# UNIFORM ESTIMATES AND BLOW-UP BEHAVIOR FOR SOLUTIONS OF $-\Delta u = V(x)e^{u}$ IN TWO DIMENSIONS

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### **Introduction**

In this paper we deal with the equation

(\*) 
$$\begin{cases} -\Delta u = V(x)e^{u} \text{ in } \Omega \in \mathbb{R}^{2}, \\ u = 0 \quad \text{on } \partial \Omega, \end{cases}$$

where  $\Omega$  is a bounded domain (except in Section II.3) and V(x) is a given function in  $L^{p}(\Omega)$  for some  $1 . We assume that <math>u \in L^{1}(\Omega)$  and  $e^{u} \in L^{p'}(\Omega)$  (where p' is the conjugate exponent of p) so that (\*) has a meaning in the sense of distributions.

A first question is whether one can conclude that  $u \in L^{\infty}(\Omega)$ . As we will see in Section II the answer is positive. Next we turn, in Section III, to a more delicate issue, namely the question of <u>uniform estimates</u>. Suppose we have a sequence  $(u_n)$  of solutions of

(\*\*) 
$$\begin{cases} -\Delta u_n = V_n(x)e^{u_n} & \text{in } \Omega, \\ u_n = 0 & \text{on } \partial\Omega, \end{cases}$$

with

$$\|\mathbf{V}_{\mathbf{n}}\|_{\mathbf{L}^{\mathbf{p}}} \leq \mathbf{C}_{\mathbf{1}}$$

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and

$$\|\mathbf{e}^{\mathbf{u}_{\mathbf{n}}}\|_{L^{\mathbf{p}'}} \leq C_{2}$$

Can one conclude that

 $\|\mathbf{u}_n\|_{L^{\infty}} \leq C_3$ 

where  $C_3$  depends only on  $C_1$ ,  $C_2$  and  $\Omega$ ? We prove that the answer is positive under a <u>smallness</u> condition, namely  $C_1C_2 < 4\pi/p'$  (see Corollary 3). The answer is also positive under a <u>domination</u> condition, namely  $|V_n| \leq W$  for a fixed  $W \in L^p(\Omega)$ ,  $1 (and then <math>C_3$  depends also on W, see Corollary 5).

A deeper result (see Corollary 6) asserts that if  $V_n \ge 0$  then  $(u_n)$  is bounded in  $L^{\infty}_{1 \circ C}(\Omega)$ , i.e. for every compact subset K of  $\Omega$  we have

$$\|\mathbf{u}_{\mathbf{n}}\|_{L^{\infty}(\mathbf{K})} \leq C_{3}$$

where  $C_3$  depends only on  $C_1$ ,  $C_2$  and K. Surprisingly such an estimate does not hold up to the boundary. Given any 1 we construct in Example 6 $(Section III.3) sequences <math>(u_n)$  and  $(V_n)$  satisfying (\*\*) with  $V_n \ge 0$ 

$$\|\mathbf{V}_{\mathbf{n}}\|_{\mathbf{L}^{\mathbf{p}^{\prime}}} \leq C_{1}$$
$$\|e^{\mathbf{u}_{\mathbf{n}}}\|_{\mathbf{L}^{\mathbf{p}^{\prime}}} \leq C_{2}$$

and  $\|u_n\|_{T^{\infty}} \rightarrow +\infty$ .

A corollary of our methods also yields the following. Suppose un satisfies

$$-\Delta u_n = V_n e^{u_n}$$
 in  $\Omega$ 

with

$$0 < a \leq V_n \leq b < \infty$$

and

$$\inf_{\Omega} u_n \ge -M > -\alpha$$

(here no boundary condition is imposed). Then for every compact subset K of  $\Omega$ , Sup  $u_n$  can be estimated just in terms of a,b,M,K and  $\Omega$  (see Corollary 8).

Finally we turn to the general case where no boundary condition is imposed and  $(u_n)$  is not bounded below. More precisely let  $(u_n)$  be a sequence of solutions of

$$-\Delta u_n = V_n e^{u_n}$$
 in  $\Omega$ 

with

$$\mathbf{V_n} \ge \mathbf{0} \quad \text{in} \quad \mathbf{\Omega}, \ \left\|\mathbf{V_n}\right\|_{\mathbf{L}^p} \le \mathbf{C_1} \quad \text{and} \quad \left\|\mathbf{e}^{u_n}\right\|_{\mathbf{L}^{p'}} \le \mathbf{C_2},$$

for some 1 .

Then we have the following alternative (see Theorem 3): either

(i)  $(u_n)$  is bounded in  $L_{loc}^{\infty}(\Omega)$ 

or

(ii) 
$$u_n \rightarrow -\infty$$
 uniformly on compact subsets of  $\Omega$ 

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(iii) there is a finite nonempty set S such that  $u_n \to -\infty$  uniformly on compact subsets of  $\Omega \setminus S$  and  $u_n \to +\infty$  on S (in a sense to be precised later). In this case  $V_n e^{u_n}$  converges to a finite sum of Dirac masses  $\Sigma \alpha_i \delta_{a_i}$  with coefficients  $\alpha_i \ge 4\pi/p'$ .

Such behavior is well illustrated by the sequence

$$u_n(x) = \log \frac{8n^2}{(1+n^2|x|^2)^2}$$

which satisfies  $-\Delta u_n = e^{u_n}$ ,  $||e^{u_n}||_{L^1} \leq C$ ,  $u_n(x) \rightarrow -\infty$  for all  $x \neq 0$  and  $u_n(0) \rightarrow +\infty$ . Here  $e^{u_n}$  converges to  $8\pi\delta_0$ .

We thank Congring Li for raising questions which led us to Theorem 2 and Corollary 3 (Theorem 2 is used in [3]). Some of our results (in particular Corollary 4 and Theorem 4) are connected to earlier works of Nagasaki and Suzuki (see [6] and [7]) who consider mostly the case where the  $V_n$ 's are constants. A. Chang and P. Yang [2] have also studied blow-up sequences for related equations on  $S^2$  (see e.g. their Concentration Lemma). However their approach involves  $H^1$  norms and is quite different from ours.

In a forthcoming work we shall consider similar issues for the equation  $-\Delta u = V(x)u^p$  in  $\Omega \in \mathbb{R}^N$ ,  $N \ge 3$ . The plan of the paper is the following:

### Introduction

- I. A basic inequality
- II.  $L^{\omega}$ -boundedness for a single solution of  $-\Delta u = Ve^{u}$ 
  - II.1. The case of a bounded domain
  - II.2. Some variants and counterexamples
  - II.3. The case  $\Omega = \mathbb{R}^2$

III. Uniform  $L^{\infty}$  bounds and blow-up behavior for solutions of  $-\Delta u = Ve^{u}$ 

- III.1. Some easy cases
- III.2. The main results
- III.3. Variants and counterexamples.

### I. <u>A basic inequality</u>

Assume  $\Omega \subset \mathbb{R}^2$  is a bounded domain and let u be a solution of

(1) 
$$\begin{cases} -\Delta u = f(x) & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega \end{cases}$$

with  $f \in L^{1}(\Omega)$ . Set  $||f||_{1} = \int_{\Omega} |f(\mathbf{x})| d\mathbf{x}$ .

<u>Theorem 1</u>. For every  $\delta \in (0, 4\pi)$  we have

(2) 
$$\int_{\Omega} \exp\left[\frac{(4\pi-\delta)|\mathbf{u}(\mathbf{x})|}{\|\mathbf{f}\|_{1}}\right] d\mathbf{x} \leq \frac{4\pi^{2}}{\delta} (\operatorname{diam} \Omega)^{2}.$$

<u>Proof.</u> Let  $R = \frac{1}{2} \operatorname{diam} \Omega$  so that  $\Omega \in B_R$  for some ball of radius R. Extend f to be zero outside  $\Omega$  and set, for  $x \in \mathbb{R}^2$ ,

$$\bar{u}(\mathbf{x}) = \frac{1}{2\pi} \int_{B_{R}} \log(\frac{2R}{|\mathbf{x}-\mathbf{y}|}) |f(\mathbf{y})| d\mathbf{y}$$

so that

$$-\Delta \bar{u} = |f|$$
 on  $\mathbb{R}^2$ .

Note that  $\bar{u}(x) \ge 0$  for  $x \in B_R$  since  $\frac{2R}{|x-y|} \ge 1 \quad \forall x, y \in B_R$ . It follows from the maximum principle that  $|u| \le \bar{u}$  on  $\Omega$  and thus

(3) 
$$\int_{\Omega} \exp\left[\frac{(4\pi-\delta)|\mathbf{u}(\mathbf{x})|}{\|\mathbf{f}\|_{1}}\right] d\mathbf{x} \leq \int_{\mathbf{B}_{\mathbf{R}}} \exp\left[\frac{(4\pi-\delta)\mathbf{\ddot{u}}(\mathbf{x})}{\|\mathbf{f}\|_{1}}\right] d\mathbf{x}.$$

We now estimate the right-hand side of (3) using Jensen's inequality

$$F(\int w(y)\varphi(y)dy) \leq \int w(y)F(\varphi(y))dy$$
with  $F(t) = \exp t$ ,  $w(y) = \frac{|f(y)|}{||f||_1}$  and  $\varphi(y) = \frac{(4\pi-\delta)}{2\pi} \log(\frac{2R}{|x-y|})$ . We obtain
$$\int_{B_R} \exp\left[\frac{(4\pi-\delta)\bar{u}(x)}{||f||_1}\right] dx \leq \int_{B_R} dx \int_{B_R} \left(\frac{2R}{|x-y|}\right)^{2-\frac{\delta}{2\pi}} \frac{|f(y)|}{||f||_1} dy$$

$$= \int_{B_R} \frac{|f(y)|}{||f||_1} \left[ \left( \int_{B_R} \left(\frac{2R}{|x-y|}\right)^{2-\frac{\delta}{2\pi}} dx \right) dy.$$

But, for  $y \in B_R$ , we have

$$\int_{B_{R}} \left(\frac{2R}{|x-y|}\right)^{2-\frac{\delta}{2\pi}} dx \leq \int_{B_{R}} \left(\frac{2R}{|x|}\right)^{2-\frac{\delta}{2\pi}} dx = \frac{4\pi^{2}}{\delta} \left(\operatorname{diam} \Omega\right)^{2}$$

and the estimate (2) follows.

A simple consequence of Theorem 1 is

<u>Corollary 1</u>. Let u be a solution of (1) with  $f \in L^{1}(\Omega)$ . Then for every constant k > 0

$$e^{\mathbf{k}|\mathbf{u}|} \in L^1(\Omega).$$

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<u>Proof.</u> Let  $0 < \epsilon < 1/k$ . We may split f as  $f = f_1 + f_2$  with  $||f_1||_1 < \epsilon$ and  $f_2 \in L^{\infty}(\Omega)$ . Write  $u = u_1 + u_2$  where  $u_i$  are the solutions of

$$\begin{cases} -\Delta u_i = f_i & \text{in } \Omega, \\ u_i = 0 & \text{on } \partial \Omega \end{cases}$$

Choosing, for example,  $\delta = (4\pi - 1)$  in Theorem 1 we find  $\int_{\Omega} \exp\left[\frac{|u_1(x)|}{||f_1||_1}\right] < \infty$ and thus  $\int_{\Omega} \exp[k|u_1|] < \infty$ . The conclusion follows since  $|u| \le |u_1| + |u_2|$ and  $u_2 \in L^{\infty}(\Omega)$ .

<u>Remark 1</u>. The conclusion of Theorem 1 could also be deduced from BMO estimates and the John-Nirenberg inequality [4].

<u>Remark 2</u>. There is a local form of Corollary 1, namely if  $u \in L^{1}_{loc}(\Omega)$  and  $\Delta u \in L^{1}_{loc}(\Omega)$ , then for every k > 0,  $e^{k|u|} \in L^{1}_{loc}(\Omega)$ . [Here we use the well-known fact that  $u \in L^{1}_{loc}(\Omega)$  and  $\Delta u \in L^{1}_{loc}(\Omega)$  imply  $\nabla u \in L^{1}_{loc}(\Omega)$ .]

<u>Remark 3</u>. In Corollary 1,  $e^{\mathbf{k} |\mathbf{u}|} \in L^1$  but  $||e^{\mathbf{k} |\mathbf{u}|}||_1$  can <u>not</u> be estimated in terms of  $\mathbf{k}$  and  $||f||_1$ . For example, we may have a sequence  $(f_n)$  such that  $||f_n||_1 \leq 1$ ,  $f_n \rightarrow \delta_{\mathbf{x}_0}$  and then  $u_n \rightarrow u$  with  $u(\mathbf{x}) \simeq \frac{1}{2\pi} \log \frac{1}{|\mathbf{x} - \mathbf{x}_0|}$  as  $\mathbf{x} \rightarrow \mathbf{x}_0$  so that  $\int e^{\mathbf{k} |\mathbf{u}|} = \mathbf{\omega}$  for  $\mathbf{k} \geq 4\pi$ .

II.  $\underline{L}^{\infty}$ -boundedness for a single solution of  $-\Delta u = Ve^{u}$ .

II.1. The case of a bounded domain.

Let u satisfy the nonlinear equation

(4) 
$$\begin{cases} -\Delta u = V(x)e^{u} & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega \end{cases}$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^2$  and V(x) is a given function on  $\Omega$ .

<u>Corollary 2</u>. Suppose u is a solution of (4) with  $V \in L^{p}(\Omega)$  and  $e^{u} \in L^{p'}(\Omega)$ for some  $1 . Then <math>u \in L^{\infty}(\Omega)$ . <u>Proof.</u> By Corollary 1 we know that  $e^{ku} \in L^{1}(\Omega) \forall k > 0$ , i.e.,  $e^{u} \in L^{r}(\Omega)$  $\forall r < \infty$ . It follows that  $Ve^{u} \in L^{p-\delta} \forall \delta > 0$  if  $p < \infty$ , and  $Ve^{u} \in L^{r}(\Omega)$  $\forall r < \infty$  if  $p = \infty$ . Standard elliptic estimates imply that  $u \in L^{\infty}(\Omega)$ .

Remark 4. The conclusion of Corollary 2 still holds for a solution u of

$$\begin{cases} -\Delta u = V(x)e^{u} + f(x) & \text{in } \Omega, \\ u = g & \text{on } \partial \Omega \end{cases}$$

,

with  $g \in L^{\infty}(\partial \Omega)$  and  $f \in L^{q}(\Omega)$  for some q > 1. Indeed let w be the solution of

$$\begin{cases} -\Delta w = f \quad \text{in} \quad \Omega, \\ w = g \quad \text{on} \quad \partial \Omega \end{cases}$$

so that  $w \in L^{\infty}(\Omega)$ . The function  $\tilde{u} = u-w$  satisfies

$$\begin{bmatrix} -\Delta \tilde{u} = (Ve^{W})e^{\tilde{u}} & \text{in } \Omega, \\ \tilde{u} = 0 & \text{on } \partial \Omega \end{bmatrix}$$

and we are reduced to the assumptions of Corollary 2.

<u>Remark 5</u>. There is a local version of Corollary 2, namely if  $u \in L^1_{loc}(\Omega)$  satisfies

$$-\Delta u = Ve^{u}$$

with  $V \in L^{p}_{loc}(\Omega)$  and  $e^{u} \in L^{p'}_{loc}(\Omega)$  for some  $1 , then <math>u \in L^{\infty}_{loc}(\Omega)$ . This follows easily from Remark 2.

### П.2. Some variants and counterexamples.

1. The conclusion of Corollary 2 fails when p = 1 (we may only say that  $u^+ \in L^{\infty}(\Omega)$ ). Here is an example:

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Example 1. Let 0 < a < 1. The function

 $u = -a \log(\log \frac{e}{r})$  with r = |x| satisfies

(5) 
$$\begin{cases} -\Delta u = V e^{u} \text{ in } \Omega = B_{1} \\ u = 0 \text{ on } \partial \Omega \end{cases}$$

with  $V = -\frac{a}{r^2(\log \frac{e}{r})^{2-a}}$ . Note that  $V \in L^1(\Omega)$ ,  $e^{U} \in L^{\infty}(\Omega)$  and nevertheless  $u \notin L^{\infty}(\Omega)$  since  $u(x) \to -\infty$  as  $x \to 0$ . The same function u with a < 0provides an example where u satisfies (5) with  $V \in L^1(\Omega)$ ,  $Ve^{U} \in L^1(\Omega)$  and nevertheless  $u^+ \notin L^{\infty}(\Omega)$  since  $u(x) \to +\infty$  as  $x \to 0$ .

2. The function  $e^{u}$  is in some sense the "critical nonlinearity" for which a statement such as Corollary 2 holds. Suppose, for example, that u satisfies

$$\begin{aligned} -\Delta u &= V(x)e^{u^{\alpha}} \quad \text{in} \quad \Omega \\ u &= 0 \qquad \text{on} \quad \partial \Omega \end{aligned}$$

with  $u \ge 0$ ,  $\alpha > 1$ ,  $V \in L^{p}(\Omega)$  and  $e^{u^{\alpha}} \in L^{p'}(\Omega)$ ,  $1 . In general, we may not infer that <math>u \in L^{\infty}(\Omega)$ .

<u>Example 2</u>. Consider first the case  $p = \infty$ . Fix  $1 < \gamma < 2 - (1/\alpha)$ . In  $\Omega = B_1$  set

$$u(\mathbf{x}) = \left| \log(r^2(\log \frac{\mathbf{e}}{r})^{\gamma}) \right|^{1/\alpha}.$$

For r small we have

$$e^{u^{\alpha}} = \frac{1}{r^2(\log \frac{e}{r})}\gamma$$

and therefore  $e^{u^{\alpha}} \in L^{1}(\Omega)$ . On the other hand u satisfies  $-\Delta u = Ve^{u^{\alpha}}$  where V is defined by  $V = (-\Delta u)e^{-u^{\alpha}}$ . An easy computation shows that

V ~ $|\log r|^{\gamma-2+(1/\alpha)}$  as  $r \to 0$ 

and hence  $V \in L^{\infty}(\Omega)$ . Nevertheless  $u \notin L^{\infty}(\Omega)$ .

When 1 we may use the function u above and write

$$-\Delta u = (\mathrm{Ve}^{\frac{1}{p}u^{\alpha}})e^{\frac{1}{p}, u^{\alpha}}$$

The function  $\tilde{u} = (p')^{-1/\alpha}u$  satisfies  $-\Delta \tilde{u} = \tilde{V}e^{\tilde{u}\alpha}$  with  $\tilde{V} = (p')^{-1/\alpha}Ve^{\frac{1}{p}u^{\alpha}}$ so that  $\tilde{V} \in L^{p}(\Omega)$  and  $e^{\tilde{u}\alpha} \in L^{p'}(\Omega)$ .

3. There is a version of Corollary 2 for subsolutions. Assume u satisfies

$$\begin{aligned} -\Delta \mathbf{u} \leq \mathbf{V}(\mathbf{x}) \mathbf{e}^{\mathbf{u}} & \text{in } \Omega , \\ \mathbf{u} \leq \mathbf{0} & \text{on } \partial \Omega , \end{aligned}$$

with  $V \in L^p(\Omega)$  and  $e^u \in L^{p'}(\Omega)$  for some  $1 . Then <math>u^+ \in L^{\omega}(\Omega)$ .

II.3. The case  $\Omega = \mathbb{R}^2$ .

The main result is the following.

<u>Theorem 2</u>. Suppose  $u \in L^{1}_{loc}(\mathbb{R}^{2})$  satisfies

$$-\Delta u = V(x)e^{u}$$
 in  $\mathbb{R}^{2}$ 

with  $V \in L^p(\mathbb{R}^2)$  and  $e^u \in L^{p'}(\mathbb{R}^2)$  for some  $1 . Then <math>u \in L^{\infty}(\mathbb{R}^2)$ .

<u>Proof.</u> Fix  $0 < \epsilon < 1/p'$  and split  $Ve^u$  as  $Ve^u = f_1 + f_2$  with  $\||f_1\||_{L^1(\mathbb{R}^2)} < \epsilon$  and  $f_2 \in L^{\infty}(\mathbb{R}^2)$ . Let  $B_r$  be the ball of radius r centered at  $x_0$ . We denote by C <u>various constants independent of</u>  $x_0$  (but possibly depending on  $\epsilon$ ). Let  $u_i$  be the solution of

$$\begin{cases} -\Delta u_i = f_i & \text{in } B_1, \\ u_i = 0 & \text{on } \partial B_1 \end{cases}$$

By Theorem 1 (applied with  $\delta = 4\pi - 1$ ) we have

$$\int_{B_1} \exp\left[\frac{1}{\epsilon} |u_1|\right] \leq C$$

and in particular  $\|u_1\|_{L^1(B_1)} \leq C$ . We also have  $\|u_2\|_{L^{\infty}(B_1)} \leq C$ . Let

 $u_3 = u - u_1 - u_2$  so that  $\Delta u_3 = 0$  on  $B_1$ . The mean value theorem for harmonic functions implies that

(6) 
$$\|u_{3}^{+}\|_{L^{\infty}(B_{1/2})} \leq C \|u_{3}^{+}\|_{L^{1}(B_{1})}$$

On the other hand we have

$$u_3^+ \le u^+ + |u_1| + |u_2|$$

and since

$$\mathbf{p}' \int_{\mathbf{R}^2} \mathbf{u}^+ \leq \int_{\mathbf{R}^2} \mathbf{e}^{\mathbf{p}'\mathbf{u}} \leq \mathbf{C}$$

we see that  $\|u_3^+\|_{L^1(B_1)} \leq C$ . Combining this with (6) we find that  $\|u_3^+\|_{L^{\infty}(B_{1/2})} \leq C$ . Finally we write

(7) 
$$-\Delta u = V e^{u} = (V e^{u_1}) e^{u_2 + u_3} = g$$

with  $\|g\|_{L^{1+\delta}(B_{1/2})} \leq C$  for some  $\delta > 0$  (since  $e^{u_2+u_3} \in L^{\infty}(B_{1/2})$ ,  $V \in L^{p}(B_1)$  and  $e^{u_1} \in L^{1/\epsilon}(B_1)$  with  $1/\epsilon > p'$ ). Using once more the mean value theorem and standard elliptic estimates we deduce from (7) that

$$\|\mathbf{u}^{+}\|_{L^{\infty}(\mathbf{B}_{1/4})} \leq C \|\mathbf{u}^{+}\|_{L^{1}(\mathbf{B}_{1/2})} + C \|\mathbf{g}\|_{L^{1+\delta}(\mathbf{B}_{1/2})} \leq C.$$

Since C is independent of  $x_0$  we conclude that  $u^+ \in L^{\varpi}(\mathbb{R}^2)$ .

III. Uniform  $L^{\infty}$  bounds and blow-up behavior for solutions of  $-\Delta u = V(x)e^{u}$ . In this section we consider a sequence  $(u_n)$  of solutions of

(8) 
$$-\Delta u_n = V_n(x)e^{u_n} \quad \text{in } \Omega$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^2$ . We seek a uniform bound for  $\|u_n\|_{T^{\overline{\Omega}}}$ 

(resp.  $\|u_n\|_{L^{\overline{w}}_{loc}}$ ) under various assumptions. We start with:

# III.1. Some easy cases

There are two different kinds of assumptions which lead easily to uniform bounds:

- a) Smallness assumption.
- b) Uniform domination.

# a) Smallness assumption

Corollary 3. Assume  $(u_n)$  is a sequence of solutions of (8) with  $u_n = 0$  on  $\partial \Omega$ , such that

(9) 
$$||V_n||_{L^p} \leq C$$
 for some  $1$ 

and

(10) 
$$\int_{\Omega} |V_n| e^{u_n} \leq \epsilon_0 < 4\pi/p' \quad \forall n.$$

Then  $\|u_n\|_{L^{\infty}} \leq C$ .

<u>Proof.</u> Fix  $\delta > 0$  such that  $4\pi - \delta > \epsilon_0(p' + \delta)$ . By Theorem 1 we have

$$\int_{\Omega} e^{(p'+\delta)|u_n|} \leq C.$$

Therefore  $e^{u_n}$  is bounded in  $L^{p'+\delta}(\Omega)$  and so  $V_n e^{u_n}$  is bounded in  $L^q(\Omega)$  for some q > 1. Hence  $u_n$  is bounded in  $L^{\infty}(\Omega)$ .

<u>Remark 6</u>. The smallness condition (10) is sharp. Given any  $1 one can construct a sequence <math>(u_n)$  of solutions of (8) satisfying (9) and

(11) 
$$\int |V_n| e^{u_n} = 4\pi/p$$

such that  $\|u_n\|_{L^{\infty}} \to \infty$ :

Example 3. Set

$$f_{n}(\mathbf{x}) = \begin{cases} \frac{4}{p}, n^{2} & \text{if } |\mathbf{x}| < 1/n, \\ 0 & \text{otherwise} \end{cases}$$

Let u<sub>n</sub> be the solution of

$$\begin{aligned} -\Delta u_n &= f_n & \text{in } B_1, \\ u_n &= 0 & \text{on } \partial B_1. \end{aligned}$$

Note that  $u_n$  satisfies (8) with  $V_n$  being defined by  $V_n = f_n e^{-u_n}$ . An easy computation shows that (9) and (11) hold. Moreover  $||u_n||_{L^{\infty}} = u_n(0) = \frac{1}{n} (2 \log n + 1)$ .

Here is a variant of Corollary 3 where no boundary condition is imposed. <u>Corollary 4</u>. Assume  $(u_n)$  is a sequence of solutions of (8) such that, for some 1 ,

$$\|\mathbf{V}_{\mathbf{n}}\|_{\mathbf{L}^{\mathbf{p}}} \leq \mathbf{C}_{1}$$

$$\|\mathbf{u}_{\mathbf{n}}^{+}\|_{L^{1}} \leq C_{2}$$

and

(14) 
$$\int_{\Omega} |V_n| e^{u_n} \leq \epsilon_0 < 4\pi/p'.$$

Then  $(u_n^+)$  is bounded in  $L_{loc}^{\infty}(\Omega)$ .

<u>Proof.</u> Without loss of generality we may assume that  $\Omega = B_R$ . Split  $u_n$  as  $u_n = u_{1n} + u_{2n}$  where  $u_{1n}$  is the solution of

(15) 
$$\begin{cases} -\Delta u_{1n} = V_n e^{u_n} & \text{in } \Omega, \\ u_{1n} = 0 & \text{on } \partial \Omega; \end{cases}$$

so that  $\Delta u_{2n} = 0$  in  $\Omega$ . By the mean value theorem for harmonic functions we have

$$\| u_{2n}^{+} \|_{L^{\infty}(B_{R/2})} \leq C \| u_{2n}^{+} \|_{L^{1}(B_{R})} \leq C \Big[ \| u_{n}^{+} \|_{L^{1}(B_{R})}^{+} \| \|_{1n} \|_{L^{1}(B_{R})}^{+} \Big]$$
  
 
$$\leq C.$$

Using (15), the smallness condition (14) and Theorem 1 we see that  $(e^{u_{1n}})$  is bounded in  $L^{p'+\delta}(B_R)$  for some  $\delta > 0$ . Therefore  $(V_n e^{u_n})$  is bounded in  $L^q(B_{R/2})$  for some q > 1. Using (15) once more we see that  $(u_{1n})$  is bounded in  $L^{\varpi}(B_{R/4})$ . Therefore  $(u_n)$  is bounded in  $L^{\varpi}(B_{R/4})$ .

b) Uniform domination

<u>Corollary</u> 5. Assume  $(u_n)$  is a sequence of solutions of (8) with  $u_n = 0$  on  $\partial \Omega$ , satisfying, for some 1 ,

$$\|\mathbf{e}^{\mathbf{u}_{\mathbf{n}}}\|_{\mathbf{L}^{\mathbf{p}'}} \leq \mathbf{C}$$

and one of the following conditions:

either

(17) 
$$|V_n(\mathbf{x})| \leq W(\mathbf{x}) \quad \forall \mathbf{n}, \text{ with } W \in L^p(\Omega)$$

or

(18) 
$$V_n \to V \text{ in } L^p(\Omega).$$

Then  $\|\mathbf{u}_{\mathbf{n}}\|_{\mathbf{L}^{\infty}} \leq \mathbf{C}.$ 

<u>Proof.</u> Assume first that (17) holds. For every  $\epsilon > 0$  we have

$$|V_n|e^{u_n} \leq We^{u_n} \leq \epsilon e^{p'u_n} + \frac{1}{\epsilon^{1/(p-1)}}W^p.$$

By (16) we may fix  $\epsilon > 0$  small enough so that

(19) 
$$\epsilon \int_{\Omega} e^{p' u_n} \leq \alpha < 4\pi/p' \quad \forall n.$$

We have  $|u_n| \leq u_{1n} + u_2$  where  $u_{1n}$  is the solution of

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$$\begin{bmatrix} -\Delta u_{1n} = \epsilon e^{p' u_n} & \text{in } \Omega, \\ u_{1n} = 0 & \text{on } \partial \Omega \end{bmatrix}$$

and u<sub>2</sub> is the solution of

$$\begin{cases} -\Delta u_2 = \frac{1}{\epsilon^1/(p-1)} W^p & \text{in } \Omega, \\ u_2 = 0 & \text{on } \partial\Omega \end{cases}$$

By Theorem 1 and (19) we see that  $e^{u_{1n}}$  is bounded in  $L^{p'+\delta}(\Omega)$  for some  $\delta > 0$  and, by Corollary 1,  $e^{u_2} \in L^k(\Omega)$  for every  $k \ge 1$ . Thus  $|V_n|e^{u_n} \le e^{u_{1n}} (e^{u_{2n}}W)$  remains bounded in some  $L^q$ , q > 1, and the conclusion follows.

Assume now that (18) holds. Suppose, by contradiction, that  $\|u_n\|_{L^{\infty}}$  is not bounded. We may then extract a subsequence such that  $\|u_n\|_{k^{\infty} \to \infty}$ . By passing to a further subsequence (still denoted  $n_k$ ) we may assume that  $|V_{n_k}| \leq W$  for some  $W \in L^p$  (see e.g. [1]), Théorème IV.9). We are therefore reduced to the previous case.

# III.2. The main results

We now turn to the study of a sequence  $(u_n)$  of solutions of (8) under the assumptions

(20) 
$$V_n \ge 0 \text{ in } \Omega, \|V_n\|_{L^p} \le C_1 \text{ and } \|e^{u_n}\|_{L^{p'}} \le C_2$$

for some 1 . A typical example is the sequence

$$u_n(x) = \log \frac{8n^2}{(1+n^2|x|^2)^2}$$

which satisfies  $-\Delta u_n = e^{u_n}$  in  $\mathbb{R}^2$  and  $\||e^{u_n}||_{L^1(\mathbb{R}^2)} = 8\pi$ . Note that  $u_n(x) \to -\infty$  for all  $x \neq 0$  and  $u_n(0) \to +\infty$ . This example provides a very good description of the blow-up mechanism in the general case under the assumption (20). In fact, if a sequence  $(u_n)$  becomes unbounded then there is a

finite set S (possibly empty) where  $u_n$  tends to  $+\infty$  and elsewhere  $u_n$  tends to  $-\infty$ .

More precisely, define the "blow-up" set as follows:

$$S = \left\{ x \in \Omega; \text{ there exists a sequence } x_n \text{ in } \Omega \text{ such that } x_n \to x \right\}$$
  
and  $u_n(x_n) \to +\infty$ 

Then we have

<u>Theorem 3.</u> Assume  $(u_n)$  is a sequence of solutions of (8) satisfying, for some 1 ,

(21) 
$$V_n \ge 0 \text{ in } \Omega$$

$$\|\mathbf{V}_{\mathbf{n}}\|_{\mathbf{r},\mathbf{p}} \leq C_{1}$$

and

$$\|e^{\mathbf{u}_{\mathbf{n}}}\|_{\mathbf{L}^{\mathbf{p}'}} \leq C_2$$

Then, there exists a subsequence  $(u_{n_k})$  satisfying the following alternative: either

(i) 
$$(u_{n_k})$$
 is bounded in  $L_{loc}^{\infty}(\Omega)$ 

or

(ii)

$$u_{n_k}(x) \longrightarrow -\infty$$
 uniformly on compact subsets of  $\Omega$ 

or

(iii) the blow-up set S (relative to  $(u_{n_k})$ ) is finite, nonempty and  $u_{n_k}(\mathbf{x}) \rightarrow -\infty$  uniformly on compact subsets of  $\Omega \setminus S$ . In addition  $V_{n_k} e^{u_{n_k}}$ converges in the sense of measures on  $\Omega$  to  $\sum_{i} \alpha_i \delta_{a_i}$  with  $\alpha_i \ge 4\pi/p' \quad \forall i$  and  $S = \bigcup_i \{a_i\}.$ 

Before giving the proof of Theorem 3 we mention some Corollaries. <u>Corollary 6</u>. Assume  $(u_n)$  is a sequence of solutions of (8) with  $u_n = 0$  on  $\partial\Omega$ , satisfying (21), (22), and (23). Then  $(u_n)$  is bounded in  $L_{loc}^{m}(\Omega)$ .

<u>Proof.</u> By the maximum principle  $u_n \ge 0$  on  $\Omega$  and therefore cases (ii) and (iii) in Theorem 3 are excluded for all subsequences. Therefore the (full) sequence  $(u_n)$  is bounded in  $L_{loc}^{\Theta}(\Omega)$ .

<u>Remark 7</u>. One may wonder whether the conclusion of Corollary 6 holds uniformly up to the boundary (since we impose here the boundary condition  $u_n = 0$  on  $\partial \Omega$ ). This is not true as is shown in Section III.3 (Example 6). However it is plausible that a stronger assumption about the  $V'_n$ 's yields an estimate up to the boundary. For example, here is an

<u>Open problem 1:</u> Suppose  $(u_n)$  is a sequence of solutions of (8) with  $u_n = 0$  on  $\partial \Omega$  satisfying (21),

(24) 
$$V_n \to V \text{ in } C^{\circ}(\overline{\Omega})$$

and

 $\|e^{\mathbf{u}_{\mathbf{n}}}\|_{\mathbf{r}^{1}} \leq \mathbf{C}.$ 

Can one conclude that  $\|u_n\|_{L^{\infty}} \leq C$ ?

<u>Remark 8</u>. The conclusion of Corollary 6 also fails if we remove assumption (21) (i.e.  $V_n \ge 0$  on  $\Omega$ ); see [8].

Another obvious consequence of Theorem 3 is:

<u>Corollary 7</u>. Assume  $(u_n)$  is a sequence of solutions of (8) satisfying (21), (22) and (23). Assume in addition

(26)  $u_n \ge -M$  in  $\Omega, \forall n$ 

for some positive constant M, or more generally

 $\|\mathbf{u}_{\mathbf{n}}^{-}\|_{\mathbf{T}^{1}} \leq \mathbf{M} \quad \forall \mathbf{n} \,.$ 

Then  $(u_n)$  is bounded in  $L_{loc}^{\omega}(\Omega)$ . <u>Corollary 8</u>. Assume  $(u_n)$  is a sequence of solutions of (8) satisfying (26) and (28)  $0 < a \le V_n \le b < \infty$  in  $\Omega$ 

for some constants a, b.

Then  $(u_n)$  is bounded in  $L_{loc}^{o}(\Omega)$ .

<u>Proof.</u> In view of Corollary 7 we have only to show that  $(e^{U_n})$  is bounded in  $L^1_{loc}(\Omega)$ . We may always assume that M = 0, i.e.  $u_n \ge 0$  (this amounts to replace  $u_n$  by  $u_n + M$ ). Let  $\varphi_1$  be the first eigenfunction of  $-\Delta$  on  $\Omega$  with zero Dirichlet conditions and let  $\lambda_1$  be the corresponding eigenvalue. Multiplying (8) by  $\varphi_1$  and integrating we obtain

$$\int_{\Omega} V_{n} e^{u_{n}} \varphi_{1} = \int_{\partial \Omega} u_{n} \frac{\partial \varphi_{1}}{\partial \nu} + \lambda_{1} \int_{\Omega} u_{n} \varphi_{1}$$

where  $\nu$  is the outward normal. Using (28),  $u_n \ge 0$  and  $\frac{\partial \varphi_1}{\partial \nu} \le 0$ we obtain

$$a \int_{\Omega} e^{u_n} \varphi_1 \leq \lambda_1 \int_{\Omega} u_n \varphi_1$$

This provides an upper bound for  $\int_{\Omega} e^{u_n} \varphi_1$ . Therefore  $(e^{u_n})$  is bounded in  $L^1_{loc}(\Omega)$  and the conclusion follows.

Remark 9. There are two natural questions suggested by Corollary 8:

<u>Open problem 2</u>: Suppose  $(u_n)$  is a sequence of solutions of (8) with  $u_n = 0$  on  $\partial \Omega$  satisfying (28). Can one conclude that  $(u_n)$  is bounded in  $L^{\infty}(\Omega)$ ? Is this true if we assume in addition that  $\|e^{u_n}\|_{r^1} \leq C$ ?

<u>Open problem 3</u>: Assume  $(u_n)$  is a sequence of solutions of (8) satisfying (26) and (28). Let K be a compact subset of  $\Omega$ . What is the optimal bound for  $\sup_{n \to \infty} u_n$  as a function of M? Does one have

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(29) 
$$\begin{array}{c} \sup_{\mathbf{K}} u_n \leq C_1 M + C_2 \end{array}$$

for some positive constants  $C_1$ ,  $C_2$  depending only on a, b, K and  $\Omega$ ? Can one take  $C_1 = 1$  if  $V_n(x) \equiv 1$ ? [Note that (29) holds with  $C_1 = 1$  for the special sequence  $u_n(x) = \log \frac{8n^2}{(1+n^2|x|^2)^2}$ ].

<u>Proof of Theorem 3</u>. Since  $(V_n e^{u_n})$  is bounded in  $L^1(\Omega)$  we may extract a subsequence (still denoted  $V_n e^{u_n}$ ) such that  $V_n e^{u_n}$  converges in the sense of measures on  $\Omega$  to some nonnegative bounded measure  $\mu$ , i.e.

(30) 
$$\int V_{\mathbf{n}} e^{\mathbf{u}_{\mathbf{n}}} \ \psi \rightarrow \int \psi \mathrm{d}\mu$$

for every  $\psi \in C_c(\Omega)$ .

It follows from Corollary 4 (applied in a small ball around  $x_0$ ) that if  $x_0$  is a regular point then there is some  $R_0 > 0$  such that

(32) 
$$(u_n^+)$$
 is bounded in  $L^{\infty}(B_{R_0}(x_0))$ .

Note that (13) holds since  $(e^{u_n})$  is bounded in  $L^{p'}(\Omega)$ .

We denote by  $\Sigma$  the set of nonregular points in  $\Omega$ . Clearly  $x_0 \in \Sigma$  iff  $\mu(\{x_0\}) \ge 4\pi/p'$ . Since  $\mu$  is a bounded measure (with  $\int d\mu \le C_1C_2$ ) it follows that  $\Sigma$  is finite and

card (
$$\Sigma$$
)  $\leq C_1 C_2 p'/4\pi$ .

We now split the proof of Theorem 3 into 3 steps.

<u>Step 1</u>:  $S = \Sigma$ .

Clearly  $S \in \Sigma$  by (32). Conversely, suppose  $x_0 \in \Sigma$ . Then we have

(33) 
$$\forall \mathbf{R} > 0, \ \lim \|\mathbf{u}_{\mathbf{n}}^{+}\|_{\mathbf{L}^{\mathbf{m}}(\mathbf{B}_{\mathbf{R}}(\mathbf{x}_{0}))} = + \mathbf{w}.$$

Otherwise there would be some  $R_0 > 0$  and a subsequence such that

$$\|u_{\mathbf{n}_{\mathbf{k}}}^{+}\|_{L^{\infty}(B_{\mathbf{R}_{0}}(\mathbf{x}_{0}))} \leq C. \text{ In particular } \|e^{u_{\mathbf{n}_{\mathbf{k}}}}\|_{L^{\infty}(B_{\mathbf{R}_{0}}(\mathbf{x}_{0}))} \leq C \text{ and therefore }$$
$$\int_{B_{\mathbf{R}}(\mathbf{x}_{0})} V_{\mathbf{n}_{\mathbf{k}}} e^{u_{\mathbf{n}_{\mathbf{k}}}} \leq CC_{1} \mathbf{R}^{2/p'} \text{ for all } \mathbf{R} < \mathbf{R}_{0}.$$

This implies (31) for some suitable  $\psi$ . Therefore  $\mathbf{x}_0$  is regular - a contradiction. Hence we have established (33). Choose  $\mathbf{R} > 0$  small enough so that  $\mathbf{B}_{\mathbf{R}}(\mathbf{x}_0)$ does not contain any other point of  $\Sigma$ . Let  $\mathbf{x}_n \in \mathbf{B}_{\mathbf{R}}(\mathbf{x}_0)$  be such that

$$u_n^+(x_n) = \max_{B_R}(x_0) u_n^+ \to +\infty$$

We claim that  $x_n \to x_0$ . Otherwise there would be a subsequence  $x_{n_k} \to \overline{x} \neq x_0$  and  $\overline{x} \notin \Sigma$ , i.e.  $\overline{x}$  is a regular point. This is impossible in view of (32). Hence we have established that  $x_0 \in S$ . This completes the proof of Step 1.

<u>Step 2</u>:  $S = \phi$  implies (i) or (ii) holds.

By (32)  $(u_n^+)$  is bounded in  $L^{\infty}_{loc}(\Omega)$  and therefore  $f_n = V_n e^{u_n}$  is bounded in  $L^{p}_{loc}(\Omega)$ . This implies that  $\mu \in L^1(\Omega) \cap L^{p}_{loc}(\Omega)$ . Let  $v_n$  be the solution of

$$\begin{cases} -\Delta \mathbf{v}_{\mathbf{n}} = \mathbf{f}_{\mathbf{n}} & \text{in } \Omega, \\ \mathbf{v}_{\mathbf{n}} = 0 & \text{on } \partial \Omega \end{cases}$$

Clearly,  $v_n \to v$  uniformly on every compact subset of  $\Omega,$  where v is the solution of

$$\begin{bmatrix} -\Delta \mathbf{v} = \mu & \text{in } \Omega, \\ \mathbf{v} = 0 & \text{on } \partial \Omega. \end{bmatrix}$$

Let  $w_n = u_n - v_n$  so that  $\Delta w_n = 0$  on  $\Omega$  and  $w_n^+$  is bounded in  $L_{loc}^{\alpha}(\Omega)$ . By Harnack's principle we find that: either

(34) a subsequence  $(w_{n_k})$  is bounded in  $L^{\infty}_{loc}(\Omega)$ ,

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or

(35)  $(w_n)$  converges uniformly to  $-\infty$  on compact subsets of  $\Omega$ .

Case (i) corresponds to (34) and case (ii) to (35).

<u>Step 3</u>. S  $\neq \phi$  implies (iii) holds.

By (32)  $(u_n^+)$  is bounded in  $L_{loc}^{\infty}(\Omega \setminus S)$  and therefore  $f_n = V_n e^{u_n}$  is bounded in  $L_{loc}^{\infty}(\Omega \setminus S)$ . This implies that  $\mu$  is a bounded measure on  $\Omega$  with  $\mu \in L_{loc}^{p}(\Omega \setminus S)$ . (A basic difference with Step 2 is that here  $\mu$  is a measure, not an  $L^1$  function, and, as will be shown later,  $\mu$  is a sum of Dirac masses). Let  $v_n$ , v and  $w_n$  be defined as in Step 2. Then  $v_n \to v$  uniformly on compact subsets of  $\Omega \setminus S$ . As above, by Harnack's principle, we find that either

(36) a subsequence  $(w_{n_k})$  is bounded in  $L^{\infty}_{loc}(\Omega \setminus S)$ 

or

(37)  $(w_n)$  converges to  $-\infty$  uniformly on compact subsets of  $\Omega \setminus S$ .

We claim that (36) does not happen. Fix some point  $x_0 \in S$  and R > 0 small so that  $x_0$  is the only point of S in  $\mathbb{B}_R(x_0)$ . Assume (36) holds, so that  $(w_{n_k})$  is bounded in  $L^{\infty}(\partial B_R(x_0))$  and similarly for  $(v_n)$ . Therefore  $(u_{n_k})$  is bounded in  $L^{\infty}(\partial B_R(x_0))$ , say by C. Let  $z_{n_k}$  be the solution of

$$\begin{bmatrix} -\Delta z_{\mathbf{n}_{\mathbf{k}}} &= \mathbf{f}_{\mathbf{n}_{\mathbf{k}}} & \text{in } \mathbf{B}_{\mathbf{R}}(\mathbf{x}_{0}), \\ \mathbf{z}_{\mathbf{n}_{\mathbf{k}}} &= -\mathbf{C} & \text{on } \partial \mathbf{B}_{\mathbf{R}}(\mathbf{x}_{0}). \end{bmatrix}$$

By the maximum principle  $u_{n_k} \ge z_{n_k}$  in  $B_R(x_0)$ . In particular

(38)  $\int e^{p' \mathbf{z}_{\mathbf{n}_{k}}} \leq \int e^{p' \mathbf{u}_{\mathbf{n}_{k}}} \leq C_{2}^{p'}.$ 

On the other hand  $z_{n_k} \to z$  a.e. (even uniformly on compact subsets of  $B_R(x_0) \setminus \{x_0\}$ ) where z is the solution of

$$\begin{aligned} -\Delta \mathbf{z} &= \mu \quad \text{in } \mathbf{B}_{\mathbf{R}}(\mathbf{x}_0), \\ \mathbf{z} &= -\mathbf{C} \quad \text{on } \quad \partial \mathbf{B}_{\mathbf{R}}(\mathbf{x}_0). \end{aligned}$$

Finally note that since  $\mathbf{x}_0 \in S$  is not a regular point we have  $\mu(\{\mathbf{x}_0\}) \ge 4\pi/p'$ . This implies that  $\mu \ge \frac{4\pi}{p'} \delta_{\mathbf{x}_0}$  and therefore

$$z(\mathbf{x}) \geq \frac{2}{p'} \log \frac{1}{|\mathbf{x} - \mathbf{x}_0|} + 0(1) \text{ as } \mathbf{x} \to \mathbf{x}_0$$

Thus  $e^{p'z} \ge C/|x-x_0|^2$  with C > 0. Hence  $\int_{B_R(x_0)} e^{p'z} = \infty$ . On the other

hand, by (38) and Fatou's Lemma we find that

$$\int e^{p'z} \leq C_2^{p'}.$$

A contradiction. Hence we have shown that (37) holds. Consequently  $(u_n)$  converges to  $-\infty$  uniformly on compact subsets of  $\Omega \setminus S$ . Therefore  $V_n e^{u_n} \longrightarrow 0$  in  $L^p_{1oc}(\Omega \setminus S)$  and hence  $\mu$  is supported on S. This means that  $\mu = \sum_i \alpha_i \delta_{a_i}$  with  $S = \bigcup_i \{a_i\}$ . The argument above gives that  $\alpha_i \ge 4\pi/p'$  for each i.

<u>Remark 10</u>. The conclusion (iii) in Theorem 3 involves a finite sum of Dirac masses  $\sum_{i} \alpha_i \delta_{a_i}$  with coefficients  $\alpha_i \ge 4\pi/p'$ . The  $\alpha_i$ 's as well as the  $a_i$ 's can be chosen arbitrarily. More precisely given any finite set  $S = \bigcup_{i=1}^{k} \{a_i\}$  and any  $\alpha_i > 4\pi/p'$  there exist sequences  $(u_n)$  and  $(V_n)$  as in Theorem 3 such that  $V_n e^{u_n}$  converges to  $\sum_{i=1}^{k} \alpha_i \delta_{a_i}$ .

To construct such sequences we proceed as follows. Set, for  $1 \leq i \leq k$ ,

$$\mathbf{v}_{i,n} = \begin{cases} -\frac{A_i}{4} n^{2\beta_i} |\mathbf{x} - \mathbf{a}_i|^2 + \frac{A_i}{4} & \text{if } |\mathbf{x} - \mathbf{a}_i| < 1/n^{\beta_i}, \\ \frac{A_i}{2} \log(\frac{1}{n^{\beta_i} |\mathbf{x} - \mathbf{a}_i|}) & \text{if } |\mathbf{x} - \mathbf{a}_i| \ge 1/n^{\beta_i} \end{cases}$$

where  $A_i = \alpha_i/\pi > 4/p'$  and  $\beta_i$  is defined by the relation  $\beta_i(\frac{A_i}{2} - \frac{2}{p'}) = 1$ . Let  $u_n = \sum_{i=1}^k v_{i,n} + \sigma_n$  where  $\sigma_n = ((k-1) + \frac{2}{p'} \sum_{i=1}^k \beta_i)\log n$ . A direct computation shows that  $V_n = (-\Delta u_n)e^{-u_n}$  satisfies (21) and (22); moreover  $(e^{u_n})$ is bounded in  $L^{p'}$  and  $V_n e^{u_n}$  converges to  $\sum_{i=1}^k \alpha_i \delta_{a_i}$ .

We believe that under additional conditions on the  $V_n$ 's the  $\alpha_i$ 's in Theorem 3 cannot take arbitrary values  $(> 4\pi/p')$ : <u>Open problem 4</u>: Assume  $(u_n)$  is a sequence of solutions of (8) satisfying  $V_n \ge 0$  on  $\Omega$ ,  $V_n \longrightarrow V$  uniformly in  $\overline{\Omega}$  with  $V_n$ ,  $V \in C^0(\overline{\Omega})$  and  $\|e^{u_n}\|_{\tau^1} \le C$ . Assume  $S \neq \phi$  so that case (iii) holds. Can one conclude that

Evidence in favor of a positive answer comes from the fact that after a blow-up near  $a_i$  we are led to a solution of  $-\Delta v = ce^v$  on  $\mathbb{R}^2$  with  $c = V(a_i)$  and  $\int_{\mathbb{R}^2} e^v < \infty$ . It follows from the result of [3] that  $\int_{\mathbb{R}^2} ce^v = 8\pi$ . On the other hand, the blow-up analysis gives (formally)  $\alpha_i = \int_{\mathbb{R}^2} ce^v$ .

In Theorem 3 the assumption  $\|e^{u_n}\|_{L^{p'}} \leq C$  provides some kind of <u>bound</u> from above for  $(u_n)$  and plays an important role in proving that the blow-up set S is finite. If we drop that assumption little can be said in the general case. For instance, we may have a sequence  $(u_n)$  of solutions of

$$-\Delta u_n = e^{u_n}$$
 on  $\Omega$ 

 $(\text{with } \|e^{u_n}\|_{L^1} \to \infty) \text{ such that }$ 

$$\begin{cases} u_n \longrightarrow +\infty & \text{ on a line } S, \\ u_n \longrightarrow -\infty & \text{ in } \Omega \backslash S. \end{cases}$$

Example 4. The sequence

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$$u_n(x,y) = 2nx - 2 \log (1 + e^{2nx}) + \log 8n^2$$

satisfies  $-\Delta u_n = e^{u_n}$ ,  $u_n(0,y) \rightarrow +\infty$  and  $u_n(x,y) \rightarrow -\infty$  for  $x \neq 0$ . However, if we assume some <u>bound from below</u> for the  $u_n$ 's then there are

only two possibilities: either  $S = \Omega$  (total blow-up) or S is (locally) finite.

<u>Theorem 4</u>. Assume  $(u_n)$  is a sequence of solutions of (8) satisfying, for some 1 , (21), (22) and

$$\|\mathbf{u}_{\mathbf{n}}^{-}\|_{-1} \leq C$$

Then, there exists a subsequence  $(u_{n_k})$  satisfying the following alternative: either

(i)  $u_{n_k} \longrightarrow +\infty$  uniformly on compact subsets of  $\Omega$ 

or

(ii) the blow-up set S (relative to  $(u_{n_k})$ ) is locally finite (i.e. for each  $x \in \Omega$ there is some neighborhood N(x) of x such that N(x)  $\cap$  S is finite). Moreover  $(u_{n_n})$  is bounded in  $L^{\infty}_{loc}(\Omega \setminus S)$ .

Remark 11. Both cases in the alternative may occur:

Example of (i). Let v be any solution of  $-\Delta v = e^{v}$  in  $\mathbb{R}^2$ . Then  $u_n = v+n$  satisfies  $-\Delta u_n = V_n e^{u_n}$  with  $V_n = e^{-n}$  and  $u_n \to +\infty$  everywhere. Example of (ii). Recall that  $v_n(x) = \log \frac{8n^2}{(1+n^2|x|^2)^2}$  satisfies  $-\Delta v_n = e^{v_n}$ . Thus  $u_n = v_n + \log n^2$  satisfies  $-\Delta u_n = V_n e^{u_n}$  with  $V_n = 1/n^2$ . Note that  $u_n(0) \to +\infty$  while  $u_n(x)$  remains bounded for  $x \neq 0$ .

Proof of Theorem 4. Without loss of generality we may assume that

(40) 
$$u_n \ge 0$$
 in  $\Omega$ .

Indeed, by Kato's inequality [5] we have

(41) 
$$\Delta u_n \ge -(\Delta u_n)\chi([u_n \le 0]) = V_n e^{u_n}\chi([u_n \le 0]) \ge -|V_n|.$$

It follows from (39), (41) and standard elliptic estimates that  $(u_n)$  is bounded in

 $L_{loc}^{\alpha}(\Omega)$ . Passing to a smaller domain and adding a constant to  $(u_n)$  we may always assume that (40) holds.

We now split the proof into 3 cases.

<u>Case 1</u>: There exists a compact subset K c  $\Omega$  and a subsequence  $(u_{n_k})$  such that

(42) 
$$\int_{\mathbf{K}} \mathbf{V}_{\mathbf{n}_{\mathbf{k}}} \overset{\mathbf{u}_{\mathbf{n}_{\mathbf{k}}}}{\longrightarrow} + \boldsymbol{\omega} \, .$$

Then (i) holds.

Indeed, let K' be any compact subset of  $\Omega$ . Using (40) we obtain

$$\mathbf{u}_{\mathbf{n}_{k}}(\mathbf{x}) \geq \int_{\Omega} \mathbf{G}(\mathbf{x},\mathbf{y}) \mathbf{V}_{\mathbf{n}_{k}}(\mathbf{y}) \mathbf{e}^{\mathbf{u}_{\mathbf{n}_{k}}(\mathbf{y})} \mathrm{d}\mathbf{u},$$

where G is the Green's function of  $-\Delta$  with Dirichlet condition on  $\partial\Omega$ . Since  $G(x,y) \geq \alpha > 0 \quad \forall x \in K', \forall y \in K$  we see that, for  $x \in K'$ ,

$$u_{\mathbf{n}_{\mathbf{k}}}(\mathbf{x}) \geq \alpha \int_{\mathbf{K}} V_{\mathbf{n}_{\mathbf{k}}} e^{u_{\mathbf{n}_{\mathbf{k}}}} \rightarrow + \omega$$

<u>Case 2</u>.  $(V_n e^{u_n})$  is bounded in  $L^1_{loc}(\Omega)$  and there exists a compact subset K  $\subset \Omega$  such that, for a subsequence,

$$\int_{\mathbf{K}} \mathbf{u}_{\mathbf{n}_{\mathbf{k}}} \to + \infty .$$

Then (i) holds.

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Indeed, let K' be any compact subset of  $\Omega$ . Let  $\omega$  be an open set such that  $K \cup K' \subset \omega \subset \Omega$ . In  $\omega$ , split  $u_n$  as  $u_n = u_{1n} + u_{2n}$  where  $u_{1n}$  is the solution of

$$\begin{bmatrix} -\Delta u_{1n} = V_n e^{u_n} & \text{in } \omega, \\ u_{1n} = 0 & \text{on } \partial \omega. \end{bmatrix}$$

Note that  $(u_{1n})$  is bounded in  $L^1(\omega)$  and  $u_{2n}$  satisfies

 $\begin{cases} -\Delta u_{2n} = 0 \quad \text{in} \quad \omega, \\ u_{2n} \ge 0 \quad \text{on} \quad \partial \omega \end{cases}$ 

Thus  $u_{2n} \ge 0$  in  $\omega$  and by Harnack's principle

(43) 
$$\begin{array}{c} \sup_{K \cup K'} u_{2n} \leq C \quad \inf_{K \cup K'} u_{2n} \leq C \quad \inf_{K'} u_{n} \\ K \cup K' \quad K' \quad K' \end{array}$$

On the other hand

$$\begin{array}{c} u_{2n} \leq C \quad Sup \quad u_{2n} \leq C \quad Sup \quad u_{2n} \\ K \quad K \cup K' \end{array}$$

and

$$\int_{\mathbf{K}} \mathbf{u}_{2n} = \int_{\mathbf{K}} \mathbf{u}_n - \int_{\mathbf{K}} \mathbf{u}_{1n} \ge \int_{\mathbf{K}} \mathbf{u}_n - \mathbf{C}$$

It follows that  $\inf_{K'} u_{n_k} \to +\infty$  and thus (i) holds.

We are left with:

<u>Case 3</u>:  $(V_n e^{u_n})$  and  $(u_n)$  are bounded in  $L^1_{loc}(\Omega)$ . Then (ii) holds.

We proceed here as in the proof of Theorem 3. We extract a subsequence (still denoted  $V_n e^{u_n}$ ) such that  $V_n e^{u_n}$  converges in the sense of measures to some nonnegative (possibly unbounded) measure  $\mu$ , i.e.

$$\int V_{\mathbf{n}} e^{\mathbf{u}_{\mathbf{n}}} \psi \longrightarrow \int \psi d\mu$$

for every  $\psi \in C_c(\Omega)$ . We say that a point  $x_0 \in \Omega$  is a <u>regular point</u> if there is a function  $\psi \in C_c(\Omega)$ ,  $0 \le \psi \le 1$ , with  $\psi = 1$  in some neighborhood of  $x_0$ , such that

$$\int \psi d\mu < 4\pi/p'$$

It follows from Corollary 4 (applied in a small ball around  $x_0$ ) that if  $x_0$  is a regular point then there is some  $R_0 > 0$  such that

(44) 
$$(u_n)$$
 is bounded in  $L^{\omega}(B_{R_0}(x_0))$ .

We denote by  $\Sigma$  the set of nonregular points in  $\Omega$ . Clearly  $x_0 \in \Sigma$  if  $\omega(\{x_0\}) \ge 4\pi/p'$ . It follows that  $\Sigma$  is locally finite and for every compact subset K of  $\Omega$ 

$$\operatorname{card}(\Sigma \cap K) \leq (p'/4\pi) \int_{K} d\mu$$

We have  $S = \Sigma$  as in the proof of Theorem 3 (Step 1). Thus S is locally finite and by (44)  $(u_n)$  is bounded in  $L^{\infty}_{loc}(\Omega \setminus S)$ , i.e. (ii) holds.

# III.3. Variants and counterexamples

1. Suppose that instead of a sequence of solutions of (8) we have a sequence of <u>subsolutions</u>, i.e.

$$-\Delta u_n \leq V_n(x)e^{u_n}$$
 in  $\Omega$ .

It is easy to adapt the arguments of Section III.1 to obtain estimates for  $\||\mathbf{u}_{\mathbf{n}}^+||_{\mathbf{L}^{\infty}}$ under smallness or uniform domination assumption. However the analogue of Corollary 6 for subsolutions does not hold as may be seen from the following:

<u>Example 5</u>. There is a sequence  $(u_n)$  satisfying

$$\begin{cases} -\Delta u_n \leq e^{u_n} & \text{in } \Omega = B_1, \\ u_n = 0 & \text{on } \partial \Omega \end{cases}$$

with

$$\int_{\Omega} e^{u_n} \leq C$$

and such that  $u_n(0) \rightarrow +\infty$ . First, note that the function

$$\varphi_{\epsilon}(\mathbf{x}) = \log \frac{8\epsilon^2}{(\epsilon^2 + |\mathbf{x}|^2)^2}$$

satisfies

$$-\Delta\varphi_{\epsilon} = e^{\varphi_{\epsilon}} \quad \forall \epsilon > 0$$

and

$$\int_{\mathbb{R}^2} e^{\varphi_{\epsilon}} = 8\pi \quad \forall \epsilon > 0.$$

Hence the function  $u_n = \varphi_{1/n}^+$  has all the required properties. The same example can be used to produce sequences  $(v_n)$  and  $(V_n)$  such that

$$\begin{aligned} -\Delta \mathbf{v}_{\mathbf{n}} &\leq \mathbf{V}_{\mathbf{n}} \mathbf{e}^{\mathbf{v}_{\mathbf{n}}} \quad \text{in} \quad \Omega = \mathbf{B}_{1} \\ \mathbf{v}_{\mathbf{n}} &= 0 \qquad \text{on} \quad \partial \Omega \end{aligned}$$

such that  $V_n \ge 0$ ,  $||V_n||_{L^p} \le C$ ,  $||e^{v_n}||_{L^{p'}} \le C$ ,  $1 , and <math>v_n(0) \rightarrow +\infty$ . It suffices to take  $v_n = \frac{1}{p}$ ,  $u_n$  and  $V_n = \frac{1}{p}$ ,  $e^{\frac{1}{p}u_n}$ .

2. The same kind of example shows that the conclusion of Theorem 2 does <u>not</u> hold <u>uniformly</u>. More precisely there are sequences  $(u_n)$  and  $(V_n)$  such that

$$-\Delta u_n = V_n e^{u_n}$$
 on  $\mathbb{R}^2$ 

with  $\|V_n\|_{L^p(\mathbb{R}^2)} \leq C$ ,  $\|e^{v_n}\|_{L^{p'}(\mathbb{R}^2)} \leq C$ ,  $1 , such that <math>u_n(0) \rightarrow +\infty$ . One may take for instance  $u_n = \frac{1}{p}, \varphi_{1/n}$  and  $V_n = \frac{1}{p}, \exp(\frac{1}{p} \varphi_{1/n})$ .

3. The conclusion of Corollary 6 cannot be strengthened to  $\|u_n\|_{L^{\infty}} \leq C$ . There are sequences  $(u_n)$  and  $(V_n)$  satisfying

$$\begin{aligned} -\Delta \mathbf{u}_{\mathbf{n}} &= \mathbf{V}_{\mathbf{n}} \mathbf{e}^{\mathbf{u}_{\mathbf{n}}} & \text{in } \Omega &= \mathbf{B}_{1} \\ \mathbf{u}_{\mathbf{n}} &= 0 & \text{on } \partial \Omega \\ \mathbf{V}_{\mathbf{n}} &\geq 0 & \text{in } \Omega \\ & \left\| \mathbf{V}_{\mathbf{n}} \right\|_{\mathbf{L}^{\mathbf{p}}} \leq \mathbf{C}, \\ & \left\| \mathbf{e}^{\mathbf{u}_{\mathbf{n}}} \right\|_{\mathbf{L}^{\mathbf{p}'}} \leq \mathbf{C}, \end{aligned}$$

with  $1 and such that <math>\|u_n\|_{L^{\infty}} \to \infty$ . It suffices to construct such an example when  $p = \infty$ . For a general  $1 we may use the <math>p = \infty$  example and note that  $\tilde{u}_n = \frac{1}{p}, u_n$  satisfies  $-\Delta \tilde{u}_n = \tilde{V}_n e^{\tilde{u}_n}$  with  $\tilde{V}_n = \frac{1}{p}, V_n \exp(\frac{1}{p}, u_n)$  so that  $\|\tilde{V}_n\|_{L^p} \leq C$  and  $\|e^{\tilde{u}_n}\|_{L^{p'}} \leq C$ .

Example 6. Let  $\Omega$  be the unit disc centered at (1,0). Set  $a_{\epsilon} = (d_{\epsilon}, 0)$  with

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 $\epsilon < d_{\epsilon} < 1$ . Let A > 1 be a constant and let

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$$\mathbf{f}_{\epsilon} = \begin{cases} \frac{4\mathbf{A}}{\epsilon^2} & \text{in } \mathbf{B}_{\epsilon}(\mathbf{a}_{\epsilon}) \\ 0 & \text{otherwise} \end{cases}.$$

Let  $u_{\epsilon}$  be the solution of

$$\begin{split} -\Delta \mathbf{u}_{\epsilon} &= \mathbf{f}_{\epsilon} \quad \text{in} \quad \Omega \;, \\ \mathbf{u}_{\epsilon} &= \; 0 \quad \text{on} \quad \partial \Omega \; \;. \end{split}$$

Let  $V_{\epsilon}$  be defined by

$$V_{\epsilon} = f_{\epsilon} e^{-u_{\epsilon}}$$

so that  $-\Delta u_{\epsilon} = V_{\epsilon} e^{u_{\epsilon}}$ . We claim that, for an appropriate choice of  $d_{\epsilon}$ , we have

$$\|V_{e}\|_{L^{\infty}} \leq C$$

(46) 
$$\int_{\Omega} e^{u_{\epsilon}} \leq C.$$

while  $u_{\epsilon}(\mathbf{a}_{\epsilon}) \rightarrow +\infty$ .

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<u>Verification of (45)</u>. Let  $v_{\epsilon}$  be the solution of

$$\begin{cases} -\Delta \mathbf{v}_{\epsilon} = \mathbf{f}_{\epsilon} & \text{in } \mathbf{B}_{\mathbf{d}}(\mathbf{a}_{\epsilon}), \\ \mathbf{v}_{\epsilon} = \mathbf{0} & \text{on } \partial \mathbf{B}_{\mathbf{d}}(\mathbf{a}_{\epsilon}) \end{cases}.$$

By the maximum principle we have  $v_{\epsilon} \leq u_{\epsilon}$  in  $B_{d_{\epsilon}}(a_{\epsilon})$  so that

$$\|\mathbf{V}_{\epsilon}\|_{\mathbf{L}^{\varpi}} = \|\mathbf{f}_{\epsilon}\mathbf{e}^{-\mathbf{u}_{\epsilon}}\|_{\mathbf{L}^{\varpi}} \leq \frac{4A}{\epsilon^{2}} \|\mathbf{e}^{-\mathbf{v}_{\epsilon}}\|_{\mathbf{L}^{\varpi}(\mathbf{B}_{\epsilon}(\mathbf{a}_{\epsilon}))}.$$

But  $\mathbf{v}_{\epsilon}$  is given explicitly by

$$\mathbf{v}_{\epsilon} = \begin{cases} -\frac{4A}{\epsilon^2} \mathbf{r}^2 + \alpha_{\epsilon} & 0 \le \mathbf{r} < \epsilon \\ \\ 2A \log(\frac{d}{\mathbf{r}}) & \epsilon < \mathbf{r} < d_{\epsilon} \end{cases}$$

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where  $\mathbf{r} = |\mathbf{x}-\mathbf{a}_{\epsilon}|$  and  $\alpha_{\epsilon} = \mathbf{A} + 2\mathbf{A}\log(\frac{\mathbf{d}_{\epsilon}}{\epsilon})$ . Thus  $\|\mathbf{e}^{-\mathbf{v}_{\epsilon}}\|_{\mathbf{L}^{\infty}(\mathbf{B}_{\epsilon}(\mathbf{a}_{\epsilon}))} = \mathbf{e}^{\mathbf{A}-\alpha_{\epsilon}} = (\frac{\epsilon}{\mathbf{d}_{\epsilon}})^{2\mathbf{A}}$ . Hence (45) holds provided (47)  $\frac{1}{\epsilon^{2}}(\frac{\epsilon}{\mathbf{d}_{\epsilon}})^{2\mathbf{A}} \leq \mathbf{C}$ .

Verification of (46). Let G be the half-plane

$$G = \{(x_1, x_2) \in \mathbb{R}^2; x_1 > 0\}.$$

Let  $w_{\ell}$  be the solution of

$$\begin{cases} -\Delta \mathbf{w}_{\epsilon} = \mathbf{f}_{\epsilon} & \text{in } \mathbf{G}, \\ \mathbf{w}_{\epsilon} = \mathbf{0} & \text{on } \partial \mathbf{G}. \end{cases}$$

By the maximum principle we have  $u_{\epsilon} \leq w_{\epsilon}$  in  $\Omega$  and thus

 $\int_{\Omega}^{u_{\epsilon}} e^{u_{\epsilon}} \leq \int_{\Omega}^{w_{\epsilon}} e^{w_{\epsilon}}.$ 

But  $w_{\epsilon}$  is given explicitly by

$$\mathbf{w}_{\epsilon} = \begin{cases} -\frac{4A}{\epsilon^2} |\mathbf{x} - \mathbf{a}_{\epsilon}|^2 + \beta_{\epsilon} + 2A \log |\mathbf{x} - \mathbf{a}_{\epsilon}'| & \text{if } |\mathbf{x} - \mathbf{a}_{\epsilon}| < \epsilon \\ 2A \log \left[\frac{|\mathbf{x} - \mathbf{a}_{\epsilon}'|}{|\mathbf{x} - \mathbf{a}_{\epsilon}|}\right] & \text{otherwise} \end{cases}$$

where  $a_{\epsilon}' = -a_{\epsilon}$  and  $\beta_{\epsilon} = A - 2A \log \epsilon$ . We have

$$\mathbf{w}_{\epsilon}(\mathbf{x}) \leq \mathbf{C} + 2\mathbf{A} \log(\frac{\mathbf{d}_{\epsilon}}{\epsilon}) \quad \text{if } |\mathbf{x}-\mathbf{a}_{\epsilon}| < \epsilon$$

 $(\text{since } |\mathbf{x}-\mathbf{a}_{\epsilon}'| < |\mathbf{x}-\mathbf{a}_{\epsilon}| + 2\mathbf{d}_{\epsilon} \leq \epsilon + 2\mathbf{d}_{\epsilon} \leq 3\mathbf{d}_{\epsilon}),$ 

$$w_{\epsilon}(x) \leq C + 2A \log(\frac{d_{\epsilon}}{|x-a_{\epsilon}|})$$
 if  $\epsilon \leq |x-a_{\epsilon}| < d_{\epsilon}$ 

(since  $|\mathbf{x}-\mathbf{a}_{\ell}'| < 3d_{\ell}$ ) and

$$\mathbf{w}_{\epsilon} \leq \mathbf{C}$$
 if  $|\mathbf{x}-\mathbf{e}_{\epsilon}| \geq \mathbf{d}_{\epsilon}$ 

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$$\int_{\Omega} e^{\mathbf{W}_{\varepsilon}} \leq C \ \epsilon^{2} \left(\frac{d_{\epsilon}}{\epsilon}\right)^{2A} + C \int_{\epsilon}^{d_{\epsilon}} \left(\frac{d_{\epsilon}}{r}\right)^{2A} r \ dr + C$$
$$\leq C \ \epsilon^{2} \left(\frac{d_{\epsilon}}{\epsilon}\right)^{2A} + C.$$

Hence (47) and (46) can be achieved by choosing  $d_{\epsilon} = \epsilon^{1-(1/A)}$ . Finally we have  $u_{\epsilon}(a_{\epsilon}) \ge v_{\epsilon}(a_{\epsilon}) = \alpha_{\epsilon} \ge 2A \log(\frac{d_{\epsilon}}{\epsilon})^2 \longrightarrow +\infty$  as  $\epsilon \longrightarrow 0$ . Note that in this Example  $\int_{\Omega} V_{\epsilon} e^{u_{\epsilon}} = 4A\pi$  can be made arbitrarily close to  $4\pi$ , showing once more that assumption (10) in Corollary 3 is sharp.

4. One may combine the techniques of Sections III.1 and Section III.2. Assume for example that all the assumptions of Corollary 6 hold with 1 and in addition

 $|V_n(x)| \leq W(x)$  in some fixed neighbourhood of  $\partial \Omega$ 

with  $W \in L^p$ . Then  $\|u_n\|_{T^{\infty}} \leq C$ .

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