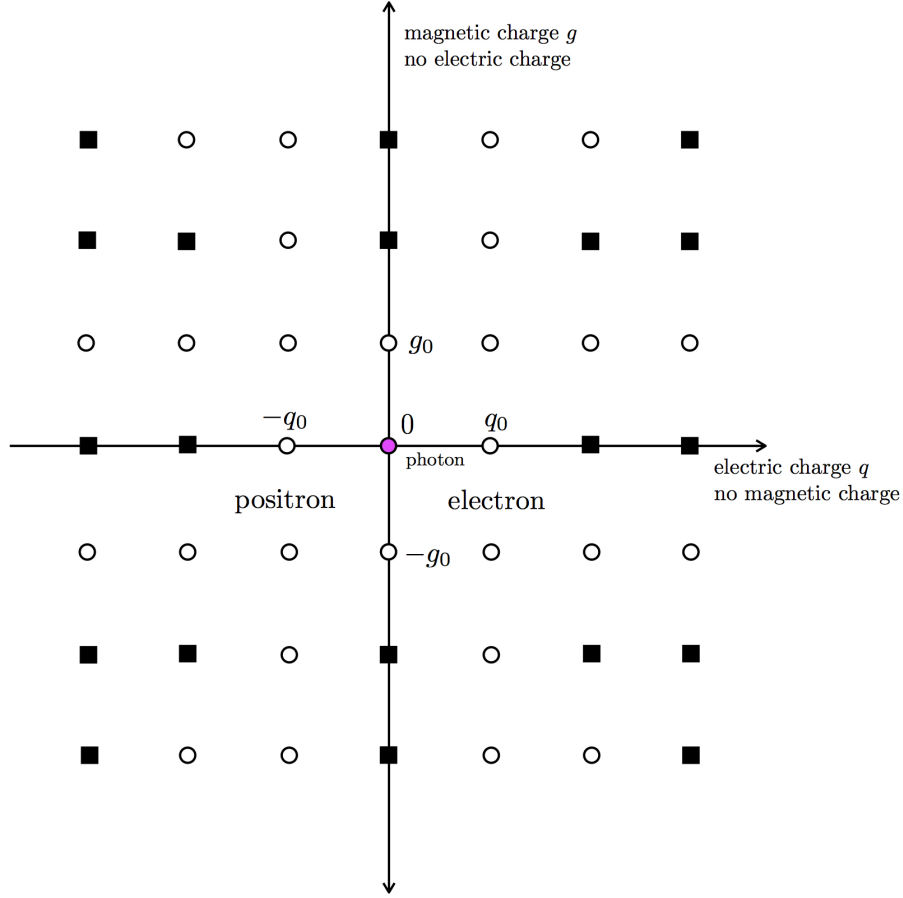


FIGURE 2. The charge lattice \mathcal{Q}

for these points in the equation $q + ig = q_0(m\tau + n)$. The action of the modular group $SL_2(\mathbb{Z})$ preserves primitive vectors in the charge lattice.

Given a specific theory, an important question is identifying the subset of the charge lattice that can be realized by single-particle states, as opposed to multi-particle states. Here it is the primitive vectors that represent single-particle states ([O2]).

We now look at symmetries of the charge lattice. Firstly, the Weyl group $W \cong \mathbb{Z}_2$ permutes the axes of the charge vector lattice, interchanging electric and magnetic charges. Secondly, the generalization of duality symmetry of Maxwell is a rotation by $\pi/2$ which exchanges electrically charged particles and magnetic monopoles. The $\pi/2$ -rotational symmetry

$$\begin{aligned} (q_0, 0) &\rightarrow (0, -g_0), & (-q_0, 0) &\rightarrow (0, g_0), \\ (0, g_0) &\rightarrow (q_0, 0), & (0, -g_0) &\rightarrow (-q_0, 0) \end{aligned}$$

preserves the magnitudes of electric and magnetic charges. The rotation group R of the elementary charges is hence the cyclic group of order 4 generated by the matrix

$$a = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix},$$

that is

$$R = \mathrm{SL}_2(\mathbb{Z}) \cap \mathrm{SO}(2).$$

The discrete subgroup R acts by rotating the elementary charge vector $(q_0, 0)$ by integral multiples of $\pi/2$, while keeping its length fixed. The group R is the full rotation group of the charge lattice, and the rotational symmetry preserves primitive vectors and interchanges electric and magnetic charges.

The charge lattice can be characterized algebraically. Let $G = \mathrm{SL}_2(\mathbb{C})$ and consider the vector spaces

$$V_{\mathbb{Z}} = \mathbb{Z} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \oplus \mathbb{Z} \begin{bmatrix} 0 \\ 1 \end{bmatrix},$$

$$V_{\mathbb{C}} = \mathbb{C} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \oplus \mathbb{C} \begin{bmatrix} 0 \\ 1 \end{bmatrix}.$$

Then $V_{\mathbb{C}}$ is a highest weight module for the Lie algebra $\mathfrak{g} = \mathfrak{sl}_2(\mathbb{C})$. The vector space $V_{\mathbb{C}}$ can be identified with the charge lattice \mathcal{Q} by setting the magnitude of the elementary electric charge q_0 equal to 1. The group $\mathrm{SL}_2(\mathbb{Z})$ can be viewed as the subgroup of $\mathrm{SL}_2(\mathbb{C})$ that stabilizes the vector space $V_{\mathbb{Z}}$.

In the absence of magnetic charges, the space of states allowed by the Dirac quantization condition, $p = p_0 n$, $n \in \mathbb{Z}$, is the set of points (p, q) with $p, q \in \mathbb{Z} - \{0\}$. In general, the charge lattice contains pairs of charges (p, q) on both the real and imaginary axes with $p, q \neq 0$. As before, primitive vectors correspond to coordinates (p, q) with $p, q \in \mathbb{Z} - \{0\}$, $\mathrm{gcd}(p, q) = 1$.

The following proposition is easy to verify.

Proposition 6.1. *The action of $\mathrm{SL}_2(\mathbb{Z})$ on the set of allowed states (p, q) with $p, q \in \mathbb{Z} - \{0\}$ is not transitive. The action of the generators of $\mathrm{SL}_2(\mathbb{Z})$ is:*

$$\begin{aligned} \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} &= \begin{bmatrix} 1 \\ 0 \end{bmatrix}, & \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} &= \begin{bmatrix} 1 \\ 1 \end{bmatrix} \\ \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} &= \begin{bmatrix} 1 \\ 1 \end{bmatrix}, & \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} &= \begin{bmatrix} 0 \\ 1 \end{bmatrix}. \end{aligned}$$

Thus there is a subset of fixed vectors. The action of $\mathrm{SL}_2(\mathbb{Z})$ on the set (p, q) with $p, q \in \mathbb{Z} - \{0\}$, $\mathrm{gcd}(p, q) = 1$ is transitive.

6.1. Spectrum-generating symmetries for BPS solitons. The concept of dyons and their corresponding charges from electromagnetism generalize to ‘higher notions’ of particles, meaning extended objects, for example, p -branes (see [DGHT]). The case $p = 0$ corresponds to the case of point particles described above. A *BPS-saturated soliton* is a solution of a supergravity theory that describes the low energy limit of a string theory in which infinite p -branes occupy a longitudinal submanifold of spacetime. Associated to such a solution is a vector space of electric or magnetic charges, or both, if dyonic p -branes are present.

One important goal is to determine the spectrum-generating symmetries for the fundamental BPS solitons. In [CLPS], it was shown that the standard global supergravity symmetry group is not sufficient for the purpose of generating complete sets of p -brane solitons. The additional

ingredient needed is a scaling transformation that allows one to map between BPS solitons with different masses. To that end, consider the matrices

$$\sigma(s) = \begin{pmatrix} s & 0 \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad \tau(t) = \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix},$$

for fixed $s, t \in \mathbb{R}_{>0}$. These elements σ and τ belong to $\text{GL}_2(\mathbb{R})$ and are the ‘trombone’ scaling symmetries of [CLPS]. We note however, that [CLPS] use only one trombone symmetry which is combination of σ and τ . These elements are symmetries of the equations of motion, but not of the action, are invariant under dimensional reduction, and preserve the dilation and axion scalar fields ([CLPS]).

Indeed, for any $s, t \in \mathbb{R}_{>0}$, we have

$$\begin{aligned} \begin{pmatrix} s & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} q_0 \\ g_0 \end{pmatrix} &= \begin{pmatrix} sq_0 \\ g_0 \end{pmatrix}, \\ \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix} \begin{pmatrix} q_0 \\ g_0 \end{pmatrix} &= \begin{pmatrix} q_0 \\ tg_0 \end{pmatrix}, \end{aligned}$$

so $\sigma(s)$ rescales the elementary electric charge and $\tau(t)$ rescales the elementary magnetic charge.

Let $P(\mathbb{R})$ be the subgroup of $\text{GL}_2(\mathbb{R})$ generated by $K = \text{SO}(2)$ and the elements

$$\sigma(\mathbb{R}) = \left\{ \begin{pmatrix} s & 0 \\ 0 & 1 \end{pmatrix} \mid s \in \mathbb{R} \right\}$$

and

$$\tau(\mathbb{R}) = \left\{ \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix} \mid t \in \mathbb{R} \right\}.$$

As in [CLPS], we have the following.

Theorem 6.2. *At the classical level, the subgroup $P(\mathbb{R}) = \langle \text{SO}(2), \sigma(\mathbb{R}), \tau(\mathbb{R}) \rangle$ of $\text{GL}_2(\mathbb{R})$ is the spectrum-generating symmetry group for the parameter space of BPS solitons that preserve 1/2 of the supersymmetry of Type IIB supergravity in $D = 10$ dimensions.*

Next consider

$$P(\mathbb{Z}) = \langle \text{SL}_2(\mathbb{Z}), \sigma(\mathbb{Z}), \tau(\mathbb{Z}) \rangle$$

where

$$\sigma(\mathbb{Z}) = \left\{ \begin{pmatrix} s & 0 \\ 0 & 1 \end{pmatrix}, s \in \mathbb{Z} \right\} \quad \text{and} \quad \tau(\mathbb{Z}) = \left\{ \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix}, t \in \mathbb{Z} \right\}.$$

Then $\text{SL}_2(\mathbb{Z})$ is naturally contained in $P(\mathbb{Z})$. However, $P(\mathbb{Z})$ is not a group, as the inverses of $\sigma(\mathbb{Z})$ and $\tau(\mathbb{Z})$ are not contained in $P(\mathbb{Z})$. We can identify $P(\mathbb{Z})$ as a *monoid* and in fact a submonoid of $\text{GL}_2(\mathbb{Z}_{\geq 0})$, where $\text{GL}_2(\mathbb{Z}_{\geq 0}) = M_2(\mathbb{Z}_{\geq 0}) \cap \text{GL}_2(\mathbb{Z})$ is defined as the monoid of matrices of determinant $\neq 0$ with entries in $\mathbb{Z}_{\geq 0}$. The elements of $\text{GL}_2(\mathbb{Z}_{\geq 0})$ are not necessarily invertible, since inverses may not exist over $\mathbb{Z}_{\geq 0}$.

To see that $P(\mathbb{Z})$ is a submonoid of $\text{GL}_2(\mathbb{Z}_{\geq 0})$, we use the following.

Theorem 6.3. ([R])

- (1) $\text{SL}_2(\mathbb{Z}_{\geq 0}) = M_2(\mathbb{Z}_{\geq 0}) \cap \text{SL}_2(\mathbb{Z})$ is a free monoid on two generators $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$, $\begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix}$.

(2) $\mathrm{SL}_2(\mathbb{Z})$ can be decomposed using 8 copies of $\mathrm{SL}_2(\mathbb{Z}_{\geq 0})$:

$$\begin{aligned} \mathrm{SL}_2(\mathbb{Z}) &= \mathrm{SL}_2(\mathbb{Z}_{\geq 0}) \sqcup \gamma \mathrm{SL}_2(\mathbb{Z}_{\geq 0}) \sqcup \mathrm{SL}_2(\mathbb{Z}_{\geq 0})\gamma \sqcup \gamma \mathrm{SL}_2(\mathbb{Z}_{\geq 0})\gamma \\ &\sqcup (-1) \mathrm{SL}_2(\mathbb{Z}_{\geq 0}) \sqcup (-1)\gamma \mathrm{SL}_2(\mathbb{Z}_{\geq 0}) \sqcup (-1) \mathrm{SL}_2(\mathbb{Z}_{\geq 0})\gamma \sqcup (-1)\gamma \mathrm{SL}_2(\mathbb{Z}_{\geq 0})\gamma, \end{aligned}$$

$$\text{where } \gamma = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \text{ and } (-1) = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}$$

(3) $\mathrm{GL}_2(\mathbb{Z})$ can be decomposed using 2 copies of $\mathrm{SL}_2(\mathbb{Z})$: $\mathrm{GL}_2(\mathbb{Z}) = \mathrm{SL}_2(\mathbb{Z}) \sqcup \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \mathrm{SL}_2(\mathbb{Z})$.

It follows that the generators for $\mathrm{GL}_2(\mathbb{Z}_{\geq 0})$ are

$$\begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix} \text{ and } \begin{pmatrix} 1 & 0 \\ v & 1 \end{pmatrix}$$

for $u, v \in \mathbb{Z}_{\geq 0}$ and

$$\begin{pmatrix} s & 0 \\ 0 & 1 \end{pmatrix} \text{ and } \begin{pmatrix} 1 & 0 \\ 0 & t \end{pmatrix}$$

for $s, t \in \mathbb{Z}_{> 0}$. Thus $P(\mathbb{Z})$ is a submonoid of $\mathrm{GL}_2(\mathbb{Z}_{\geq 0})$.

Let $(p_0, 0)^T$, $(0, q_0)^T$ be the elementary vectors of electric charge. It is convenient to choose $p_0 = 1$. Then we must have $q_0 \neq 1$ since the vector (p_0, q_0) is primitive. We can generate all primitive vectors using the $P(\mathbb{Z})$ action on $(p_0, 0)^T$ and $(0, q_0)^T$ in the following way. We choose matrices

$$m_1 = \begin{pmatrix} n & b \\ q_0 & d \end{pmatrix}, m_2 = \begin{pmatrix} a & p_0 \\ c & m \end{pmatrix}, m_3 = \begin{pmatrix} s & 0 \\ 0 & t \end{pmatrix} \in P(\mathbb{Z})$$

with $a, b, c, d, p_0, q_0, m, n, s, t \in \mathbb{Z}$, in such a way that

$$\mathrm{gcd}(n, q_0) = \mathrm{gcd}(p_0, m) = \mathrm{gcd}(s, t) = 1.$$

Then

$$\begin{pmatrix} n & b \\ q_0 & d \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} n \\ q_0 \end{pmatrix}, \quad \begin{pmatrix} a & p_0 \\ c & m \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} p_0 \\ m \end{pmatrix}, \quad \text{and} \quad \begin{pmatrix} s & 0 \\ 0 & t \end{pmatrix} \begin{pmatrix} p_0 \\ q_0 \end{pmatrix} = \begin{pmatrix} sp_0 \\ tq_0 \end{pmatrix}.$$

Thus all the primitive vectors are generated by the action of m_1, m_2 or $m_3 \in P(\mathbb{Z})$.

Even though the matrices m_1, m_2 and m_3 have integral entries and are contained in $P(\mathbb{Z})$, their inverses cannot be contained in $P(\mathbb{Z})$. This it is a subtle question to determine what it means to ‘invert’ a charge under the action of $P(\mathbb{Z})$. One can always invert charges under the action of the group $\mathrm{SL}_2(\mathbb{Z})$ which preserves the primitive vectors and hence maps between allowed charge vectors. Hence the group $\mathrm{SL}_2(\mathbb{Z})$ generates the spectrum of physically distinct states in a single fixed vacuum.

Theorem 6.3 also ensures that charges can be inverted under the action of the group $\mathrm{GL}_2(\mathbb{Z})$.

However, we have a quantum anomaly of the theory: the trombone symmetries $\sigma(\mathbb{Z})$ and $\tau(\mathbb{Z})$ cannot be inverted over \mathbb{Z} in the monoid $P(\mathbb{Z})$ that they generate.

We may summarize the above discussion in the following.

Lemma 6.4. *Elements of the cosets*

$$\mathrm{SO}(2)\sigma(\mathbb{Z}), \quad \mathrm{SO}(2)\tau(\mathbb{Z})$$

generate the full charge lattice from the elementary charge vector $(1, 0)^T$ and preserve the moduli of the scalar fields.

The elements of $\{\mathrm{SO}(2)\sigma(\mathbb{Z}), \mathrm{SO}(2)\tau(\mathbb{Z})\}$ are not integral matrices and moreover do not generate a group, but rather cosets of $\mathrm{SO}(2)\backslash\mathrm{GL}_2(\mathbb{R})$.

This shows that the usual duality group symmetries at the classical level are not inherited in the same form at the quantum level and one must take care when describing the discrete symmetries in the quantum theory.

7. U-DUALITY GROUPS IN GENERAL

Our method for computing Iwasawa coordinates of the quotient $K\backslash G$ and the action of the \mathbb{Z} -form $G(\mathbb{Z})$ on the coset $K\backslash G$ do not depend on the choice of group G . Our techniques can hence be applied to other U-duality groups, including the hyperbolic Kac–Moody group E_{10} , which is conjectured to be the U-duality group of 11-dimensional supergravity in 1 dimension ([HT], [DHN]).

In [AC], the authors obtained a finite presentation of an important class of hyperbolic Kac–Moody groups, including $E_{10}(\mathbb{R})$ and $E_{10}(\mathbb{Z})$. Hence the action of $E_{10}(\mathbb{Z})$ on the coset $K(E_{10}(\mathbb{R}))\backslash E_{10}(\mathbb{R})$ is given by finitely many rules which can be determined by a generalization of the methods given here.

The \mathbb{Z} -form $G(\mathbb{Z})$, for any simple and simply connected Chevalley group or Kac–Moody group G , may be defined as the stabilizer

$$G(\mathbb{Z}) = \{g \in G(\mathbb{R}) \mid g \cdot V_{\mathbb{Z}} \subseteq V_{\mathbb{Z}}\}$$

of a lattice $V_{\mathbb{Z}}$ in a suitable highest weight module V for the Lie algebra \mathfrak{g} of G ([St], [BC]). In many cases of physical interest, a fundamental representation serves as a suitable choice of V .

For a general discrete duality group $G(\mathbb{Z})$, we may also determine the action of $G(\mathbb{Z})$ on \mathbb{Z} -forms $V_{\mathbb{Z}}$ of fundamental modules V for Lie algebras and Kac–Moody algebras in general. The charge lattice \mathcal{Q} of Section 6 can be normalized to coincide with the lattice $V_{\mathbb{Z}}$. Once a basis for $V_{\mathbb{Z}}$ has been chosen, the $G(\mathbb{Z})$ orbits on the charge lattice can be computed explicitly in terms of this basis. We hope to take up these open questions in forthcoming work.

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